

# A Study of Translational Diffusion Constants of Paramagnetic Species In Solution

by

Hasan Bakor Abdullah Balkhoyor

A Thesis Presented to the

FACULTY OF THE COLLEGE OF GRADUATE STUDIES

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DHAHRAN, SAUDI ARABIA

In Partial Fulfillment of the  
Requirements for the Degree of

**MASTER OF SCIENCE**

In

**CHEMISTRY**

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**Balkhoyor, Hasan Bakor Abdullah, M.S.**

**King Fahd University of Petroleum and Minerals (Saudi Arabia), 1993**

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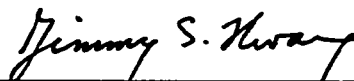
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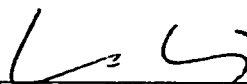
**COLLEGE OF GRADUATE STUDIES**

This thesis, written by **HASAN BAKOR BALKHOYOR** under the direction of his Thesis Advisor and approved by his Thesis Committee, has been presented to and accepted by the Dean of the College of Graduate Studies, in partial fulfillment of the requirements for the degree of **MASTER OF SCIENCE in Chemistry**.


**Thesis Committee**



Prof. Jimmy S. Hwang  
Thesis Advisor



Prof. Uwe K. A. Klein  
Member



Prof. Samuel J. Lyle  
Member

  
Dr. Abdul-Rahman A. Al-Arfaj  
Department Chairman

  
Dr. Ala H. Al-Rabeh  
Dean, College of Graduate Studies

Date: Aug 14, 1993



**THIS THESIS IS DEDICATED TO MY**  
***PARENTS,***  
***WIFE,***  
***and***  
***CHILDREN.***



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## خلاصة الرسالة

اسم مقدم الرسالة : حسن بكر عبدالله بالخير

عنوان الرسالة : دراسة معامل الانتشار الانتقالي للمركبات ذات الخواص المقاطيسية في المحاليل

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في هذا البحث تمت دراسة معامل الانتشار الانتقالي لمركبين لهما خواص مقاطيسية ، الأول : ثنائي كلورو -  
ثين (١- بتنادينيل حلقي) فثاديوم ، ويختصر (  $Cp_2VCl_2$  ) في الكلوروفورم كمذيب، والثاني : S -  
ميثيل - N - O - مئوكسي - مالبسيلدين - هايدرازين - كاريودايشوتو - فنانشولين - لكسي الفثاديوم ،  
ويختصر VO[5-MeO Sal SB] (Phen)، في ثنائي كلورو ميثان كمذيب. وذلك عند درجة حرارة  
الغرفة. وقد استخدمت طريقة الانتشار خلال الأنبوب الشعرية لأيجاد هذه المعاملات عن طريق تتبع شدة الطيف  
المقاطيسي مع مرور الزمن بواسطة جهاز طيف الطنين المقاطيسي الإلكتروني . ويرسم علاقة للتغير بين  
الزمن وشدة الطيف ثم إيجاد أفضل منحنى نظري يحاكي المنحنى التجريبي من خلال حل القانون الانتشار  
الثاني ل-Fick's والتعويض في هذا القانون لأيجاد المعاملات بواسطة برنامج كمبيوتر كتب لهذا  
الغرض. وبأستخدام علاقة Stokes - Einstein التي تربط بين معامل الانتشار الانتقالي و  
اللزوجة وجد نصف القطر الحركي للجزيئات ومنه حسبت حجوم هذه الجزيئات. فكان معامل الانتشار الانتقالي  
للمركب الأول  $0.520 \times 10^{-5}$  سم<sup>٢</sup>/ث ، وحجمه  $250 \times 10^{-8}$  سم<sup>٣</sup>. ومعامل الانتشار للمركب الثاني  
 $655 \times 10^{-5}$  سم<sup>٢</sup>/ث وحجمه  $275 \times 10^{-8}$  سم<sup>٣</sup>.

في الجزء الثاني من الرسالة استخدمت أربع مجموعات مختلفة لمعاملات مقاطيسية للمركب الأول ذكرت في  
الدوريات العالمية لتحقيقها على قيم عرض الخطوط للطيف المقاطيسي عند درجات حرارة مختلفة في مادة  
التولوين ، لأيجاد معامل التباين الدوراني (N) وكذلك محور الدوران لهذا المركب. وقد وجد أن مجموعة  
واحدة من المعاملات المقاطيسية حققت قيم عرض خطوط الطيف المقاطيسي بمعامل تباين دوراني  
(N)  $2.7 \pm 0.5$  وذلك على المحور  $Z' = Y$  ، وتطابقت هذه القيمة مع نموذج  
Stokes-Einstein ومعادلة القيمة المانونة ( Allowed - Values )  
( Equation ) و من معامل التباين الدوراني حسب زمن الدوران الذي يرتبط بعلاقة رياضية مع اللزوجة  
والحرارة والحجم الجزيئي ومعامل التجانب بين جزيئات المذاب والمذيب (k) ، الذي حسبت قيمته 0.15  
وهذه القيمة تدل على أن المركب أكثر اقتراباً من المركب التومونجي فناديل اسيتل اسيتونات (VOAA)  
( والذي يبلغ مقدار التجانب له في التولوين 0.53 ، ويؤيد ذلك معامل الانصاق (S)  
الذي يأخذ القيم (  $0 < S < 1$  ) فعندما يكون S يساوي صفر فنلك يدل على الانزلاق  
الكامل ، والواحد الانصاق الكامل ، وبحساب معامل الانصاق للمركب الأول وجد أن قيمته 0.0875 .  
وبمراجعة أشكال هذه الجزيئات يمكن تقدير ذلك فكلماً كان الجزيئ منتظماً قل معامل التجانب بين جزيئات  
المذيب والمذاب وكذلك قل معامل الانصاق ، والعكس كلما كان الشكل الجزيئ غير منتظم ازداد معامل التجانب  
و كذلك معامل الانصاق .

## ABSTRACT

The translational diffusion constants of two paramagnetic species, dichlorobis( $\eta$ -cyclopentadienyl)vanadium, abbreviated as  $(Cp_2VCl_2)$  in chloroform, and S-methyl-N-5-methoxy-salicylidenehydrazine-carbodithioatophenanthroline-oxovanadium(IV), abbreviated as  $VO[5-MeO-Sal SB](Phen)$  in methylene chloride were measured at room temperature ( $25 \pm 1$  °C) by a capillary diffusion method. The translational diffusion processes were monitored by electron spin resonance spectral intensities as a function of time. The total spectral intensities were plotted as a function of time and from the best least squares fit a single value of the translational diffusion coefficient was obtained. A computer program based on Fick's second law of diffusion for the capillary diffusion cell has been written to get a single value of the diffusion constant. Using Stokes-Einstein for translational diffusion, the hydrodynamic radii and the molecular hydrodynamic volumes of the two paramagnetic species in solution were found to be  $0.520 \times 10^{-3}$  and  $0.655 \times 10^{-3}$  cm<sup>3</sup>/s, and 250 and 275 Å<sup>3</sup>, respectively.

From an analysis of the temperature dependent ESR line widths of  $Cp_2VCl_2$  in toluene in the motional narrowing region by using four sets of the magnetic parameters available in the literature, only the set by Peterson and Dahl could explain the line width data. It was found that  $Cp_2VCl_2$  was undergoing anisotropic reorientation with  $N = 2.7 \pm 0.5$  at an axis  $z' = Y$ , where  $N$  is the ratio of  $R_{\parallel}/R_{\perp}$  and  $R_{\parallel}$  is the rotational diffusion constant along the molecular  $Z$  axis and  $R_{\perp}$  is the rotational diffusion constant perpendicular to the molecular  $Z$  axis. This value is supported by the Stokes-Einstein model and the allowed-values equation.

In addition, the rotational correlation times as a function of temperature were determined from the line width data. The anisotropic interaction parameter  $\kappa$ , determined from the plot of  $\tau_c$  vs  $\eta/T$  was found to be 0.150. The anisotropic interaction parameter is a measure of the anisotropy of intermolecular interactions. The values show that  $Cp_2VCl_2$  in toluene is more slip than VOAA in toluene which is more stick and its  $\kappa$  is 0.53. The Stickiness factor ( $S$ ) is calculated to be 0.0875 for  $Cp_2VCl_2$ , which is close to the slip limit ( $S=0$ ). Therefore, it seems that  $\kappa$  value decreases with decreasing solute corrugation.

# **CHAPTER ONE**

## **INTRODUCTION**

## CHAPTER 1

### INTRODUCTION

#### 1.1 INTRODUCTION

The purpose of this research is to study the translational motion of two paramagnetic species in solution by the capillary diffusion method. The first complex is S-methyl-N-5-methoxy salicylidene hydrazine carbodithioatophenanthroline oxovanadium(IV), in methylene chloride, written simply as, VO[5-MeO-Sal SB](Phen), which has a monomeric octahedral structure as shown in Fig.(1.1) [1]. The second complex is dichlorobis( $\eta$ -cyclopentadienyl)vanadium in chloroform, written as  $\text{Cp}_2\text{VCl}_2$ , its structure shown in Fig.(1.2). The translational motion will be measured at room temperature by a capillary diffusion method [2,3,4,5]. It has been previously shown that the translational diffusion processes can be monitored by electron spin resonance spectral intensities as a function of time [6]. The total spectral intensities are plotted as a function of time and from the best least-squares fit a single value of translational diffusion coefficient can be obtained [3]. Using the Stokes-Einstein equation the hydrodynamic radii of the two paramagnetic species in solution can be calculated.

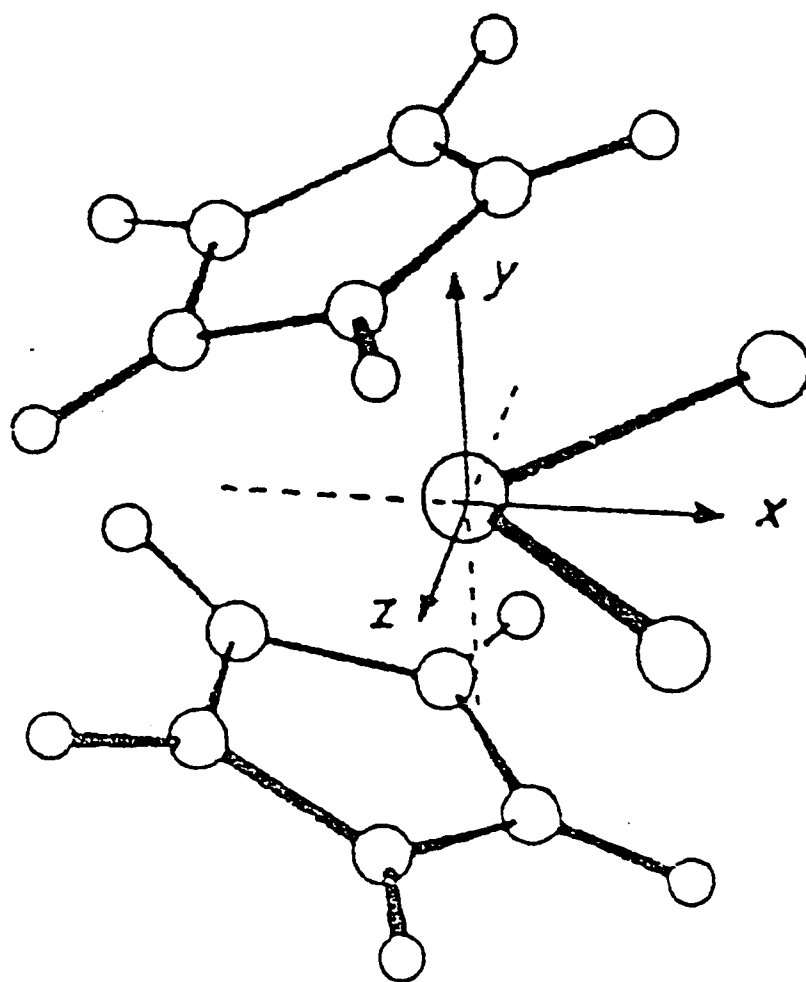
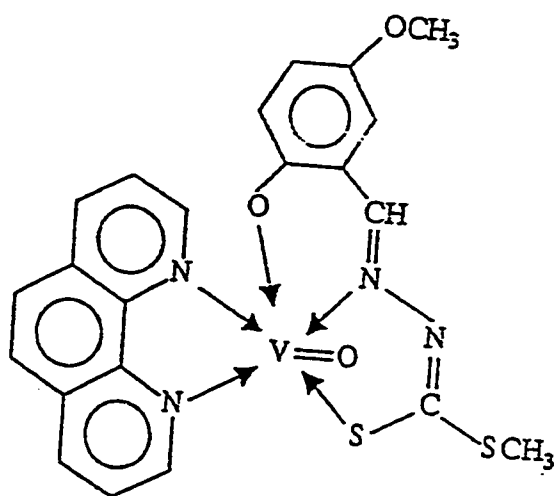


Fig. 1.1:  $\text{Cp}_2\text{VCl}_2$



**Fig. 1.2: VO[5-MeO-Sal SB](Phen)**

The translational diffusion coefficient  $D$  is related to the viscosity by the Stokes-Einstein equation

$$D = \frac{kT}{6\pi\eta r_0} \quad (1.1)$$

where  $k$  is the Boltzmann constant,  $T$  is the absolute temperature,  $\eta$  is the shear viscosity, and  $r_0$  is the hydrodynamic radius of an equivalent spherical volume [7], from which the molecular volume of the paramagnetic species can be determined. In the past the molecular volume has been determined using the Dreiding model [8], a computer program (Desktop Molecular Modeller) [1], and experimentally using such techniques as the porous disk method [9], and the electron spin resonance/spin exchange method [10].

The value of  $r_0$  which determined experimentally by the capillary diffusion method will be used in recalculating the  $\kappa$  value from the correlation time,  $\tau_1$ , which is analyzed from molecular reorientation of paramagnetic species in solution as a function of temperature [1]. For spherical top or linear molecule this correlation time can be expressed as [11],

$$\tau_R = \left(\frac{4\pi r_0^3}{3}\right) \left(\frac{\kappa}{k_B}\right) \left(\frac{\eta}{T}\right) \quad (1.2)$$

Where  $\eta$  is the coefficient of shear viscosity of the solvent [12], and  $\kappa$  is an experimentally determined dimensionless coupling parameter called the anisotropic interaction parameter which measures the coupling of rotational motions of the probe to the translation modes of the fluid [11].

From the electron spin resonance relaxation studies [13], which have been analyzed in terms of the hydrodynamic free space model for molecular relaxation in liquids at various microwave frequencies [8,13,14], the analysis of the temperature dependent ESR line shapes in the motional narrowing and slow tumbling regions, the stickiness factors  $S$  were measured [15].

The stickiness factor is independent of molecular geometry and is zero in the slip and one in the stick limit,  $0 < S < 1$ . The stickiness factor can be calculated from the relationship

$$S = \frac{\kappa - C_{slip}^{hyd}}{1 - C_{slip}^{hyd}} \quad (1.3)$$

where  $C_{slip}^{hyd}$  is the empirical factor which depends on the non-hydrodynamic character of the friction and on the shape of the probe molecules [11], then from the recalculated value of  $\kappa$  we can obtain a more accurate value of  $S$ .



In the remainder of this chapter we give an overview of the transport process. The state of a system under diffusion is defined, then Fick's law of diffusion is set up mathematically and a solution of this law is derived. In Chapter II the apparatus used and its function in the experiment are described. In Chapter III we describe the experimental procedure used to set up the capillary diffusion cell using the ESR technique. In Chapter IV the analysis of the paramagnetic parameters are given. In Chapter V the conclusions of the experimental results and discussion are given.

## 1.2 TRANSPORT PHENOMENA

Transport phenomenon arises because any system in nonequilibrium will always try to approach equilibrium. A system is in equilibrium if the intensive state variables or parameters (temperature and density, for example) are constant at all times and have the same magnitude at all positions in the system.

From this definition follow other familiar properties of equilibrium state: at equilibrium, *gradients* of the parameters are absent, spontaneous processes do not occur, and there is no net exchange of matter or energy between the system and the surroundings, or between one part of the system and another.

If a system is in a state in which these conditions are not satisfied, the system is in nonequilibrium state. Suppose an equilibrium system is perturbed to a nonequilibrium state by some external or internal disturbance. At least one of the system parameters becomes a function of position. When the disturbance is removed, an irreversible decay process occurs spontaneously, and the system advances through a series of nonequilibrium state until equilibrium is reached.

The decay process is a transport process, the name arising because some quantity is transferred throughout the system as the system attempts to make all parameters independent of position.

A spatial dependence of the parameters is thus associated with a *flux* of some kind. It is logically assumed that the relation between a gradient  $X$  and its conjugate flux  $J$  is linear thus,

$$J = L X \quad (1.4)$$

If Eq. (1.4) is valid at all times, the system is said to be "close to equilibrium" at all times.

The rate at which the system approaches equilibrium is determined by the proportionality constant of Eq.(1.4) that is by the *transport coefficient*  $L$ .

It is a matter of experience that transport processes taking place in a system with  $n$  independent gradients can be represented by the relation,

$$J = \sum_{k=1}^n L_{ik} X_k \quad (i=1,2,3,\dots,n) \quad (1.5)$$

In the nomenclature of nonequilibrium thermodynamics,  $X$  is called a thermodynamic driving force, and  $L$  a phenomenological coefficient.

Equation (1.5) can be written in matrix notation,

$$\begin{bmatrix} J_1 \\ J_2 \\ \vdots \\ J_n \end{bmatrix} = \begin{bmatrix} L_{11} & L_{12} & \dots & L_{1n} \\ L_{21} & L_{22} & \dots & L_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ L_{n1} & L_{n2} & \dots & L_{nn} \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \\ \vdots \\ X_n \end{bmatrix}$$

which expresses the fact that if the off-diagonal elements of  $L$  exist, then the  $i$ th flux depends in a linear manner on the  $k$ th force as well as on its conjugate  $X_i$ . The relation between  $J_i$  and  $X_k$  is called thermodynamic coupling.

### 1.2.1 Uncoupled Transport Processes

When  $L_{ik} = 0$  ( $i \neq k$ ), Eq.(1.5) becomes  $J_i = L_{ii}X_i$ , a phenomenological representation of familiar empirical laws,

(1) Fourier's law of heat conduction,

$$\mathbf{q} = -\lambda \nabla T \quad (1.6)$$

The vector  $\mathbf{q}$  is the heat flux, the coefficient  $\lambda$  is the coefficient of thermal conductivity, and  $\text{grad } T$  is the gradient of temperature.

(2) Fick's law of diffusion,

$$J_i = -D \nabla C_i \quad (1.7)$$

where  $J_i$  is the flux of chemical species  $i$ ,  $C_i$  is the concentration of  $i$ , and  $D$  is the diffusion coefficient.

(3) Newton's law of viscosity,

$$\Pi_{yx} = -\eta \frac{\partial u_x}{\partial y} \quad (1.8)$$

Here  $\Pi_{yx}$  represents a momentum flux in the y direction caused by a fluid flow in the x direction with velocity component  $u_x$ ,  $\eta$  is the coefficient of shear viscosity [16].

#### 1.2.2 Coupled Transport Processes

The coupled phenomena arising if  $L_{ik} \neq 0$ , *thermal diffusion* of gases is an example of this processes in which the diffusion of gases depends not only on the concentration difference but also on the temperature difference. Also the *thermoelectricity* with coupling between flows of heat and of electrical current is another example [16].

#### 1.2.3 Continuous and Discontinuous System

The fluid can be regarded as a continuum and the intensive state variables as continuous function of space and time. Such a system is called continuous. Also the phenomenological or transport coefficients, for both coupled and single processes, appertain to the fluid itself. At a microscopic level they could be described by the temperature

and pressure of the system, together with the microscopic properties of the components of the fluid, for example by the molecular dimensions and the intermolecular potential function, *thermal diffusion* is an example here for continuous system in which the vessel plays no part in the effect apart from containing the fluid. In a discontinuous system, the vessel plays an essential role, the fluid-vessel system is not a continuum, and the state variable are not continuous function of space because of discontinuities of the state variables. Here the phenomenological or transport coefficients are now function of the fluid and of the vessel [16].

#### 1.2.4 Transport Processes At The Steady State

A process which involves time-dependent parameters is not easy to setup experimentally or to analyze theoretically. The process becomes far simpler to deal with if the state parameters are time independent. This can be achieved when the parameter gradients are externally constrained to constant values. In fact, the empirical transport laws, Fourier's law, etc., were conceived for this case.

Today, as well, most transport experiments are planned so that the state parameters are independent of time. The simplification applies especially to coupled phenomena. It is an experimental fact that, the case of a system with  $n$  gradients present, if one gradient is constrained by external

means, all gradients become constant in the course of time. The state that the system eventually reaches is called the *steady or stationary state*. The steady state has considerable theoretical significance. The steady state is the time-invariant ideal state for an open system. It is the counterpart of the time-invariant ideal state for a closed system-*the equilibrium state* [16].

### 1.3 Mass-Transfer

#### 1.3.1 Modes of Mass Transfer

The movement of material from one location in solution to another, arises from differences in electrical or chemical potential at the two locations, or from movement of a volume element of solution. The modes of mass transfer are:

(1) *Migration*. Movement of a charged body under the influence of an electric field ( a gradient of electrical potential).

(2) *Diffusion*. Movement of a species under the influence of a gradient of chemical potential ( a concentration gradient).

(3) Convection. Stirring or hydrodynamic transport.

Generally fluid flow occurs because of natural convection ( convection caused by density gradients) and forced convection, and may be characterized by *laminar flow* and *turbulent flow* [17].

### 1.3.2 Laminar And Turbulent Flow Through Tubes

The flow through tubes can be laminar or turbulent, depending on the magnitude of the dimensionless Reynolds number  $R_e$ , which is given by

$$R_e = \frac{\rho v d}{\eta} \quad (1.9)$$

where  $v$  is the fluid velocity (cm/sec),  $d$  is the tube diameter (cm),  $\rho$  is the fluid density (g/cm<sup>3</sup>), and  $\eta$  is the fluid viscosity (poise). Physically the Reynolds number reflects the relative strength of the inertial and viscous forces in the moving fluid. Inertial forces increase with density and with square of velocity (  $\rho v^2$  ). Viscous forces increase with viscosity and velocity and decrease with the increasing tube diameter (  $\eta v/d$  ). The Reynolds number is simply the ratio of these two terms.

At low Reynolds number (  $R_e < 2100$  ) the flow called *laminar* or *viscous*, above this critical Reynolds number *turbulence flow* occurs [18].



#### 1.4 The Diffusion Process

Mass transport on a molecular scale is called *diffusion* and on the scale of the apparatus is termed *bulk motion*, or more precisely, *convection* [18].

Diffusion is a process which leads to an equalization of concentrations within a single phase [4], in which matter is transported from one part of a system to another as a result of random molecular motions "*random walk*" [19] , or *Brownian motion* [7].

##### 1.4.1 Types of Diffusion Coefficients

Fick studied the interdiffusion of two solutions of the same components but with different concentrations, and also diffusion from a solid into its saturated solution. His diffusion coefficients were therefore *interdiffusion coefficients* or *mutual diffusion coefficients*. In a pure liquid the individual particles move with random motion and have an equal opportunity of taking up any point in the total space occupied by the liquid. If it were possible to label a particle without otherwise changing its properties and to follow its motion through the unlabelled molecules, a *self diffusion coefficient* could be defined. Equally, in a multicomponent system, some labelled molecules of component 1 could be introduced so as to form a concentration gradient of labelled and unlabelled species 1 with an otherwise uniform environ-

ment. Application of Fick's law to the interdiffusion of labelled and unlabelled molecules of species 1 gives the *intradiffusion coefficient* of this species within the solution. A self diffusion coefficient is a special case of intradiffusion in a system which contains only the two distinguishable (by isotopic labelling, nuclear magnetic vector orientation, etc.) forms of the same chemical species. An other related experiment comprises adding at some point in space a trace of labelled species 1 to an otherwise homogeneous medium which may, or may not, be multicomponent, and may, or may not, contain unlabelled species 1. Attention is focused on the decay of the initially formed concentration gradient of the labelled species when it is at low concentration in an otherwise homogeneous environment. The resulting diffusion coefficient is termed a *tracer diffusion coefficient*. Also, the orientation of any molecule undergoes random fluctuations and so rotational motion also occurs. This motion ensures that, in the absence of any external field, the molecules do not take up any preferred orientation. When an external field is present the extent of rotational order produced depends upon a balance between the magnitude and direction of the applied field and the intensity of the rotational motion, which in turn depends upon the thermal energy of the system. Removal of the field causes the degree of rotational order which had been induced to decay back to zero at a characteristic rate, then a *rotational diffusion*

*coefficient* can be defined. All these diffusion coefficients can be related to the translational Brownian motion of the molecules which arises from random fluctuations in the positions of molecules in space as a result of the decay of a concentration gradient, and a *translational diffusion coefficient* can be termed for this motion [5].

#### 1.4.2 Basic Hypothesis Of Mathematical Theory

Transfer of heat by conduction is also due to random molecular motion, and there is an obvious analogy between the two processes. This was recognized by Fick, who first put diffusion on a quantitative basis by adopting the mathematical equation of heat conduction derived some years earlier by Fourier. The mathematical theory of diffusion in isotropic substances is therefore based on the hypothesis that the rate of transfer of diffusing substance through unit area of a section is proportional to the concentration gradient measured normal to the section,

$$F = - D \frac{\partial C}{\partial X} \quad (1.10)$$

where  $F$  is the rate of transfer per unit area of section,  $C$  the concentration of diffusing substance,  $X$  the space coordinate measured normal to the section, and  $D$  is the diffusion coefficient.

In some cases, e.g. diffusion in dilute solutions,  $D$  can reasonably be taken as constant, while in other, e.g. diffusion in high polymers, it depends very markedly on concentration. If  $F$ , the amount of material diffusing, and  $C$  the concentration, are both expressed in terms of the same unit of quantity, e.g. grams or grams molecules, then it is clear from Eq. (1.10) that  $D$  is independent of this unit and has dimensions ( $\text{length}^2 \text{ time}^{-1}$ ), e.g.  $\text{cm}^2 \cdot \text{sec}^{-1}$ . The negative sign in Eq.(1.10) arises because diffusion occurs in the direction opposite to that of increasing concentration.

It must be emphasized that the statement expressed mathematically by Eq. (1.10) is in general consistent only for an isotropic medium, whose structure and diffusion properties in the neighborhood of any point are the same relative to all directions. Because of this symmetry, the flow of diffusing substance at any point is along the normal to the surface of constant concentration through the point. This need not be true in an anisotropic medium for which the diffusion properties depend on the direction in which they are measured [19].

### 1.5 Differential Equation of Diffusion

The fundamental differential equation of diffusion in an isotropic medium is derived from Eq.(1.10) as follows. Consider an element of volume in the form of a rectangular parallelepiped whose sides are parallel to the axes of coordinates and are of lengths  $2dx$ ,  $2dy$ ,  $2dz$ . Let the center of the element be at  $P(x, y, z)$ , where the concentration of diffusion substance is  $c$ . Let ABCD and A'B'C'D' be the faces perpendicular to the axes of X as in Fig.(1.3).

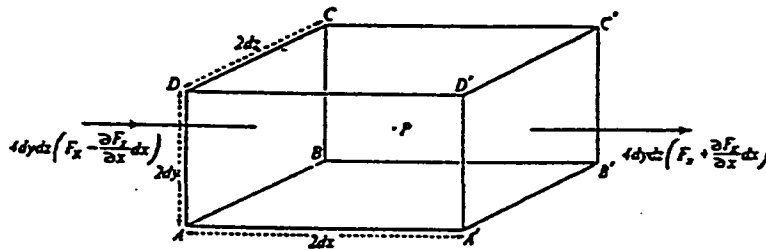


Fig.1.3: Element of Volume

Then the rate at which diffusing substance enters the element through the face ABCD in the plane  $(x - dx)$  is given by

$$4 \, dydz \left( F_x - \frac{\partial F_x}{\partial x} dx \right)$$

where  $F_x$  is the rate of transfer through unit area of the corresponding plane of through P. Similarly the rate of loss of diffusion substance through the face A'B'C'D' is given by

$$4 \, dydz \left( F_x + \frac{\partial F_x}{\partial x} dx \right)$$

The contribution to the rate of increase of diffusion substance in the element from these two faces is thus equal to

$$- 8 \, dx dy dz \frac{\partial F_x}{\partial x}$$

Similarly from the other faces we obtain

$$- 8 \, dx dy dz \frac{\partial F_y}{\partial y} \quad , \quad - 8 \, dx dy dz \frac{\partial F_z}{\partial z}$$

But the rate at which the amount of diffusing substance in the element increases is also given by

$$\delta \, dx dy dz \, \frac{\partial C}{\partial t}$$

and hence we have immediately

$$\frac{\partial C}{\partial t} + \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z} = 0 \quad (1.11)$$

If the diffusion coefficient is constant,  $F_x, F_y, F_z$  are given by Eq.(1.10), and Eq.(1.11) becomes

$$\frac{\partial C}{\partial t} = D \left( \frac{\partial^2 C}{\partial x^2} + \frac{\partial^2 C}{\partial y^2} + \frac{\partial^2 C}{\partial z^2} \right) \quad (1.12)$$

reducing simply to

$$\frac{\partial C}{\partial t} = D \, \frac{\partial^2 C}{\partial x^2} \quad (1.13)$$

if diffusion is one dimensional, i.e. if there is a gradient of concentration only along the X-axis, expressions of Eqs.(1.10 & 1.13) are usually referred to as Fick's first and second laws of diffusion [19,20].

## 1.6 Solution Methods of the Diffusion Equation when the Diffusion Coefficient is Constant

General solution of the diffusion equation can be obtained for a variety of initial and boundary conditions provided the diffusion coefficient is constant. Such a solution usually has one of two standard forms. Either it is comprised of a series of error function or related integrals, in which case it is most suitable for numerical evaluation at small times, i.e. in the early stages of diffusion, or it is the form of a trigonometrical series which converges most satisfactorily for large values of time. Another solution used the Laplace transform, is essentially an operator method by which both types of solution may be obtained. It is the most powerful of the three, particularly for more complicated problems [19].

### 1.6.1 Method of Separation of Variables [19,20,21,22,23]

A standard method of obtaining a solution of a partial differential equation is to assume that the variables are separable. Thus we may attempt to find a solution of Eq. (1.13) by putting

$$C(x, t) = F(x) \cdot f(t) \quad (1.14)$$



where  $F(x)$  and  $f(t)$  are function of  $x$  and  $t$  respectively .

Therefore

$$\begin{aligned}\frac{\partial C}{\partial t} &= F(x) \cdot f'(t) \\ \frac{\partial C}{\partial x} &= F'(x) \cdot f(t) \\ \frac{\partial^2 C}{\partial x^2} &= F''(x) \cdot f(t)\end{aligned}\tag{1.14 a}$$

substitution of Eq.(1.14 a) into Eq.(1.13) gives

$$F(x) f'(t) = D F''(x) \cdot f(t)\tag{1.14 b}$$

By separation of variables of Eq.(1.14 b) we get

$$\frac{f'(t)}{D f(t)} = \frac{F''(x)}{F(x)}\tag{1.14 c}$$

Since the left and right sides of Eq.(1.14 c) are respectively function of  $t$  only and  $x$  only, the equation can be satisfied if each side is separately equal to the same constant, let it be  $(-k^2)$ , and then we have two equations

$$F''(x) = -k^2 F(x)\tag{1.14 d}$$

and

$$f' = -k^2 D f(t) \quad (1.14 \text{ e})$$

In the capillary tube diffusion method, which will be discuss in the next chapter, the initial and boundary conditions for a tube close at ( $x = 0$ ) and open at ( $x = a$ ) are

$$\text{At } t=0, c=c_0 \text{ for } 0 < x < a, c=0 \text{ for } x > a$$

$$\text{At } t>0, c=0 \text{ at } x=a \text{ and } \frac{\partial c}{\partial x}=0 \text{ at } x=0$$

To solve Eq.(1.14 d),

$$F''(x) + k^2 F(x) = 0 \quad (1.14 \text{ f})$$

the auxiliary equation for a second order linear differential equation is  $m^2 + k^2 = 0$ . Solving for  $m$  we get two complex roots  $m = \pm ik$ . Then the general solution for Eq. (1.14 f) has the form

$$F(x) = d_1 \cos kx + d_2 \sin kx \quad (1.14 \text{ g})$$

To find  $d_1$ ,  $d_2$ , and  $k$  we use the initial and boundary conditions,

$$F' = (d_1 \sin kx)(-k) + kd_2 \cos kx \quad (1.14 \text{ h})$$

but we know that at  $x = 0$ ,  $F'(x)=0$  so that  $F'(0) = kd_2 = 0$ . Therefore  $d_2 = 0$ , and

$$F(x) = d_1 \cos kx \quad (1.14 \text{ i})$$

We also have at  $x = a$ ,  $F(a) = 0$ , therefore  $F(a) = d_1 \cos ka = 0$ , but  $d_1 \neq 0$  (if so  $F(x) = 0$ ). Then  $\cos ka = 0$ , and solving for  $ka$  with respect to  $\pi$  at which the  $\cos$  function equal to zero we get

$$\begin{aligned} ka &= \frac{2n+1}{2}\pi, \\ k &= \frac{2n+1}{2a}\pi, \quad n=0,1,2,\dots \end{aligned} \quad (1.14 \text{ j})$$

Then  $F(x)$  becomes

$$F_n(x) = d_1 \sum_{n=0}^{\infty} \cos \frac{\pi(2n+1)}{2a} x \quad (1.14 \text{ k})$$

To solve the second equation, Eq. (1.14 e),

$$\frac{f'(t)}{f(t)} = -k^2 D$$

by integration

$$\begin{aligned} (\ln f(t))' &= (-k^2 D t + c_1)' \\ \ln f(t) &= -k^2 D t + c_1 \end{aligned}$$

we get for the function  $f(t)$ ,

$$f(t) = c_1 \exp(-k^2 D t) \quad (1.14 \text{ l})$$

Substituting Eqs. (1.14 k & 1.14 l) into Eq. (1.14) we get

$$C_n(x, t) = \sum_{n=0}^{\infty} B_n \exp(-k^2 D t) \cos kx \quad (1.14 \text{ m})$$

where  $B_n = c_1 d_1$ . To find it's value we use the condition  
 $c(x, 0) = 0$  at  $x > a$  and  $c(x, 0) = c_0$  at  $0 < x < a$   
 putting the value of  $k$  from Eq.(1.14 j), Eq.(1.14 m) will be

$$c(x, 0) = \sum_{n=0}^{\infty} B_n \cos \frac{\pi(2n+1)}{2a} x \quad (1.14 \text{ n})$$

which can be solved by Fourier analysis as

$$B_n = \frac{2}{a} \int_0^a c_o \cos \frac{\pi(2n+1)}{2a} x \, dx \quad (1.14 \text{ o})$$

Upon integrating we obtain

$$\begin{aligned} B_n &= \frac{4c_o}{\pi(2n+1)} \sin\left[\frac{\pi(2n+1)}{2}\right] \\ &= \frac{4c_o}{\pi(2n+1)} (-1)^n \end{aligned} \quad (1.14 \text{ p})$$

Substitution of  $B_n$  and  $k$  in Eq. (1.14 m) yields

$$c(x, t) = \sum_{n=0}^{\infty} \frac{4c_o(-1)^n}{\pi(2n+1)} \text{EXP}\left[\frac{-\pi^2(2n+1)^2Dt}{4a^2}\right] \cos \frac{\pi(2n+1)}{2a} x \quad (1.14 \text{ q})$$

By integration of Eq. (1.14 q) in the limit of  $0 < x < a$  to get the total concentration at any time  $t$  we get the final result

$$c(x, t) = \frac{8ac_o}{x} \sum_{n=0}^{\infty} \frac{(-1)^n}{\pi^2(2n+1)^2} \text{EXP}\left[\frac{-\pi^2(2n+1)^2Dt}{4a^2}\right] \sin \frac{\pi(2n+1)x}{2a} \quad (1.15)$$

From Beer's law one can relate the concentration( $c$ ) to the absorption of electromagnetic radiation by a compound and to the intensity( $I$ ) of the spectrum, so  $c(x,t)$  will be equal to  $I(t)$  (the intensity of the spectrum at any time), and  $c_0$  will be equal to the  $I_0$  (the initial intensity of the spectrum or the reference one).

In Eq. (1.15), we can replace the ratio of  $x/a$  in which  $x$  is the total length of the capillary tube ( $\ell_t$ ) or the maximum distance that the molecules can go,  $a$  is the small length of the tube ( $\ell_s$ ) that is under study. Because of the uncertainty in the length of ( $\ell_s$ ) which is under the electromagnetic radiation, a ratio  $R$  ( $R = \ell_s / \ell_t$ ) can be substituted in the equation, we can rewrite Eq. (1.15) as

$$\frac{I(t)}{I(0)} = \frac{8}{\pi^2 R} \sum_{n=0}^{\infty} (-1)^n f_n(t) \sin \left[ \frac{(2n+1)\pi R}{2} \right]$$

where

$$f_n(t) = \frac{1}{(2n+1)^2} \exp \left\{ \frac{-(2n+1)^2 \pi^2 D t}{4 \ell_s^2} \right\}$$

(1.16)

A computer program for Eq. (1.16) has been written and it is listed in Appendix 1.

**C H A P T E R     T W O**  
**A P P A R A T U S**

## CHAPTER # 2

### APPARATUS

#### 2.1 Introduction

Electron Spin Resonance spectroscopy (ESR), also called Electron Paramagnetic Resonance spectroscopy (EPR), discovered by Zavoisky in 1944, has become an essential tool for the study of the structure and dynamics of molecular systems containing one or more unpaired electrons [24].

##### 2.1.1 Systems Studied by ESR

A large number of systems have been studied by ESR:

- (1) Biological system : hemoglobin, nucleic acid, enzymes.
- (2) Chemical system : polymers, catalysts, free radicals.
- (3) Conduction electrons : alkali and alkali earth metals.
- (4) Stable free radicals : in solid and solution, biradicals triplet states.
- (5) Gases :  $\text{NO}$  ,  $\text{NO}_2$  ,  $\text{SO}_2$ .
- (6) Irradiated substances : ionic crystals solid and liquid.
- (7) Naturally occurring substance : petroleum, rubber, dolomite, coal.
- (8) Semiconductors : doped semiconductors.
- (9) Transition elements : single crystals, chelates, sandwich compound, alloys [25].



### 2.1.2 The Resonance Condition in ESR

Electron is a charged particle with angular momentum (orbital and spin) that possesses a magnetic moment,  $\mu_e$ , given by

$$\mu_e = -g \beta J \quad (2.1)$$

Here,  $J$  is the total angular momentum vector in units of  $\hbar$  where  $\hbar$  is the Planck's constant,  $g$  is a dimensionless constant (the  $g$ -value,  $g$  factor, or spectroscopic splitting factor), and  $\beta$  is a constant called the Bohr magneton. The negative sign in the equation is a consequence of negative electronic charge.

Neglecting orbital angular momentum and considering only the total spin angular momentum,  $S$ , Eq.(2.1) can be written as

$$\mu_e = -g \beta S \quad (2.2)$$

The  $g$ -value for the free electron,  $g_e$ , is 2.0023 for an unpaired electron in an organic molecule or ion,  $g$  is generally within a few percent of this value.

The approximation made in Eq.(2.2) is valid for most discussions of the ESR spectra of the organic free radicals and transition metal complexes whose orbital angular momentum can be considered to be "quenched". Treating the  $g$ -value as an experimental quantity does not harm the present discussions, since deviations of the  $g$ -value from  $g_e$  can be accounted for by introduction of spin-orbital coupling.

Magnetic moments can be detected by their interactions with magnetic fields. In zero field, the magnetic moments of unpaired electrons in a sample are randomly oriented. In the presence of a magnetic field  $H$ , electron moments assume orientations with respect to the applied field, giving rise to  $(2S+1)$  energy states (Zeeman splitting). The measurable components of  $\mu_e$  are  $(g\beta m_s)$ , where  $m_s$  is the magnetic spin quantum number, which can take the values

$$+S, +(S-1), \dots, -(S-1), -S.$$

The application of a magnetic field to an  $S = 1/2$  (or larger) system is said to remove the spin degeneracy (i.e., the equal energy values of  $m_s$  in the absence of an applied magnetic field).

The energy of an electron moment in a magnetic field is given by

$$E = - \mu_e \cdot H \quad (2.3)$$

Combining eqs.(2.2 & 2.3), the expression

$$E = g \beta H m_s \quad (2.4)$$

results, assuming the direction of the applied field defines the Z-axis. When  $S = \frac{1}{2}$  there are two energy levels, whose energy is linearly dependent on  $H$ . The separation between these energy levels (Fig. 2.1) at a particular value of the magnetic field,  $H_z$ , is

$$\Delta E = g \beta H_z \quad (2.5)$$

In an ESR experiment, an oscillating magnetic field perpendicular to  $H_0$  induces transitions between the  $m_s = 1/2$  and  $m_s = -1/2$  levels, provided the frequency,  $\nu$ , is such that the resonance condition

$$\Delta E = h \nu = g \beta H_0 \quad (2.6)$$

is satisfied. In ESR the frequency is held constant and the magnetic field is varied. At a particular value of the magnetic field,  $H_0$ , resonance absorption of energy occurs, resulting in a peak in the spectrum (Fig. 2.1). The frequencies commonly employed in ESR experiments are in the microwave region, these frequencies and magnetic field strengths for  $g_e$  resonance absorption signals are given in Table 2.1.

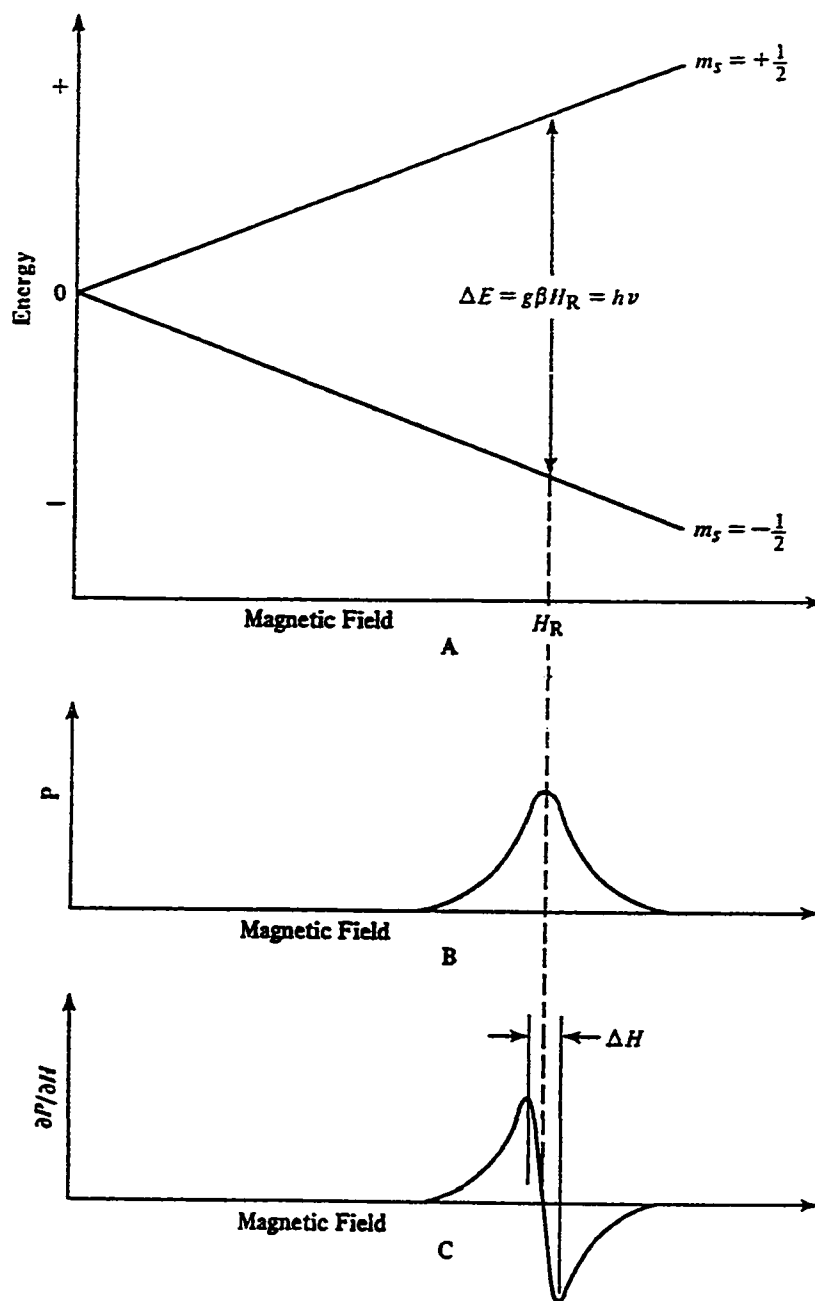


Fig. 2.1 Energy levels and spectra in ESR

**Table 2.1 Spectrometer Frequencies and  $g_e$  Resonance Field**

Designation	$\nu(\text{Hz})$	$\lambda(\text{cm})$	$\nu(\text{cm}^{-1})$	$g_e(\text{gauss})$
X-band	$\sim 9.5 \times 10^9$	3.156	0.317	3,390
K-band	$\sim 23 \times 10^9$	1.303	0.767	8,207
Q-band	$\sim 35 \times 10^9$	0.856	1.168	12,489

### 2.1.3 ESR Instrumentation

The instrumentation used in ESR spectroscopy consists of a source of electromagnetic radiation, a sample holder, and appropriate detection equipment for monitoring the amount of radiation absorbed by the sample. In ESR spectroscopy, a magnetic field provided by an electromagnet is also required. Monochromatic radiation of the various frequencies used in ESR work (Table 2.1) is obtained from *klystrons*, which are electronic oscillators producing microwave energy. Spectrometers operating at X-band (3-cm wavelength) are the ones most commonly employed. The microwave radiation is transmitted along hollow rectangular metal pipes called *waveguides*.

Figure 2.2 gives a block diagram of a simple ESR spectrometer. The sample is placed at the center of the cavity, where the magnetic vector is at a maximum. Quartz tubes (approx. 3mm o.d.) are generally used to contain solid and solution samples. Unlike the NMR technique, the sample tubes are not rotated. The magnetic field is slowly and linearly increased until the resonance condition [Eq.(2.6)] is satisfied, at which point power is absorbed by the sample and a change in current in the detector crystal is monitored. A pair of Helmholtz coils is mounted around the cavity to increase sensitivity. Feeding the coils from an oscillator superimposes a variable amplitude sinusoidal modulation on the slowly varying magnetic field.

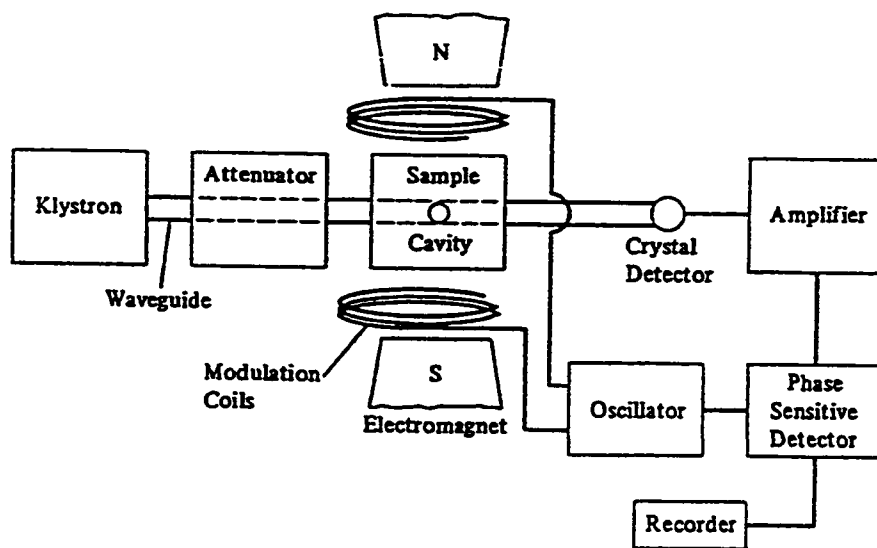


Fig. 2. 2 Block diagram of a simple ESR spectrometer



The signal detected by the phase-sensitive detection system is proportional to the slope of the ESR absorption as the magnetic field passes through resonance. The recorder then presents the first-derivative spectrum. Many spectrometers are also equipped to present second-derivative spectra.

The spectral sensitivity of ESR depends on a variety of factors, but with a response time of 1 sec, as few as  $10^{11}$  spins (approx.  $10^{-12}$  moles) can be detected with currently available spectrometers. This sort of sensitivity suggests that ESR spectroscopy would be useful for trace analysis. A minimum detectable concentration is perhaps  $10^{-9}$  M is more reasonable, where for qualitative measurement in aqueous solutions,  $10^{-7}$  M is more reasonable, where for quantitative measurements the sample concentration should be greater than about  $10^{-6}$  M. Concentrations greater than about  $10^{-4}$  M should not be used in order to avoid spin-spin (exchange) broadening. For high-resolution studies, dissolved oxygen should be removed because it is paramagnetic and can contribute to line broadening.

In any ESR experiment, it is important to monitor the microwave frequency at which the spectrometer operates and the magnetic-field range swept during the experiment. Although the frequency is constant during an experiment, the frequency available from a given klystron will vary a little with tuning of the instrument. The frequency can be determined using a built-in frequency counters.

The magnetic field can be monitored using an NMR gaussmeter or by using samples of known  $g$ -value, for instance, the DPPH free radical, for which  $g = 2.0036$ . The magnetic sweep can also be checked using Fremy's radical or oxobis(2,4-pentanedionato)vanadium(IV) [24]. More theoretical and experimental details are given in [25,26,27,28,29].

## **2.2 VARIAN E-109 ESR SPECTROMETER**

### **2.2.1 General Description**

The individual units of the E-109 spectrometer are: An electronic (operator's) console, a microwave bridge, a sample cavity with connecting waveguide, a magnet, and a magnet power supply. The system is completely solid state in design with the exception of reflux klystron tube in the microwave bridge and the cathode ray tube in the oscilloscope. The basic electronic console contains the following modules and assemblies:

- (1) The E-203B Magnetic Field Controller provides direct control and regulation of the magnetic field in the magnetic air gap.
- (2) The E-207 High Frequency Module provides magnetic field modulation frequencies of 25 kHz 100 kHz.
- (3) The E-200 oscilloscope displays the ESR signals for visual observation and the klystron power mode for spectrometer tuning.

(4) The E-080B digital drive recorder graphically displays the ESR spectrum as a function of magnetic field or time. The E-102 microwave bridge provides controlled 9.5 GHz microwave frequency to the cavity and detection of ESR signals. The absorption reference arm design permits operation at microwave power levels to one  $\mu$ W. The E-231 series sample cavity is a multipurpose rectangular cavity with an 11 mm sample access. The cavity is designed to permit up to 40 G peak-to-peak field modulation at the sample.

The magnet used is 9-inch in size with the magnet power supply providing the d.c. power to the magnet. The magnet and microwave bridge are water cooled the chilled water supplied by a mechanically refrigerated closed circuit water chiller.

To measure the field value up to 6 significant figures a Varian E-500 NMR gauss meter is used. Similarly, to get the accurate value of the frequency, HP model 5342A microwave frequency counter is used.

## **2.2.2 Data Acquisition**

### **2.2.2.1 Introduction**

The E-935 Data Acquisition system consists of the following four units:

- (1) HP 9835 computer
- (2) E-935 Data Acquisition System
- (3) HP 1350 Graphics Translator
- (4) HP 1311A CRT Display

With the use of this computer, one can store the experimental data on a cartridge and later retrieve the data for various analyses, like plotting, manipulation of data, and calculations like g-factor, etc. The data collection becomes fast and efficient with the use of computer.

### **2.2.2.2 E-935 System Software**

The software (provided by Varian Instrument Co.) is divided into four functional modules. Only one of these modules can be present in the memory of the HP 9835 computer at a time. The four functions are: calculation, manipulations, plotting, and scanning.

#### **2.2.2.3 Magnet Scan**

ESR data may be acquired while the larger magnet is swept. Magnet scan acquires ESR spectra containing 512, 1024, 2048 or 4096 data points, performs multiscan averaging, and continuously displays the current results on the display unit. Scan time may be in the range of 4-17,000 seconds.

#### **2.2.2.4 Data File Capability**

An ESR spectrum may be stored as a data file on a tape cartridge along with a file header containing descriptions of the sample. Experiment and the instrumental settings like field, frequency, power, modulation frequency modulation amplitude, number of scans, scan time, time constant, microwave power, etc. are stored in the file.

Such a data file may be retrieved from a tape cartridge, loaded into computer memory, and subsequently displayed on the display unit. Additional operations can be done on the display with the manipulation module. Also a printed catalogue of the files on a given tape cartridge may be obtained. The displayed spectrum can be immediately plotted.

#### 2.2.2.5 Data Manipulation

With the help of the manipulation module , the following manipulations of ESR spectra may be performed using the interactive capabilities of the display unit and key board:

- (1) Base line adjustment of a spectrum.
- (2) Combination of two spectra.
- (3) Detection of erroneous feature from a spectrum.
- (4) Differentiation of a spectrum.
- (5) Digital filtration of a spectrum.
- (6) Double integration of a spectrum.
- (7) X , Y axes scale and shift modifications.
- (8) Editing of the file header and recording the spectrum after editing.

#### 2.2.2.6 Tape Cartridge and File System

The tape cartridge serves two main purposes. First, the system tape cartridge contains all of the programs required for proper system operation. The second purpose is to provide for long or short term off-line storage of ESR data sets.

Before use, all tapes must be initialized as per given procedures. An initialized tape is divided into a total of

852 records of 256 bytes per record. There are two tracks per tape. Records 0-425 on the first track, while records 426-851 are on the second track. All information stored on a tape is saved in form of files. All data are contained in a particular file name. To retrieve any data later, a file name must be read to the computer.

### 2.3 VACUUM SYSTEM

For meaningful ESR line width studies, the sample should be evacuated in order to get rid of oxygen present in the air which is paramagnetic and can interfere with the line width measurements from the paramagnetic sample of interest.

A normal system manufactured by Pope Scientific Inc. was used with slight modification to suit our needs. One of the modifications is a Quick Fit arrangement to insert and remove the sample easily. Also, to measure the vacuum readily, we used a digital vacuum gauge manufactured by Granville-Phillips. Vacuum could be easily reached as low as 1 mTorr and most of the samples were sealed when the vacuum was of that order.

## 2.4 Capillary Diffusion Apparatus

This technique first introduced by Anderson and Saddington to study the diffusion of ions in solution by radioactive isotopes [30], the capillary cell is charged with the active solution and immersed in the inactive liquid for a known length of time so that normal diffusion can take place. The cell contents are then transferred to a suitable tray, evaporated to dryness, and the radioactive content measured in standard counting equipment. A reference tray is similarly prepared, a known volume of the active solution which has not undergone the diffusion process being used. The activity of this tray is measured immediately after that of the experimental tray, and from the values so obtained, that fraction of the original cell content which has passed out of it as a result of diffusion can be calculated. Knowing the length of the capillary cell and time of diffusion, the self diffusion coefficient of the ion may be calculated [30]. This method has been extensively used for determining self-and tracer diffusion coefficients of electrolytes [31,32,33]. Other studies on theoretical, concentration gradients effect, stirring rate effect, and temperatures effect have been done on this technique [34,35,36,37,38]. This method is shown to be reliable for ESR study, since it reproduces the values of diffusion coefficients that have been reported by other technique [3,9].



To use the capillary diffusion method with ESR spectrometer, a diffusion cell that can be used under solvent vapor pressure is shown in Fig. 2.3. The capillary diffusion cell was made of 1 mm i.d. heavy wall Pyrex tubing whose open end was ground flat. A small distortion of the capillary is unavoidable in joining the capillary to the 11 mm o.d. tubing (A). The total length of the capillary,  $l_t$ , was measured with a micrometer. The arm (D) was used to introduce a solid sample to the reservoir (C). Solvent was vacuum distilled into reservoirs (C) and (B) through the O-ring joint (E). The sample preparation and the experimental procedure will be described in the next chapter.

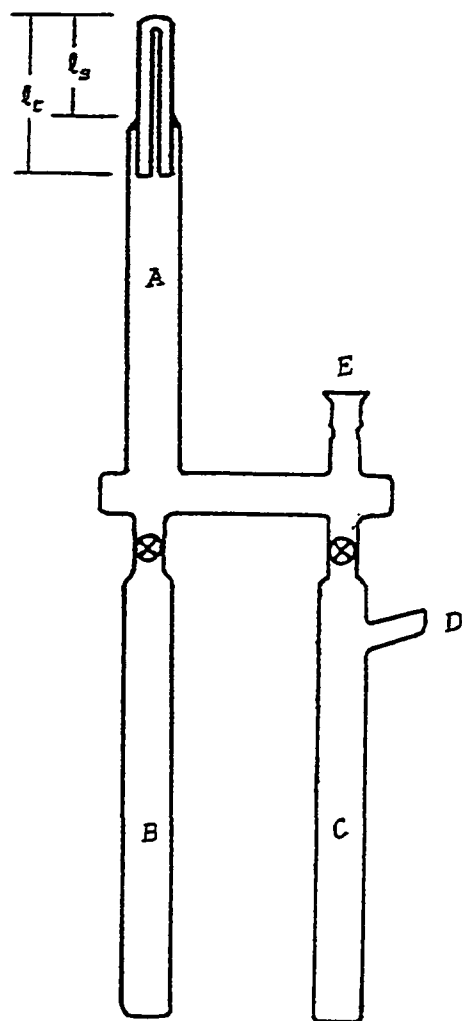


Fig. 2.3 A Capillary Diffusion Apparatus

# **CHAPTER THREE**

## **PROCEDURE**

## CHAPTER 3

### PROCEDURE

#### 3.1 SAMPLE PREPARATION

##### 3.1.1 VO[5-MeO-Sal SB] (Phen)

This compound was obtained through private communications [39]. The synthesis was carried out following the procedure outlined in reference [40]. The solvent used without further purification is methylene chloride,  $\text{CH}_2\text{Cl}_2$ , (spectroscopic grade purchased from Fluka-AG). A weighted amount is introduced to the cell, and under vacuum the solvent is added the volume of the solvent measure from the reservoir (C) in the cell which is a 10 ml graduate cylinder, the concentration is about  $(9.7 \times 10^{-3} \text{ M})$ .

##### 3.1.2 $\text{Cp}_2\text{VCl}_2$

This compound is available commercially ( PFALTZ & BAUER, INC), the solvent used is the chloroform,  $\text{CHCl}_3$ , (spectroscopic grade purchased from ARISTAR BDH). The preparation is similar to the above compound, the concentration is about  $(8.7 \times 10^{-3} \text{ M})$ .

### 3.2 Diffusion Cell Preparation

The following steps done to prepare the cell which contains the sample to be used in the ESR spectrometer:

1. Cleaning and drying the cell carefully.
2. Introduction of the weighted sample through the arm (D), then seal (D).
3. Check that all ports stop cock are close in the vacuum line, connect the cell from the O-ring joint to the vacuum line in one port.
4. Prepare the solvent needed (10-15 Ml) in a clean container and connect it to another port in the vacuum line.
5. Switch on the vacuum pump and open the main stop cock of the vacuum line, to let the pressure goes down.
6. Freeze the sample in reservoir (C) and the solvent container by two liquid nitrogen container.
7. Open the stop cock port at the two ports used, and also the stop cock in the cell at (B) and (C).
8. Degassing the cell and the solvent container until a good, low, stable pressure is obtained.
9. Close the main vacuum line stop cock, and the stop cock at (B) in the cell.
10. Remove the liquid nitrogen from the solvent, to start the evaporation of the solvent from the solvent container to condense at the freeze (C) reservoir which contains the sample.

11. One can repeat steps 6,7,8,9,10 to get more solvent in (C) for certain concentration.
12. Close the stop cock at (C) and open the one on (B), then repeat steps 6,7,8,9,10 for (B) , until good amount is collected from the solvent.
13. Close the stop cock at the solvent port, freeze (B) and (c), then open the stop cock at (B), (C), and the main vacuum line stop cock to remove the vapor of the solvent in the system.
14. Seal below the O-ring joint carefully, then close the two stop cock at (B) and (C) .
15. Close the main stop cock of, switch off the vacuum pump, and release the vacuum system slowly from any unused port to equalize the pressure.
16. Leave the cell to come to room temperature, then check the cell slowly to dissolve the sample with the solvent in (C).
17. Handle the cell pointing the capillary downward, then open slowly the stop cock at (C) to transfer the solution into the capillary through (A).
18. If a bubble of the vapor block the solution to inter the capillary, remove it by rubbing the capillary with fingers or small cloth slowly and carefully.
19. The excess liquid is poured back into (c), and the cell gently shaken to disperse any paramagnetic solution coating the wall of (a).

20. Put the cell in the ESR cavity at the close end to set up the parameters for the ESR spectrometer, use collars available to set the length of the capillary in the cavity and fix it to the cell by tape.
21. Maximum lines height should be obtained from the ESR parameter setting because these lines will decrease by time, so by the end of the experiment it will be readable.
22. Remove the cell from the cavity, open the stop cock at (B) slowly to transfer the solvent above the capillary cell in (A), put back the cell in the cavity, record the time, and take the first spectrum run.
23. The cell should not disturb in the cavity, three to four runs should be taken every day, with the same setting of the ESR spectrum parameter. Each run record the time and find the difference of time of this run and the time when the experiment start in seconds.
24. Analyze the spectrum of each run by finding the intensity of each line by the computer manipulation program, the first run can be taken as a reference, find the total intensity of the reference spectrum, enter these values into a data file for the computer program (in the appendix the program and the data files are listed for the two compounds) called Diff.
25. The data are collected for three to four weeks.

26. Run the program at the end will give two output files one to calculate the diffusion coefficient, and the other you can plot the total intensity as a function of time from the output file plot generated.

### 3.3 ESR Spectrometer Operation

The following is a brief description of how to operate the ESR spectrometer:

- (i). Put on the chiller water for circulation .
- (ii). Turn the console power on and bring the field set slowly to the central field value ( $\approx 3300 - 3400$  G), through the zero value.
- (iii). In the microwave bridge the STANDBY setting to TUNE position, while keeping the microwave power around 40 Db.
- (iv). Adjust the iris and turn on the frequency counter.

To record the spectra the capillary must be at the active region of the cavity, this could be adjusted with the collars provided. Proper tuning procedure should be followed on the microwave bridge.

Keep the field set at 3300 G, scan range of 1000 G for the room temperature spectra. For good line shape analysis, the spectra taken must be free from artifact broadening due to the wrong selection of spectrometer parameters like scan speed,



time constant, microwave power, modulation amplitude, etc. To get the right signal, first increase the modulation amplitude until no peak no peak distortion or line width increase is noticed. This value of modulation amplitude could be used for the experiments subsequently. Second, increase the microwave power until the saturation of the peak occurs. Once the right spectra are plotted, then they are stored to the tape for Varian Spectrometer, and plotted also by the plotter, the ESR parameters for the compounds are tabulated in Table 3.1.

**Table 3.1 The ESR Parameters For Experiments**

Parameters	Compound I	Compound II
Number of Point	2048	2048
Number of Scans	4	1
Field Set (G)	3440.73	3410.00
Scan Range (G)	1000	800
Scan Time (sec)	120	180
Modulation Frequency (kHz)	100	100
Modulation Amplitude(G)	2	2
Receiver gain	$1.25 \times 10^4$	$4.0 \times 10^4$
DC Gain	Normal	Normal
Time Constant	0.25 s	0.25 s
Microwave Power	2 mW	20 mW
Microwave Frequency (GHz)	9.489787	9.51050

Compound I : VO[5-MeO-Sal SB] (Phen)

Compound II :  $\text{Cp}_2\text{VCl}_2$

## **CHAPTER FOUR**

### **ESR MAGNETIC PARAMETERS**

## CHAPTER #4

### ESR MAGNETIC PARAMETERS

#### 4.1 Introduction

The analysis of the spectra to extract the magnetic parameters of the complex is obtained from the rigid limit and room temperature spectra according to Wilson and Kivelson's method [41]. The parameters are necessary for a proper interpretation of the observed ESR spectra at various temperatures in terms of correlation time  $\tau_1$  and anisotropic rotational diffusion ( $N$ ). The ESR spectra of  $\text{Cp}_2\text{VCl}_2$  in chloroform ( $\text{CHCl}_3$ ) at various temperatures, in the motional narrowing region, is analyzed to study the anisotropy of molecular rotation ( $N$ ) defined as the ratio  $R_1/R_\perp$  where  $R_\perp$  is the rotational diffusion constant along the molecular principal axis of rotation, and  $R_1$  is the rotational diffusion constant perpendicular to the molecular principal axis of rotation [42,43,44]. The results are compared with the anisotropy of rotation ( $N$ ) calculated from Stokes-Einstein model [11] and that calculated from the Allowed-Values Equation (AVE) [45]. This comparison demonstrates the applicability of these narrowing motion region equations to experimentally observable system as well as giving some valuable information about the structure. From the analysis of the temperature dependent ESR

line shapes in the motional-narrowing regime the stickiness factor ( $S$ ) which is a measure of all non-hard-core interactions is determined [15,46,47]. Also we will show how the dependence of line widths on temperature is utilized to obtain the anisotropic factor ( $N$ ) and the correlation time  $\tau_c$ . These experimental  $N$  values are compared with theoretical values [11,45].

The vanadyl radical has an electron spin of half (unpaired electron,  $d^1$ -system) which interacts with the vanadyl nucleus which has the spin  $I=7/2$  therefore the ESR spectrum consists of  $(2I+1)$ , that is eight lines.

#### 4.2 Spin Hamiltonian and Definition of the Parameters

The appropriate spin Hamiltonian,  $H_s$ , required to describe the vanadyl paramagnetic systems in general is [48,49],

$$H_s = B \beta_0 g \cdot S + S \cdot A \cdot I + (\text{electron-electron dipole term}) + (\text{electron-electron exchange term}) + (\text{nuclear Zeeman term})$$

The first term is the electron Zeeman term and the second term is the electron-nuclear hyperfine interaction term, where  $B$  is the laboratory magnetic field,  $\beta_0$  is electron Bohr magneton,  $g$  the electron  $g$ -value tensor,  $S$  is the electron spin angular

momentum operator,  $A$  is the electron-nuclear hyperfine tensor, and  $I$  is the nuclear spin angular momentum operator. The nuclear Zeeman term is not important in most work [50] and will be omitted.

The electron-electron dipole and exchange terms are significant only when very high concentration [51] and single crystals studies [52] are considered, so our Hamiltonian will be,

$$H_s = B\beta_0 g \cdot S + S \cdot A \cdot I \quad (4.1)$$

#### 4.2.1 ESR Spectra in Solution

The ESR spectra of the vanadyl complexes in liquid solution can be described by the isotropic electron Zeeman and electric-nuclear hyperfine interactions included in the spin Hamiltonian [41,53] as in eq.(4.1), the allowed transitions when microwave field at frequency  $\omega_0$  is applied, occur at the field  $B$  which satisfies the resonance condition [41,53]

$$\omega_0 = g_0 \beta_0 \frac{B}{h} + A_0 M + \frac{1}{2} A_0^2 \left[ \frac{I(I+1) - M^2}{g_0 \beta_0 (B/h)} \right] \quad (4.2)$$

where  $\omega_0$  is the microwave frequency in radians per second,  $h$  is the Planck's constant divided by  $2\pi$ ,  $I = 7/2$ ,  $M$  is the

nuclear spin magnetic moment,  $B$  is the resonant value of the magnetic field corresponding to  $M$ ,

$A_0$  is the isotropic hyperfine constant in radians per second,

$$A_0 = \frac{1}{3} (A_x + A_y + A_z) \quad (4.3)$$

and  $g_0$  is the isotropic g-factor,

$$g_0 = \frac{1}{3} (g_x + g_y + g_z) \quad (4.4)$$

where X, Y, and Z refer to the molecule fixed coordinates, and it has been assumed that  $g$  and  $A$  can be simultaneously diagonalized.

The contact hyperfine constant  $A_0$  can be obtained to second order by taking the difference of  $B_M$  for line  $M$  and  $B_{-M}$  for line  $-M$

$$A_0 = \frac{g_0 \beta_0 (B_M - B_{-M})}{2Mh} \quad (4.5)$$

furthermore,

$$g_o - g_s = g_s \left[ \frac{B_s - \frac{1}{2}(B_M + B_{-M})}{\frac{1}{2}(B_M + B_{-M})} \right] - \frac{2A_o^2 h^2 [I(I+1) - M^2]}{g_s \beta_o^2 (B_M + B_{-M})^2} \quad (4.6)$$

where  $g_s$  and  $B_s$  are the isotropic  $g$ -value and the resonant value of the magnetic field (in Gauss), for standard of known  $g$ -value. For our system, the standard used is diphenylpicrylhydrazyl (DPPH), with  $g$ -value equal to 2.0037.

The first term in eq.(4.1) the electron Zeeman term, represents the large interaction of the electron spin with the laboratory magnetic field. This term yields the relation,

$$h \nu = g_o \beta_o B \quad (4.7)$$

where  $h$  is Planck's constant and  $\nu$  is the microwave frequency of the ESR spectrometer. This equation can be reduced to [26],

$$\nu \text{ (in GHz)} = 2.80247 (g_o / g_e) B \text{ (in kG)} \quad (4.8)$$

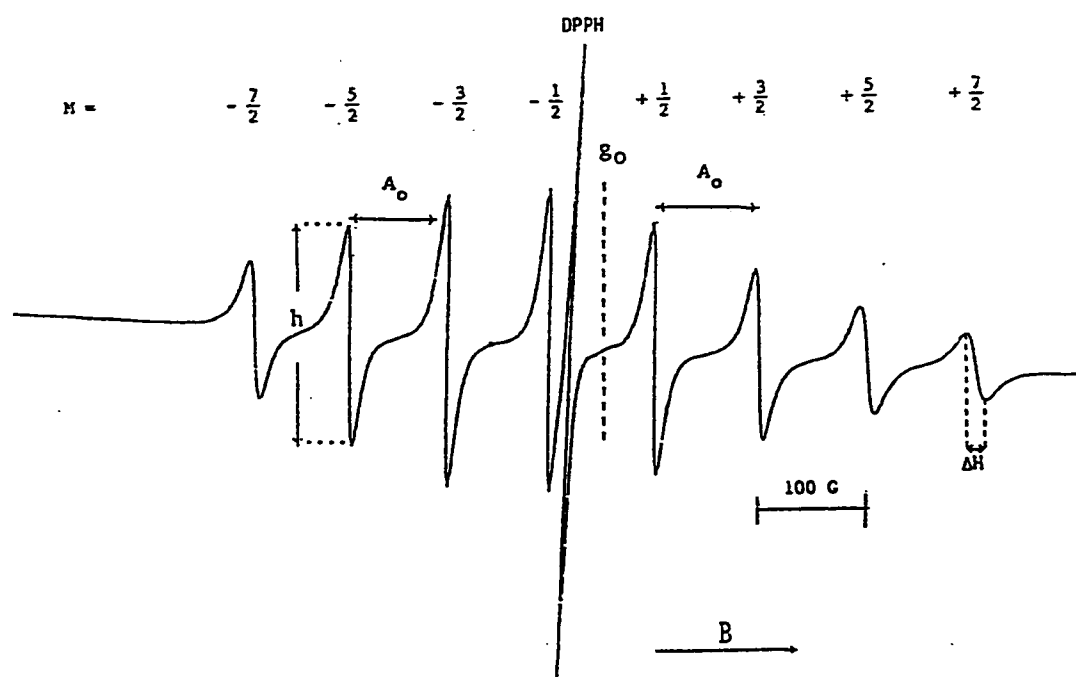


For vanadyl complexes eq.(4.8) can be further reduced to,

$$\nu \text{ (in GHz)} = 2.847 \text{ } B \text{ (in kG)} \quad (4.9)$$

This relation serves as a useful conversion factor. Most ESR spectrometers operate at  $\nu = 9.5$  GHz (X-band), and thus  $B = 3.3$  kG. The unpaired electron responsible for the ESR spectrum is confined largely to the vanadyl center, which interacts with the nuclear spin of  $^{51}\text{V}$  ( $I = 7/2$ ), the result is  $(2I + 1)$  or 8 lines separated by a coupling constant with different intensities. This nuclear interaction is represented by the second term in eq.(4.1). It is apparent from the variation in the intensities of the lines that the line widths of the different hyperfine components are not equal and depend on the  $^{51}\text{V}$  nuclear spin quantum number [41,53].

This line width variation indicates that the anisotropic hyperfine and Zeeman interaction are not completely averaged out by the tumbling of the molecules in the liquid, and they contribute significantly to the lines width. The relation of the important  $A_0$  and  $g_0$  parameters are shown in (Fig. 4.1) along with the lines widths ( $\Delta H$ ) and the peak heights ( $h$ ) in a first derivative representation of the absorption curve. The spectra can be recorded with the opposite phase from that of Fig. 4.1 producing an upside down effect. The presentation most consistent with current trends in spectroscopy involves positive phase and increasing  $B$  from left to right.



**Fig. 4.1: A Typical X-band Spectrum Of Oxovanadium (IV) in Solution at Room Temperature.**

## 4.2.2 ESR Spectra in the Absence of Molecular Motions

### 4.2.2.1 Anisotropic Effects

The glass spectrum of vanadyl probes can be described in terms of the Zeeman and nuclear-electronic hyperfine interactions by the spin Hamiltonian [41,49,54],

$$H_s = \beta_o (g_x B_x S_x + g_y B_y S_y + g_z B_z S_z) + A_x S_x I_x + A_y S_y I_y + A_z S_z I_z \quad (4.10)$$

a special case of eq.(4.1), both  $g$  and  $A$  depend on the orientation of the magnetic field in terms of molecular coordinates and the direction of  $B$ . The extremes in  $g$  and  $A$  lie along the principal X, Y, and Z axes. The glass ESR spectrum of a typical  $VO^{2+}$  at (77 k) is illustrated in Fig. 4.2 this glass spectrum of vanadyl complexes is governed by [41],

$$\omega_o = \frac{g\beta_o B_M}{h} + AM + \frac{(A_x^2 + A_y^2)(A_z^2 + A^2)}{8A^2} \frac{[I(I+1) - M^2]}{g\beta_o B_M/h} \quad (4.11)$$

where

$$g = (g_x^2 \cos^2 \theta + g_x^2 \sin^2 \theta \cos^2 \phi + g_y^2 \sin^2 \theta \sin^2 \phi)^{\frac{1}{2}}$$

and

$$gA = (A_x^2 g_x^2 \cos^2 \theta + A_x^2 g_x^2 \sin^2 \theta \cos^2 \phi + A_y^2 g_y^2 \sin^2 \theta \sin^2 \phi)^{\frac{1}{2}}$$

$\theta$  is the angle between  $B$  and unique molecular axes, and  $\phi$  is the azimuthal angle. Note that the line for which  $\theta = 0$  gives rise to the spectrum corresponding to  $g_x$  and  $A_x$  in (Fig. 4.2) while those for which  $\theta = \phi/2$ ,  $\phi = 0$  and  $\pi/2$  give rise to the "doubled spectra" corresponding to  $(g_x, A_x)$  and  $(g_y, A_y)$  as in Fig. 4.2. The anisotropic tensors can be determined from ridged limit spectra, Fig. 4.2, as Wilson and Kivelson have done, for vanadyl acetylacetonate VOAA [41]. The spacings I, II, III, IV, and V in Fig. 4.2 correspond to,

$$A_x = \frac{II + III}{2} \quad (4.12)$$

$$A_x + A_y = IV + V \quad (4.13)$$

$$I = \frac{\beta_o B^2 \delta g}{h} + 4C\left(\frac{7}{2}\right) \quad (4.14)$$

where

$$\delta_g = \frac{1}{2} (g_x - g_y) \quad (4.15)$$

$$c = \frac{1}{4} (A_x - A_y) \quad (4.16)$$

although the values of  $g_x$ ,  $g_y$ ,  $g_z$  and  $A_x$ ,  $A_y$ ,  $A_z$  obtained from the glassy spectrum are compatible with the isotropic values obtained from the liquid spectrum, the most accurate determinations of the magnetic parameters are obtained by measuring  $g_0$  and  $A_0$  from the liquid spectrum,  $g_z$ ,  $g_x - g_y$ ,  $A_z$  and  $A_x - A_y$  from the glass spectrum, and  $g_x$ ,  $g_y$ , and  $A_x$ ,  $A_y$  from this data and eqs. (4.14 and 4.15) [41].

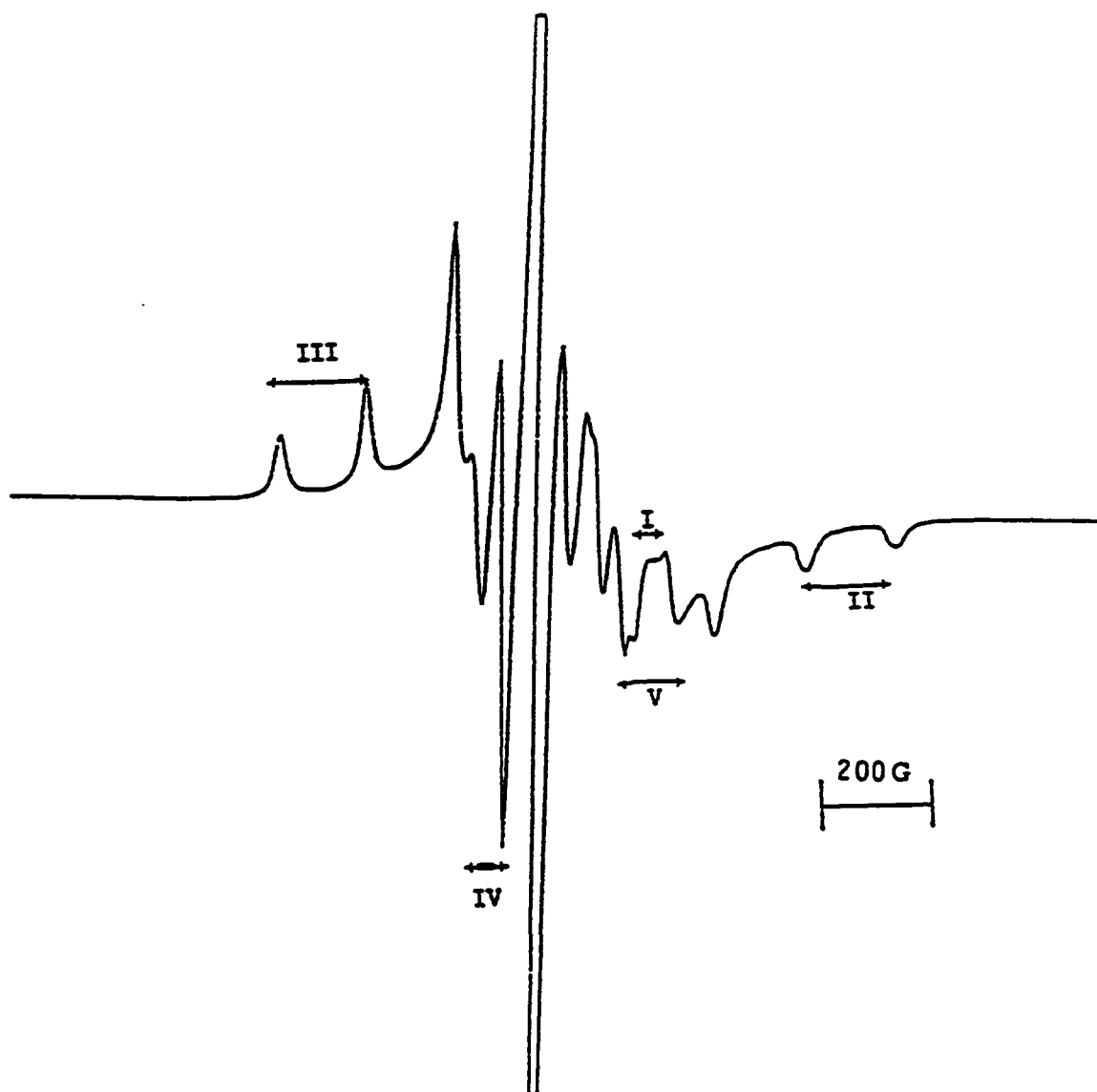


Fig. 4.2 Ridgid Limit spectra at 77 k

#### 4.2.2.2 The Effects of Rotational Motion on ESR Spectra

The shape of an EPR spectrum is markedly dependent upon the rotational mobility of the paramagnetic spin probe. For example, in a nonviscous solvent nearly all vanadyl probes exhibit eight sharp lines. The rapid isotropic tumbling motion averages out all anisotropic effects discussed above, therefore sharp lines are obtained. If the rotational motion is slowed down by increasing the solvent viscosity, averaging is incomplete. The result is unequal broadening of the eight lines. Molecular motion is not necessarily isotropic. Preferential motion about one axis is likely to occur for any asymmetric spin probe, but the effects are particularly evident when surrounding medium is anisotropic. The ESR spectrum is very sensitive to the probe mobility. The best way to make use of this sensitivity is to describe the motion quantitatively in terms of the shape of the spectrum.

As discussed above, both the anisotropic hyperfine electron-nuclear interaction and the anisotropic spin-orbit interaction in a spin probe depend on the mutual orientation of the external magnetic field and of  $\pi$ -orbital of the unpaired electron. Rotations of the probe modulate these interactions, leading to fluctuations in local magnetic fields and changes of the ESR line widths. Rotational motion is characterized by correlation time ( $\tau_R$ ) which determines the effective frequency of rotation,  $f_{\text{eff}}$  (in Hz)  $= (2\pi\tau_R)^{-1}$ .

The principal  $g$  and  $A$  values are related to the isotropic  $g_0$  and  $A_0$  values through eqs (4.3 and 4.4).

The anisotropy of  $g$  and  $A$  tensors is applied in dynamic spin probe studies to get information about the motional state of the probe. Between the two limits of very rapid motion and a rigid glass, the spectra are quite complex. This complexity provides much of the information that can be obtained about molecular motion using spin probe techniques.

### 4.3 The Correlation Time Group

The correlation time data of vanadyl complexes usually divided into three groups,

Rapid rotations	$10^{-11}$	-	$10^{-9}$	sec.
Slow rotations	$10^{-9}$	-	$10^{-6}$	sec.
Very Slow rotations	$10^{-6}$	-	$10^{-3}$	sec.

The limits of these regions are essentially determined by the anisotropy of the magnetic interactions occurring in the probes.



#### 4.4 Molecular Tumbling

The spin Hamiltonian can be divided into,

$$H = H_0 + H_1(\Omega) + e(t) \quad (4.17)$$

where

$$H_0 = \beta_0 g_0 B S_z + A_0 I \cdot S \quad (4.18)$$

where  $S_z$  is the electron spin along the Z-axis,  $S$  and  $I$  are the electron and the nuclear spin operator,  $B$  is the laboratory magnetic field, and  $\beta_0$  is electron Bohr magneton, and  $H_1$  is equal to,

$$H_1 = \beta_0 B \cdot g S + I \cdot A \cdot S \quad (4.19)$$

where  $g$  and  $A$  contain the an isotropic components of the electron  $g$ -tensor and electron-nuclear dipole hyperfine tensor after the isotropic components  $g_0$  and  $A_0$  have been subtracted. These tensors are defined in terms of molecular fixed axis whose orientation varies with time, with respect to the laboratory magnetic field due to molecular motion.

Thus the time dependence of  $H_1$  is modulated by the reorientational motion of the molecule ( $\Omega$ ), and

$$H_1(t) = \frac{1}{2} h \gamma_e B [S_+ \exp(-i\omega t) + S_- \exp(i\omega t)] \quad (4.20)$$

where

$$S_+ = S_x + iS_y$$

is the raising operator,

$$S_- = S_x - iS_y$$

is the lowering operator,  $\gamma_e$  is the magnetogyric ratio of a free electron and  $\omega$  is the microwave frequency.  $H_1(t)$  gives the interaction of the spin with the microwave radiation field which must be included in the Hamiltonian if the saturation effects are to be considered. Saturation effects are known to be negligible for experimental spectra discussed, only terms linear in  $H_1(t)$  must be considered [54]. The term "fast" and "slow" refers to the randomly fluctuating part  $H_1(t)$  spin Hamiltonian (where  $H_1(t)$  has a time averaged value zero) in relation to a correlation time  $\tau_c$  which characterizes the underlying stochastic process.

Fast motional theories are applicable if  $|H_1(t)|\tau_R \ll 1$  where  $H_1(T)$  is expressed in frequency units. On the other hand, the slow tumbling is attained when  $|H_1(t)|\tau_R \gg 1$ . Since the slow and the very slow rotations lie outside the scope of this work, only the fast tumbling rotations will be discussed.

#### 4.5 Fast Tumbling

It is convenient at this point to define a line width parameter  $T_2$ . If the peak to peak width is  $\Delta H$  (in gauss), then  $1/T_2 = \Delta H$ . In general, the dependence of  $T_2$  upon  $M$ , the component of the nuclear spin along the direction of the applied magnetic field is given by,

$$T_2^{-1}(M) = A + BM + CM^2 + DM^3 \quad (4.21)$$

where  $A$ ,  $B$ ,  $C$ , and  $D$  depend on the magnetic tensors. The line width contributions can be calculated with the Redfield relaxation matrix theory as applied by Freed and Fraenkel [55,56] or with the theory of Kubo and Tomita as applied by Kivelson [57].

In order to determine the anisotropy of rotation ( $N$ ) and the axis of rotation, the motional narrowing analysis of Freed *et al.* [42,43,44] has been used and eq.(5) in reference [44] has been used.

In practice  $\tau_R$  is obtained from  $B$ ,  $C$ , and  $D$  coefficients because of the uncertainty in  $A$  [51]. The theoretical expressions for  $B$  and  $C$  coefficients will be discussed in the next chapter [42,43].

For isotropic Brownian  $\tau_R = (6R)^{-1}$  where  $R$  is the rotational diffusion coefficient. For the eight lines spectrum for vanadyl probe, eq.(4.21) will give rise to four experimentally determined unknowns,  $A$ ,  $B$ ,  $C$ , and  $D$ . Since the line width coefficients  $B$  and  $C$  are function of  $\tau_R$  values, in the case of two values of  $\tau_R$  value not agreeing, the assumption of anisotropic rotation must be suspected. Fortunately, this restriction has been resolved by Freed and his co-worker [43,58]. The cubic dependence of  $(1/T_2)$  on  $M$  given in eq.(4.21) is retained, but the motion is described by a rotational tensor  $R$ . This is often axially symmetric and is defined by two components  $R_{\parallel}$  and  $R_{\perp}$ , where  $R_{\parallel}$  is the rotational diffusion constant along the principal axis of rotation, and  $R_{\perp}$  is the rotation diffusion constant perpendicular to the principal axis of rotation. The anisotropy of rotation ( $N$ ) is defined as the ratio of  $R_{\parallel}/R_{\perp}$ .

Two related correlation times  $\tau(0)$  and  $\tau(2)$  are given by,

$$\tau(0) = (6R_{\perp})^{-1} \quad (4.22)$$

$$\tau(2) = (2R_{\perp} + 4R_{\parallel})^{-1} \quad (4.23)$$

Reference should be made to the original literature [43,58] and more recent work [59] for details of this analysis.

## **CHAPTER FIVE**

### **RESULTS AND DISCUSSION**

## CHAPTER # 5

### RESULTS AND DISCUSSION

#### 5.1 DETERMINATION OF THE MAGNETIC PARAMETERS

##### 5.1.1 Isotropic Magnetic Parameters

The spectrum of  $\text{Cp}_2\text{VCl}_2$  is shown in (Fig.5.1). The isotropic hyperfine constants,  $A_0$ , and  $g_0$  were obtained from  $B_H$  for line  $M$  and  $B_{-H}$  for line  $-M$  by eq.(4.5) and eq.(4.6).  $A_0$  and  $B_0$  were determined for each pair of  $M$  and  $-M$  lines and average over all pairs.

##### 5.1.2 Anisotropic Magnetic Parameters

The rigid limit of  $\text{Cp}_2\text{VCl}_2$  is shown in (Fig.5.2). The anisotropic tensors were determined from this rigid limit spectrum using second - order perturbation theory , and from the isotropic  $A_0$  and  $g_0$  determined from the liquid spectra, as Wilson and Kivelson have done for VOAA [41]. Spacings II and III , from the rigid limit spectrum are used in eq.(4.12) to determine  $A_i$

$$A_i \text{ (G)} = ( \text{ II} + \text{III} ) / 2$$

as the resonant lines for the magnetic field parallel to Z-axis ( $H \parallel Z$ ) are concentric around  $g_I$  value . From the two terminal resonant lines in the glass spectrum, (Fig.5.2), which corresponds to  $H \parallel Z$  ( $B_{1/2}$  and  $B_{7/2}$ ) value of  $g_I$  can be obtained as follows,

$$g_z = \frac{h\nu}{\beta_o} / \left[ \frac{B_{7/2} - B_{1/2}}{2} \right] \quad (5.1)$$

where  $h$  is Planck's constant,  $\nu$  is the microwave frequency in Hz,  $\beta_o$  is the Bohr magneton. combining the values of  $A_I$  and  $g_I$  obtained from rigid limit spectrum with those of  $A_o$  and  $g_o$  obtained from the liquid spectra in eq.(4.3) and eq.(4.4), values of  $A_x + A_y$  and also  $g_x + g_y$  were obtained as follows,

$$A_x + A_y = 3A_o - A_z \quad (5.2)$$

$$g_x + g_y = 3g_o - g_z \quad (5.3)$$

the rigid limit resonant lines for  $H \parallel Z$  are governed by,

$$\omega_o = \frac{g_z \beta_o B_M}{h} + A_z M + \frac{(A_x^2 + A_y^2) \left( \frac{63}{4} - M^2 \right)}{4 g_z \beta_o B_M / h} \quad (5.4)$$

which is the case in eq.(4.11) where  $\theta = 0$  . Substituting  $A_y$  with respect to  $A_x$  from eq.(5.2) in eq.(5.4) will yields  $A_x$  value which can be used in eq.(5.2) again to get  $A_y$  . Using eq.(5.4) and resonant lines for  $H \parallel Z$  , the tensors obtained above were further adjusted. Spacing I is

$$I \text{ (Hz)} = \left[ \frac{\beta_o B (g_x - g_y)}{h} \right] + \frac{7}{2} (A_x - A_y) \quad (5.5)$$

where  $B$  (in gauss) is calculated from the microwave frequency (in Hz) and  $g_o$  by eq.(4.7), spacing I measured from the rigid limit (Fig.5.2) in gauss, should be used in eq.(5.5) in Hz. substituting  $A_x$  and  $A_y$  values obtained from eq.(4.5) and values of  $g_y$  from eq.(5.3) the following expression obtained,

$$I \text{ (Hz)} = \left[ \frac{\beta_o B (2g_x + g_x - 3g_o)}{h} \right] + \frac{7}{2} (A_x - A_y) \quad (5.6)$$

which is solved to give the  $g_x$  value and then  $g_y$  is also obtained using eq.(5.3).



Again when  $\theta = \pi/2$  and  $\phi = 0$  and  $\pi/2$  in eq.(4.11) yield the following two equations which govern the resonant lines for  $H \parallel X$ -axis and  $H \parallel Y$ -axis respectively,

$$\omega_o = \frac{g_x \beta_o B_H}{h} + A_x M + \frac{(A_x^2 + A_y^2) (A_x^2 + A_z^2) (\frac{63}{4} - M^2)}{8 A_x^2 g_x \beta_o B_H / h} \quad (5.7)$$

$$\omega_o = \frac{g_y \beta_o B_H}{h} + A_y M + \frac{(A_x^2 + A_y^2) (A_y^2 + A_z^2) (\frac{63}{4} - M^2)}{8 A_y^2 g_y \beta_o B_H / h} \quad (5.8)$$

The above two equations are used to verify the calculated values of the  $g$  and  $A$  tensors.

Several magnetic parameter values for  $Cp_2VCl_2$  have been reported [62,65,66,67], these values are tabulated in table 5.1, the experimental result obtained from temperature variation and the line width measurements [13], will be analyzed by using the reported values to find the best reported value, from which we determine the anisotropy of rotation  $N$ .

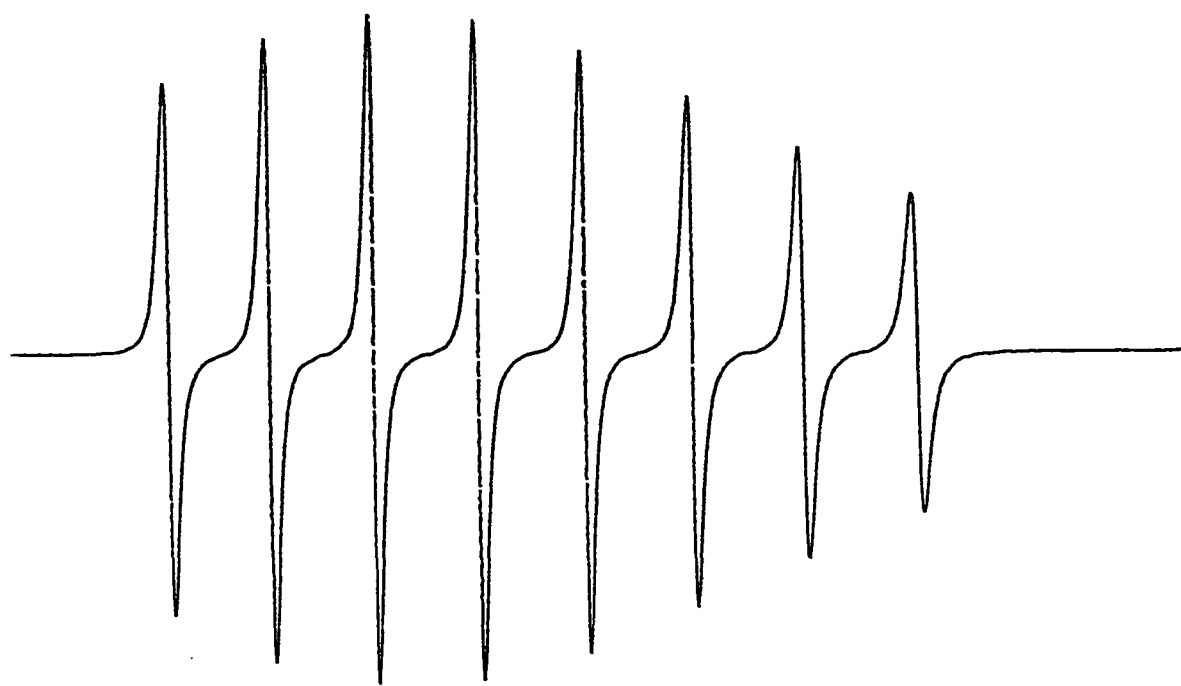
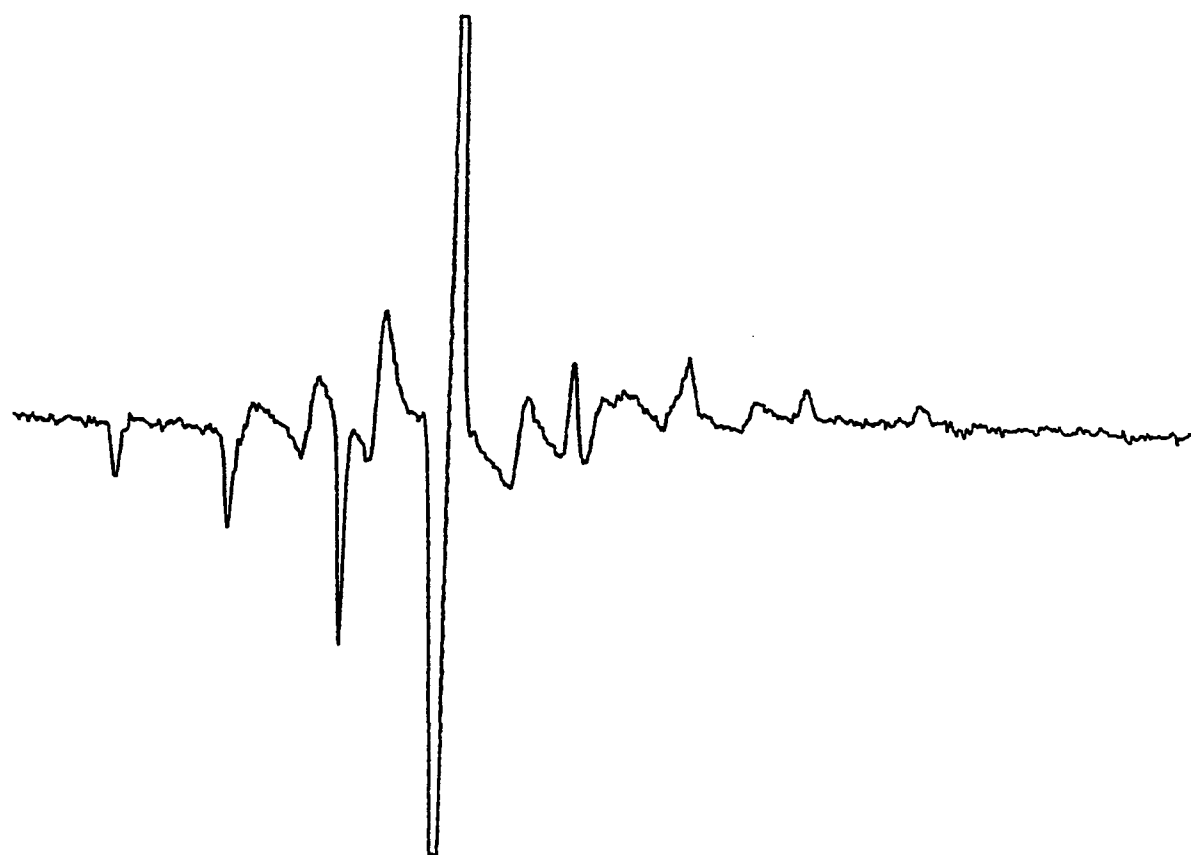


Fig. 5.1: X-band Spectrum of  $\text{Cp}_2\text{VCl}_2$  at room temperature.



**Fig 5.2: X-band Spectrum of  $\text{Cp}_2\text{VCl}_2$  at  $T = 77 \text{ K}$ .**

Table 5.1: Reported Magnetic Parameters for  $\text{Cp}_2\text{VCl}_2$

Mag.Par in (G)	(a)	(b)	(c)	(d)
$g_x$	1.986	1.965	1.986	1.9802
$g_y$	1.971	1.946	1.971	1.9695
$g_z$	2.000	1.967	2.010	2.0013
$g_0$	1.98567	1.95933	1.98900	1.98367
$A_x$	79.808	81.974	76.444	80.600
$A_y$	126.183	130.065	123.818	125.500
$A_z$	16.177	29.511	20.457	20.600
$A_0$	74.05611	79.44862	75.13467	75.56670

(a) C. P. Stewart and A. L. Porte [66],  $\text{Cp}_2\text{VCl}_2$  in  $\text{CH}_2\text{Cl}_2$  + toluene.

(b) D. P. Bakalik and R. G. Hayes [67],  $\text{Cp}_2\text{VCl}_2$  in  $\text{CH}_2\text{Cl}_2$  + toluene.

(c) A. T. Casey and J. B. Raynor [65],  $\text{Cp}_2\text{VCl}_2$  in toluene.

(d) J. L. Petersen and L. F. Dahl [62], Single-Crystal EPR Study of  $\text{V}(\eta^5\text{-C}_5\text{H}_4\text{CH}_3)_2\text{Cl}_2$  Doped in  $\text{Ti}(\eta^5\text{-C}_5\text{H}_4\text{CH}_3)_2\text{Cl}_2$ .

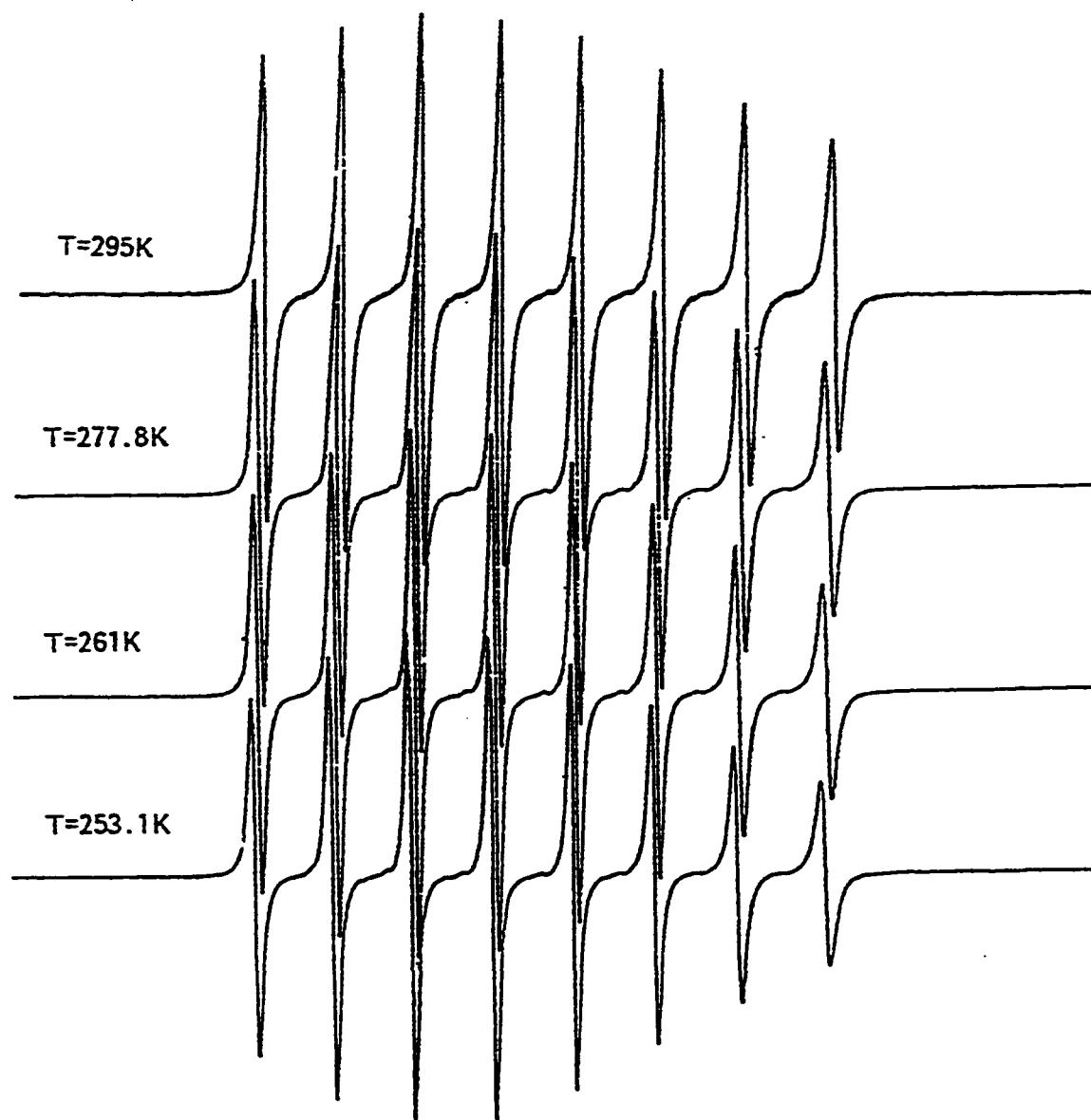
## 5.2 MOTIONAL NARROWING ANALYSIS

The spectra in Fig. 5.3 are those of  $\text{Cp}_2\text{VCl}_2$  in toluene taken over a temperature range (294 - 178 K). This free radical has an electron spin (1/2) which interacts with the vanadyl nucleus which has spin  $I = 7/2$ , therefore the ESR spectrum consists of eight lines each corresponding to a nuclear spin magnetic quantum number  $M$ :

$$M = I, I-1, \dots, -I$$

The line width parameter,  $T_2$  (in gauss<sup>-1</sup>), of each of these lorentzian hyperfine lines can be described by an expansion in  $M$ , the component of the nuclear spin along the direction of the applied magnetic field given by eq.(4.21).

The line widths and the peak to peak heights are obtained experimentally utilizing the computer graphic display which facilitates on screen line width measurements [13]. After measuring the line widths of each of these line, the widths are then entered into the computer program written for eq(4.21) for polynomial regression (called POLREG in Watfive computer language) where the coefficients of the best polynomial fit are output for the experimental values by using least square fitting, since our fitting is cubic four coefficients will be produce  $A(0)=A$ ,  $A(1)=B$ ,  $A(2)=C$ , and  $A(3)=D$ . The B and C values for different temperatures are shown in (Table 5.2) with the lines widths, also (Fig.5.4) shows the B Vs C plot of the experimental values.



**Fig.5.3 : Experimental spectra of  $\text{Cp}_2\text{VCl}$  at Different Temperature**

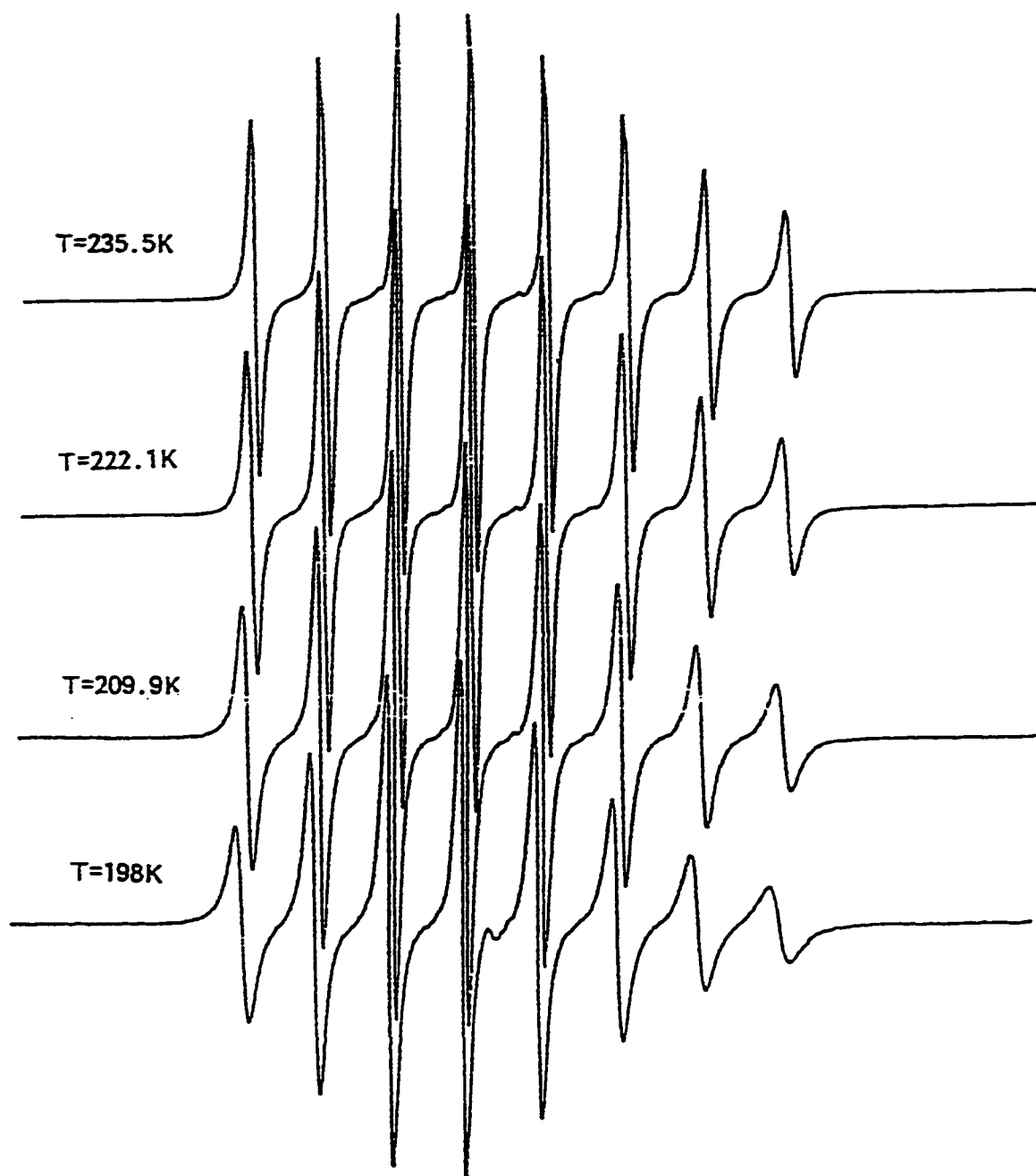


Fig.5.3(cont.) : Experimental Spectra of  $\text{Cp}_2\text{VCl}_2$  at Different Temperature

Table 5.2 : Experimental Values of Temperature & Lines Widths

Measurement & B & C values

T.°C $\Delta H(M)$	21.7	4.6	-12.2	-20.1	-37.7	-51.1	-63.3	-75.2
$\Delta H-7/2$	6.523	6.475	6.679	6.902	7.598	8.771	10.647	13.757
$\Delta H-5/2$	6.147	6.033	6.081	6.208	6.540	7.216	8.396	10.434
$\Delta H-3/2$	5.982	5.846	5.824	5.888	6.026	6.463	7.223	8.736
$\Delta H-1/2$	6.026	5.900	5.880	5.934	6.042	6.418	7.123	8.497
$\Delta H+1/2$	6.258	6.147	6.223	6.328	6.549	7.066	7.992	9.692
$\Delta H+3/2$	6.664	6.697	6.884	7.094	7.583	8.457	9.922	12.285
$\Delta H+5/2$	7.274	7.448	7.813	8.154	9.088	10.538	12.847	16.433
$\Delta H+7/2$	8.050	8.382	9.056	9.589	11.101	13.361	16.709	21.930
B	.2312	.2864	.3541	.4002	.5202	.6681	.9077	1.207
c	.09568	.1168	.1512	.1759	.2544	.3601	.5103	.7302
C/B	.41384	.4078	.4270	.4395	.4890	.5390	.5622	.6050



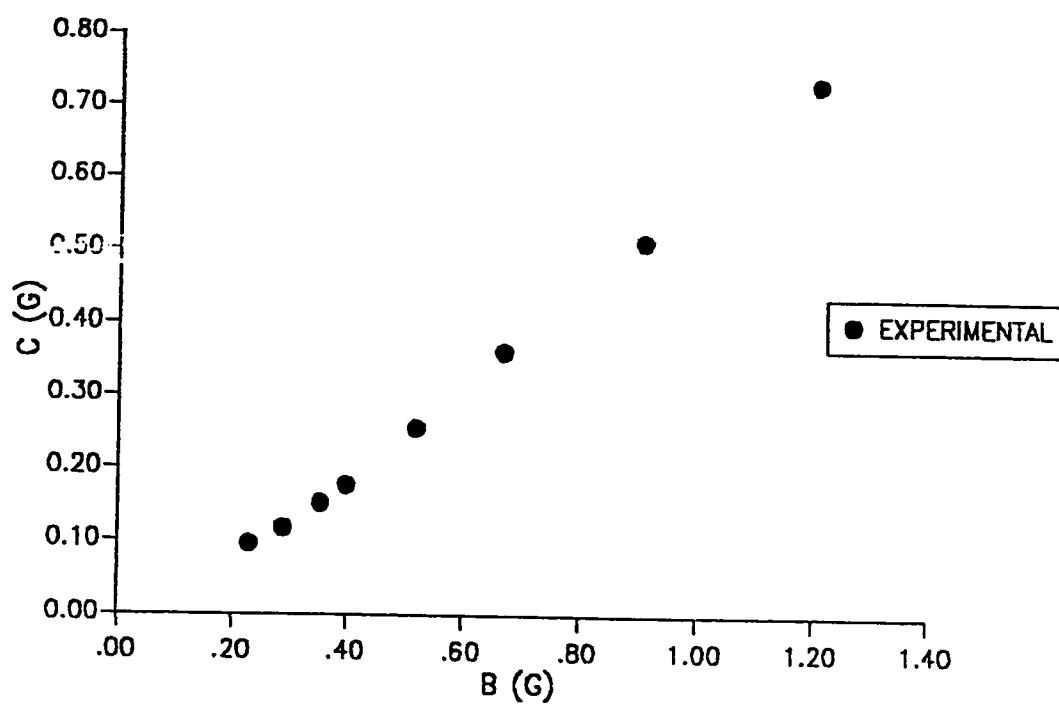


Fig. 5.4 : B Vs C Experimental Values for  $\text{Cp}_2\text{VCl}_2$  at Different Temperatures.

In order to determine the anisotropy of rotation (N) and the axis of rotation, the theoretical calculations of B and C coefficients which are given by the following equations [59,60] were used. These equations were adopted from the motional-narrowing of Freed et al.[42,43,44].

$$C_o = \frac{8}{3} - [1 + (\omega_a \tau_o)^2]^{-1} - \frac{1}{\{3 [1 + (\omega_o \tau_o)^2]\}} \quad (5.9)$$

$$C_2 = \frac{8}{3} - [1 + (\omega_a \tau_2)^2]^{-1} - \frac{1}{\{3 [1 + (\omega_o \tau_2)^2]\}} \quad (5.10)$$

$$C = \frac{2}{\sqrt{30}} (0.8\pi)^2 (D_o \tau_o C_o + 2D_2^2 \tau_2 C_2) \quad (5.11)$$

$$B_o = \frac{16}{3} + \frac{4}{1 + (\omega_o \tau_o)^2} \quad (5.12)$$

$$B_2 = \frac{16}{3} + \frac{4}{1 + (\omega_2 \tau_2)^2} \quad (5.13)$$

$$B = \frac{-2}{\sqrt{30}} (0.1\pi \omega_o) (g_1 D_o \tau_o B_o + 2g_2 D_2 \tau_2 B_2) \quad (5.14)$$

where

$$\nu = g_o (8.7940147 \times 10^6) \text{ Sec}^{-1} \text{ G}^{-1}$$

$$\nu_p = \nu \frac{g_o}{g_e}$$

$$g_e = 2.00232$$

$$g_1 = (g_x - g_o) \left(\frac{3}{2}\right)^{\frac{1}{2}}$$

$$g_2 = \frac{(g_x - g_y)}{2}$$

$$\omega_o = 2 \pi \nu$$

$$\omega_a = \frac{A_o \nu_p}{2}$$

$$D_o = (A_x - A_y) \left(\frac{|\nu_p|}{2\pi}\right) \left(\frac{3}{8}\right)^{\frac{1}{2}}$$

$$D_2 = (A_x - A_y) \left(\frac{|\nu_p|}{8\pi}\right) \text{ MHz}$$

$$\tau_o = \tau_R N^{\frac{1}{2}}$$

$$\tau_2 = \frac{3 \tau_o}{(1+2N)}$$

and the magnetic parameters  $g$ 's and  $A$ 's are in table 5.1 .

The theoretical calculations of  $B$  and  $C$  coefficients are performed by varying  $\tau_p$  in eqs(5.9, 5.14). The program written for eqs(5.9-5.14) is known as the ABC program (in Basic language). While performing theoretical simulation at the beginning we mainly look at the C/B ratio with the

anisotropic rotational diffusion and reorientation. Basically, the simulation of these curves using eqs(5.9-5.14) involves the fitting of three parameters the C/B ratio, the principal axis of rotation ( Z'-axis) and N. Usually the approximate ratio of C/B gives some hint on the Z'-axis . Value of N is then varied to move the upper portion of the curve to approach the experimental points. Appendix 2 shows N,  $\tau_1$  values and the calculated B and C values for Z' = X, Y, Z, by using the reported magnetic parameters in table 5.1, at frequency of 9.511 GHz . plotting B Vs C for the different values of N for each Z'-axis with the B and C values of the experiment will show for which N value and on which Z'-axis the experimental values is fitted. In Fig. 5.5 we use the magnetic parameter set (a) in table 5.1, it is clear that the experimental points are below the calculated values even by increasing the N values for higher number. Similarly by using the values of set (b) in Fig. 5.6 the experimental points below the calculated one. In Fig. 5.7, by using set (c) values, the experimental points come above the calculated values. Fig. 5.8 by using set (d), shows that at Z'=X and Z the experimental points come above the calculated line, but at Z'=Y and for N=2.7 the experimental points are on the calculated line.

This suggest that  $\text{Cp}_2\text{VCl}_2$  has a preferential rotation axis  $Z'=Y$  at  $N=2.7$  over the other two axes. In Fig. 5.9 we compare all sets (a,b,c,d) at  $Z'=Y$  and  $N=2.7$  to show clear fitting of the experimental values with the calculated one. Also in Fig. 5.10 we show the effect of changing the frequency on the fitting which is small effect.

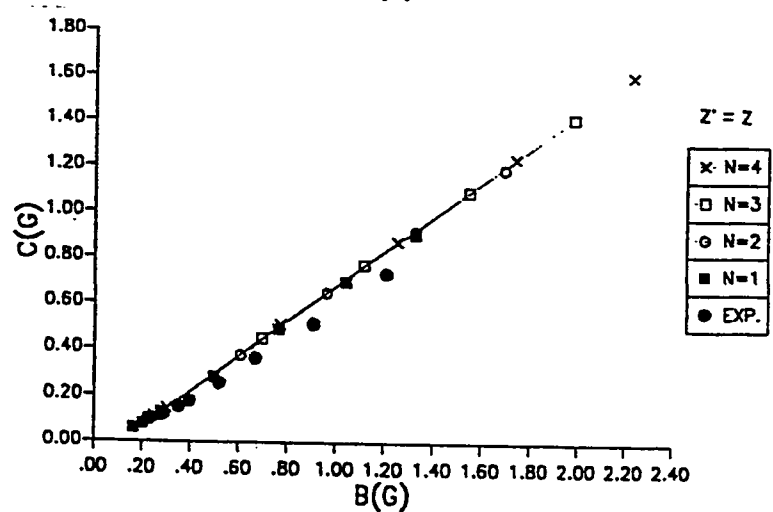
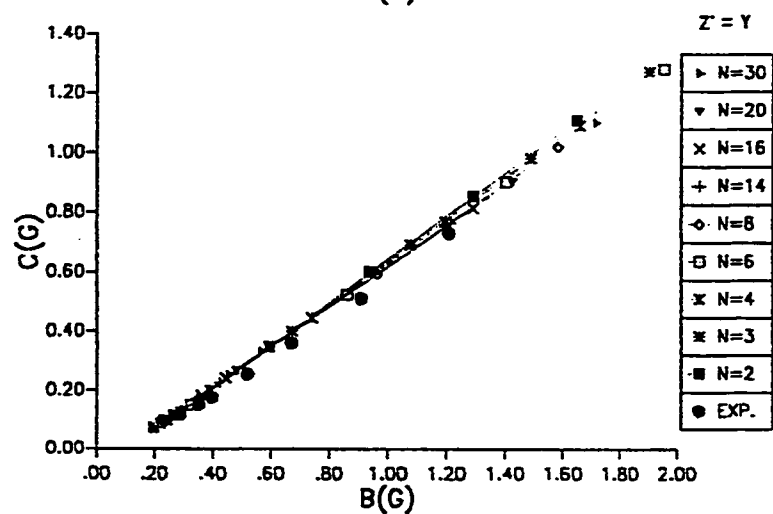
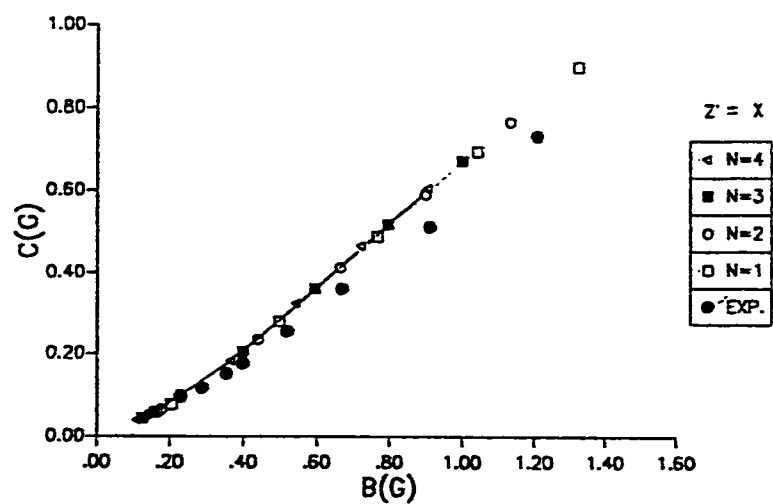


Fig. 5.5 : Experimental fit using the magnetic parameter  
set(a) for  $Z'=X,Y,Z$

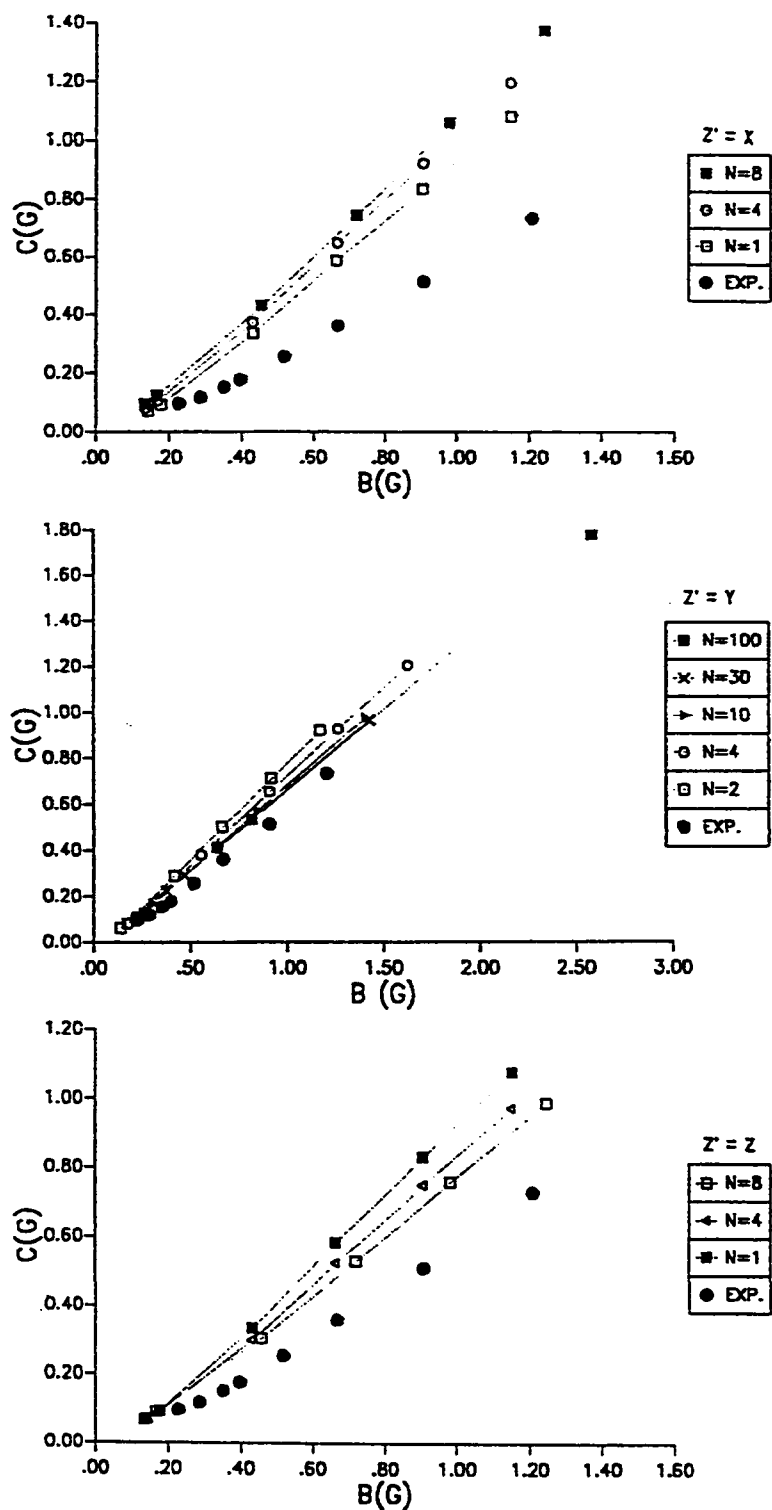


Fig. 5.6 : Experimental fit using the magnetic parameter  
set(b) for  $Z'=X,Y,Z$

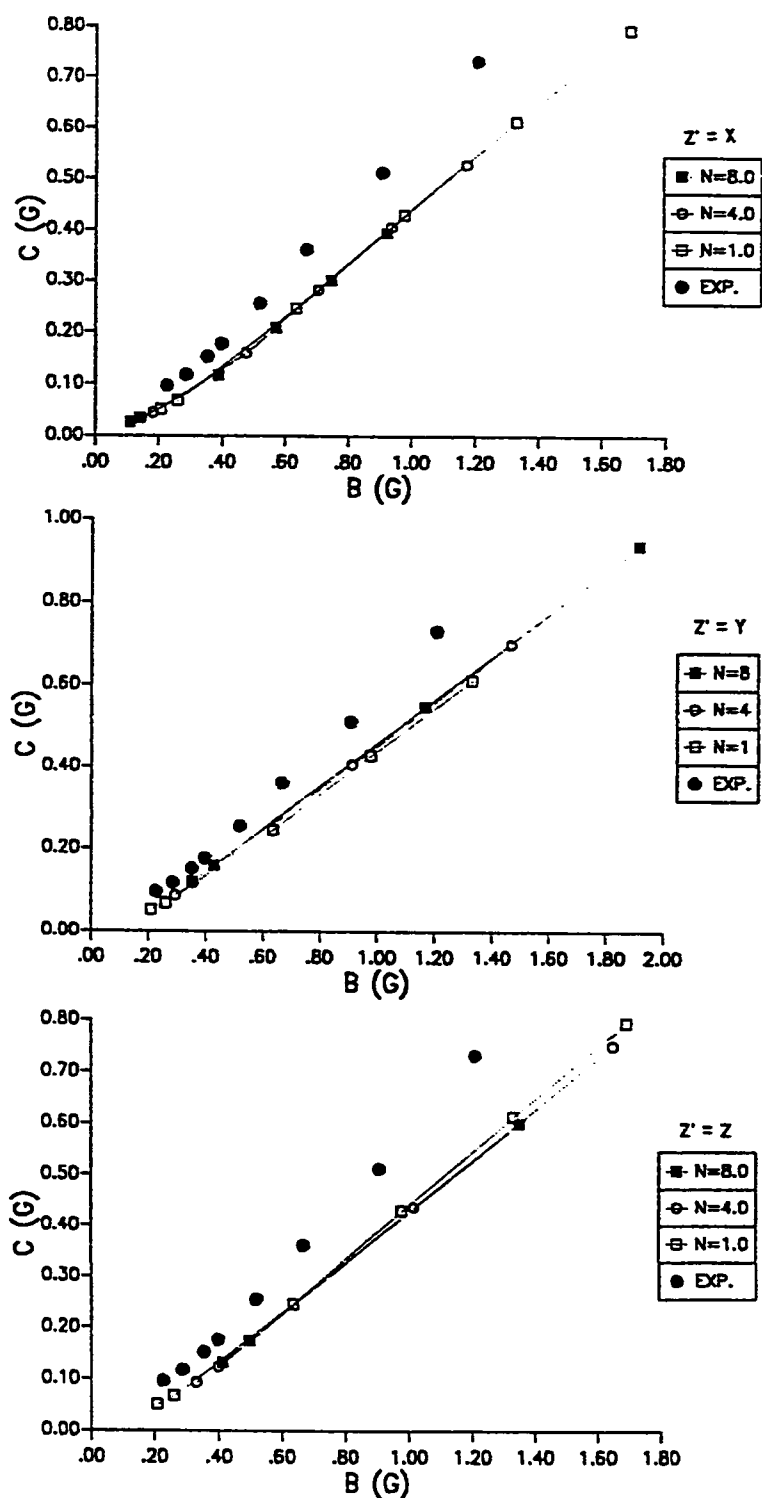


Fig. 5.7 : Experimental fit using the magnetic parameter  $\text{set}(c)$  for  $Z'=X,Y,Z$



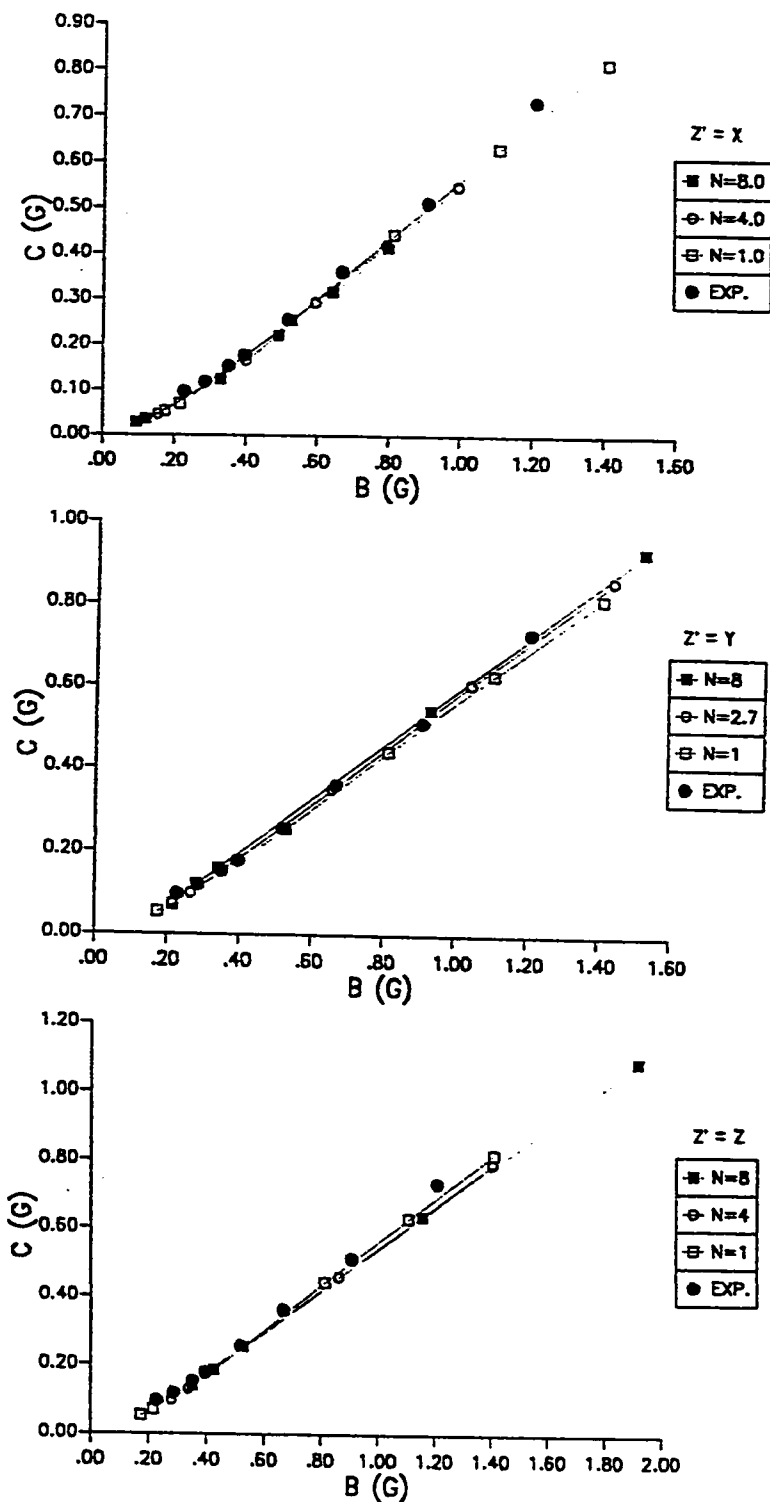


Fig. 5.8 : Experimental fit using the magnetic parameter set(d) for  $Z'=X,Y,Z$  the best fit at  $Z'=Y$  and  $N=2.7$

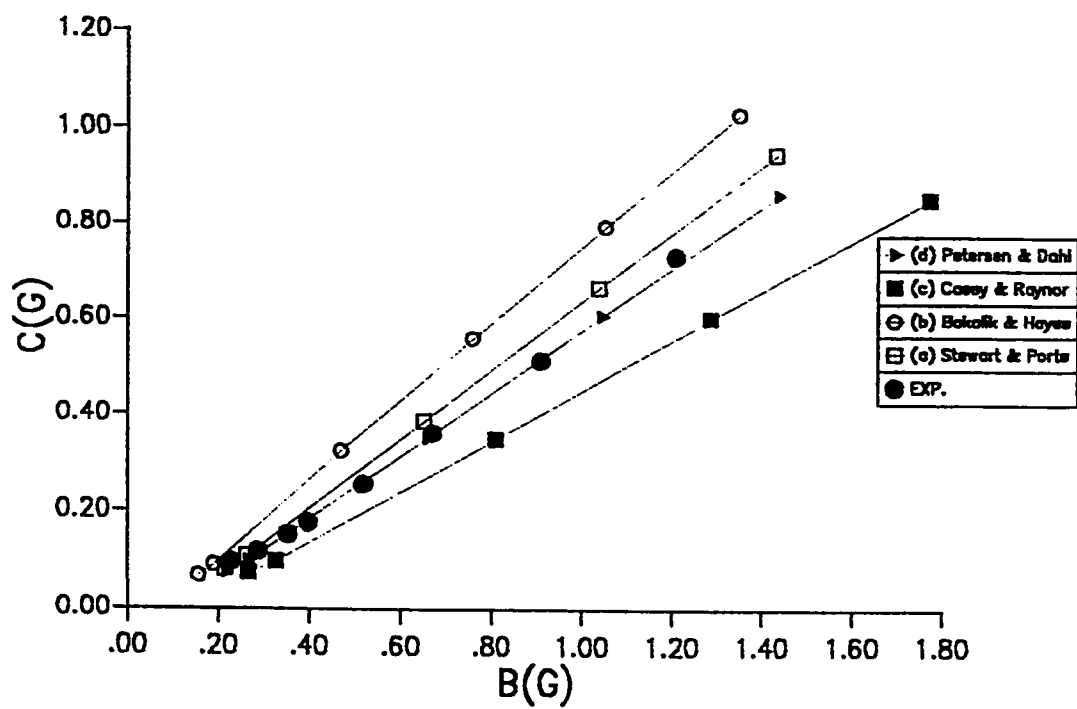


Fig. 5.9 : Experimental fit using the magnetic parameter  
sets(a,b,c,d) for  $Z'=Y$  and  $N=2.7$

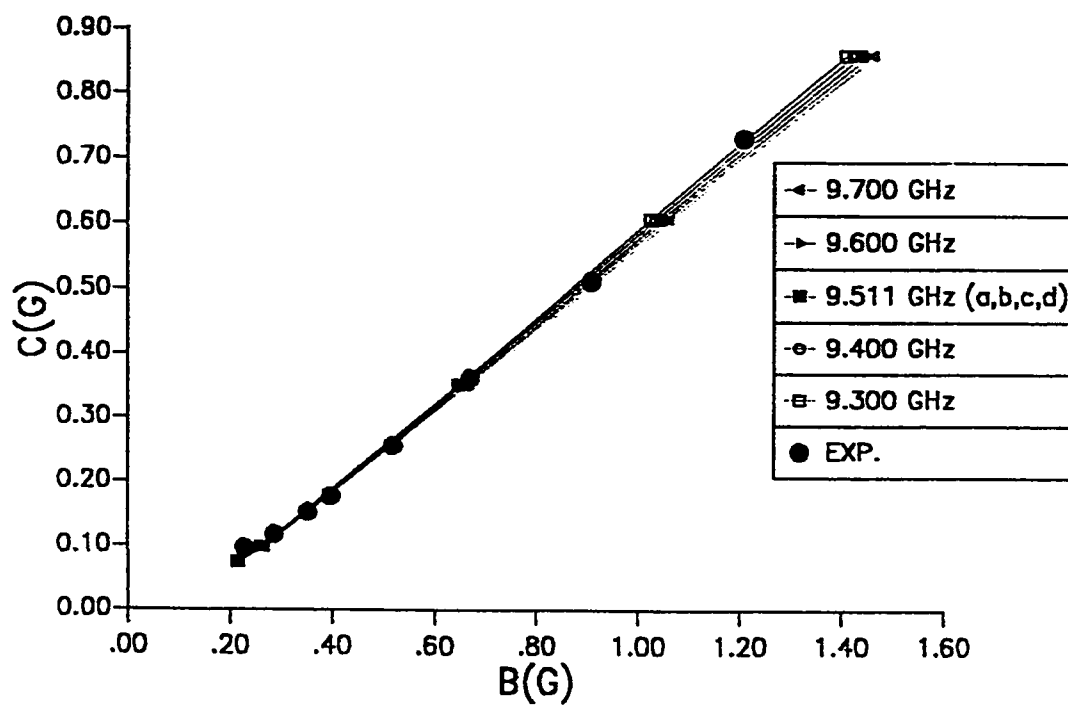


Fig. 5.10 : Experimental fit using the magnetic parameter set(d) for  $Z'=Y$  and  $N=2.7$  at different frequency.

Form the best fit plot in (Fig.5.8) one can generate an other plot between the  $C$  theoretical values Vs.  $\tau_p$  (table 5.3), then we can extrapolate the experimental values  $C$  to fined  $\tau_p$  experimental values from the straight line equation of  $C$  Vs.  $\tau_p$  of the theoretical values, as in Fig.5.11 . By knowing the experimental  $\tau_p$  at different temperature and the viscosity at different temperature we can use eq.(1.2) to find the relation (slope), if we plot the correlation time  $\tau_p$  Vs.  $\eta/T(K)$  for this experimental values as in Fig.5.12 . The viscosity values for toluene is corrected for different temperature using [61] the following equation,

$$\eta = \frac{a}{(b + t)^n} \quad (5.15)$$

where,  $a = 18.954$  ,  $b = 112.99$  ,  $n = 1.6522$ ,  $t$  the temperature in( $^{\circ}C$  ), the calculated value of  $\eta$  is in poise. Table 5.4 show the results. The slope we get is equal to

$$\text{The Slope} = \left( \frac{4}{3} \pi r_o^3 \right) \left( \frac{\kappa}{k_B} \right)$$

where the first bracket term is the volume of the molecule and the value of  $r_o$  will be determine in the next experimental result from the translational diffusion constant, from which we can recalculate  $\kappa$  (the anisotropic interaction parameter) instead of estimating its volume by computer program or other method. The value of the slope is  $= 2.6823 \times 10^{-7}$  sec  $KP^{-1}$ , by using the first seven point (corr.=0.9993) instead of (corr.=0.9951) for all eight points.

**Table 5.3 : Theoretical values of B & C &  $\tau_p$  for Z'=Y and  
N=2.7**

$\tau_p(\text{sec } \times 10^{-11})$	B (G)	C (G)
0.688	0.2168	0.07376
0.888	0.2658	0.09757
2.880	0.6578	0.3502
4.880	1.042	0.6053
6.880	1.435	0.8599
8.880	1.833	1.115

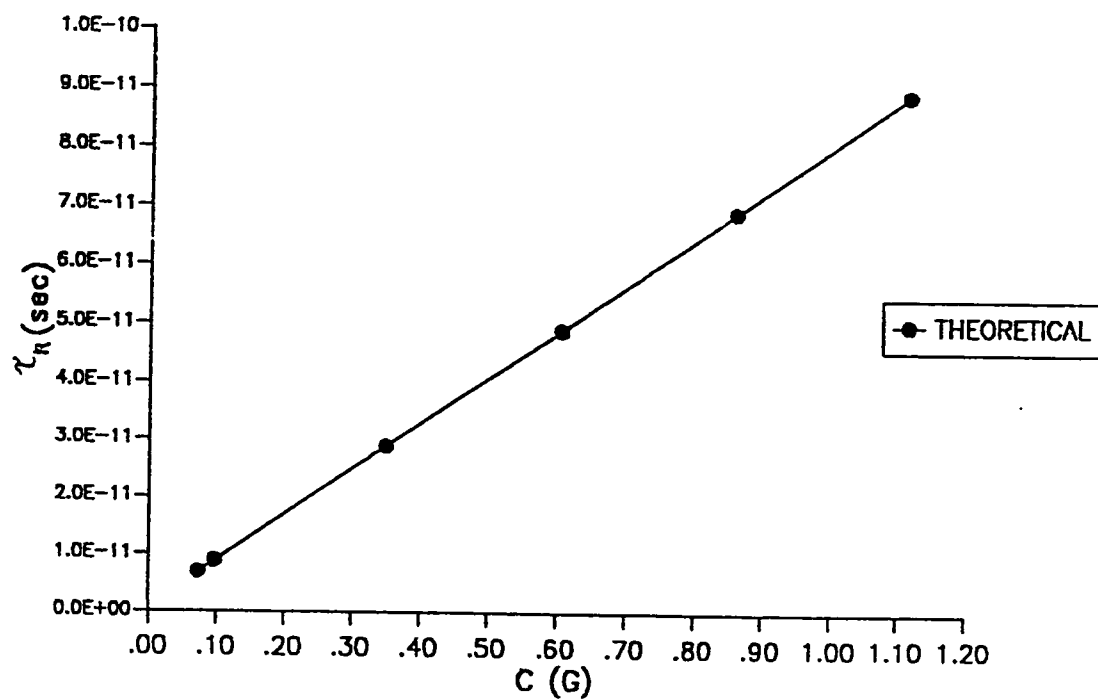


Fig. 5.11 :  $C$  Vs.  $\tau_r$  theoretical values

**Table 5.4 :  $\eta$  ,  $\eta/T$  and  $\tau_f$  Values for  $Cp_2VCl_2$  Experiment**

$t$ ( $^{\circ}C$ )	$\eta$ ( $P \times 10^{-2}$ )	$\eta/T(PK^{-1} \times 10^{-4})$	$\tau_f$ (Sec. $\times 10^{-11}$ )
21.7	0.57493	0.19500	0.87059
4.6	0.71951	0.25906	1.03664
-12.2	0.92823	0.35573	1.30709
-20.1	1.0622	0.41979	1.50128
-37.7	1.5030	0.63837	2.11845
-51.1	2.0777	0.93574	2.94947
-63.3	2.9862	1.4231	4.13035
-75.2	4.6942	2.3715	5.85921

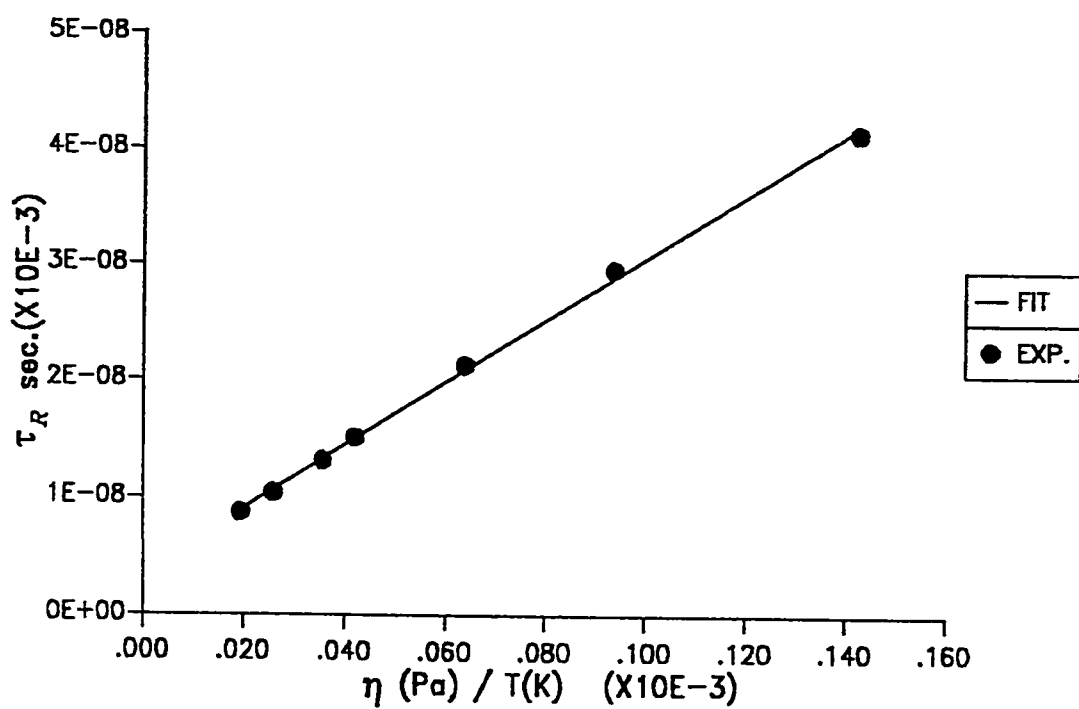


Fig. 5.12 :  $\eta/T$  Vs.  $\tau_R$  for  $\text{Cp}_2\text{VCl}_2$  in toluene.



### 5.3 : DETERMINATION OF THE TRANSLATIONAL DIFFUSION CONSTANT

#### 5.3.1 $\text{Cp}_2\text{VCl}_2$

The capillary diffusion experiment which described in chapter three is done for  $\text{Cp}_2\text{VCl}_2$  as mention. The experiment takes abut 22 days , from which 57 spectra are collected. The time is recorded for each run and each spectra is analysis to find the intensities of the eight lines (if the intensities dropped much at short time this give a result that high concentration difference between the solution in the capillary and the solvent above it, so the experiment will continue until a steady state start, then readjust the ESR spectrum parameter to get maximum intensities again Fig.5.13 show this effect). The data are introduce to the program " DIFF.FORTRAN " and the best least square fit is result form which the time and intensity for this spectra is substituted in eq.(1.13) to get the diffusion constant, and the value is recalculated by the fitting coefficients , also the ratio R ( $R = \ell_s / \ell_t$ ) is recalculated to check the experimental data. The line widths is calculated and adjusted by the following ,

$$\Delta\omega_i = \Delta\omega_o (I_o/I_i)^{\frac{1}{2}} \quad (5.16)$$

$\Delta\phi_0$  and  $I_0$  is the line width and intensity of a reference line and in  $\text{Cp}_2\text{VCl}_2$  it the line of  $M = -1/2$ ,  $\Delta\phi_i$  and  $I_i$  is the line width and intensity for a particular line, the average line widths is calculated and given by the program. The relative total intensity for each spectra and the time is in table 5.5 . In (Fig.5.14) some full spectra are shown. From the diffusion constant the program generate the fitting curve points (time Vs. intensity) which is plotted with the experimental values in (Fig. 5.15). In Appendix 1 the computer program, the input data and the output calculation is given for  $\text{Cp}_2\text{VCl}_2$  . The diffusion constant calculated by the DIFF. computer program for  $\text{Cp}_2\text{VCl}_2$  is  $= 0.520 \times 10^{-5} \text{ cm}^2/\text{s}$   
The value of  $R( = \ell_s / \ell_t )$  simulated is with in experimental error  $R(\text{exp})=0.57143$ ,  $R(\text{cal.})=0.5700$  .

### 5.3.2 VO[5-MeO-Sal SB](Phen)

For this compound the experiment goes for about 25 days and 62 spectra were taken , some of the spectra are shown in Fig.5.16. The full data are in appendix.1 as computer output also the calculation of total intensity and time is given . A plot of time vs. intensity is in Fig.5.17 with a simulation of the result, the value of diffusion constant obtained for VO[5-MeO-Sal SB](Phen) is  $= 0.655 \times 10^{-5} \text{ cm}^2/\text{sec}$  .  
The value of  $R( = \ell_s / \ell_t )$  simulated is with in experimental error  $R(\text{exp})=0.57143$ ,  $R(\text{cal.})=0.6800$  .

Table 5.5 : Time and Intensity of  $\text{Cp}_2\text{VCl}_2$  diffusion

p. #	Time(Sec)	Intensity	p. #	Time(Sec)	Intensity
* 1	$0.240 \times 10^3$	1.000	30	$0.860 \times 10^6$	0.451
2	$0.480 \times 10^5$	0.966	31	$0.901 \times 10^6$	0.439
3	$0.744 \times 10^5$	0.961	32	$0.927 \times 10^6$	0.416
4	$0.876 \times 10^5$	0.951	33	$0.940 \times 10^6$	0.407
5	$0.138 \times 10^6$	0.912	*34	$0.986 \times 10^6$	0.394
6	$0.153 \times 10^6$	0.909	35	$0.101 \times 10^7$	0.384
* 7	$0.173 \times 10^6$	0.893	36	$0.104 \times 10^7$	0.370
8	$0.212 \times 10^6$	0.862	37	$0.107 \times 10^7$	0.359
9	$0.241 \times 10^6$	0.826	38	$0.107 \times 10^7$	0.337
10	$0.259 \times 10^6$	0.831	39	$0.112 \times 10^7$	0.335
11	$0.297 \times 10^6$	0.806	*40	$0.117 \times 10^7$	0.312
12	$0.318 \times 10^6$	0.791	41	$0.119 \times 10^7$	0.298
13	$0.346 \times 10^6$	0.756	42	$0.121 \times 10^7$	0.307
*14	$0.381 \times 10^6$	0.751	43	$0.125 \times 10^7$	0.291
15	$0.409 \times 10^6$	0.718	44	$0.127 \times 10^7$	0.284
16	$0.431 \times 10^6$	0.706	45	$0.130 \times 10^7$	0.273
17	$0.469 \times 10^6$	0.687	46	$0.133 \times 10^7$	0.266
18	$0.498 \times 10^6$	0.657	47	$0.136 \times 10^7$	0.256
19	$0.512 \times 10^6$	0.651	*48	$0.138 \times 10^7$	0.244
20	$0.535 \times 10^6$	0.634	49	$0.142 \times 10^7$	0.238
*21	$0.579 \times 10^6$	0.615	50	$0.144 \times 10^7$	0.222
22	$0.604 \times 10^6$	0.590	51	$0.147 \times 10^7$	0.226
23	$0.667 \times 10^6$	0.556	52	$0.151 \times 10^7$	0.226
24	$0.690 \times 10^6$	0.543	53	$0.155 \times 10^7$	0.208
25	$0.727 \times 10^6$	0.522	54	$0.159 \times 10^7$	0.192
26	$0.758 \times 10^6$	0.510	55	$0.169 \times 10^7$	0.172
*27	$0.778 \times 10^6$	0.497	*56	$0.180 \times 10^7$	0.151
28	$0.815 \times 10^6$	0.472	57	$0.188 \times 10^7$	0.145
29	$0.843 \times 10^6$	0.466			

\*The spectra shown in Fig. 5.14 .

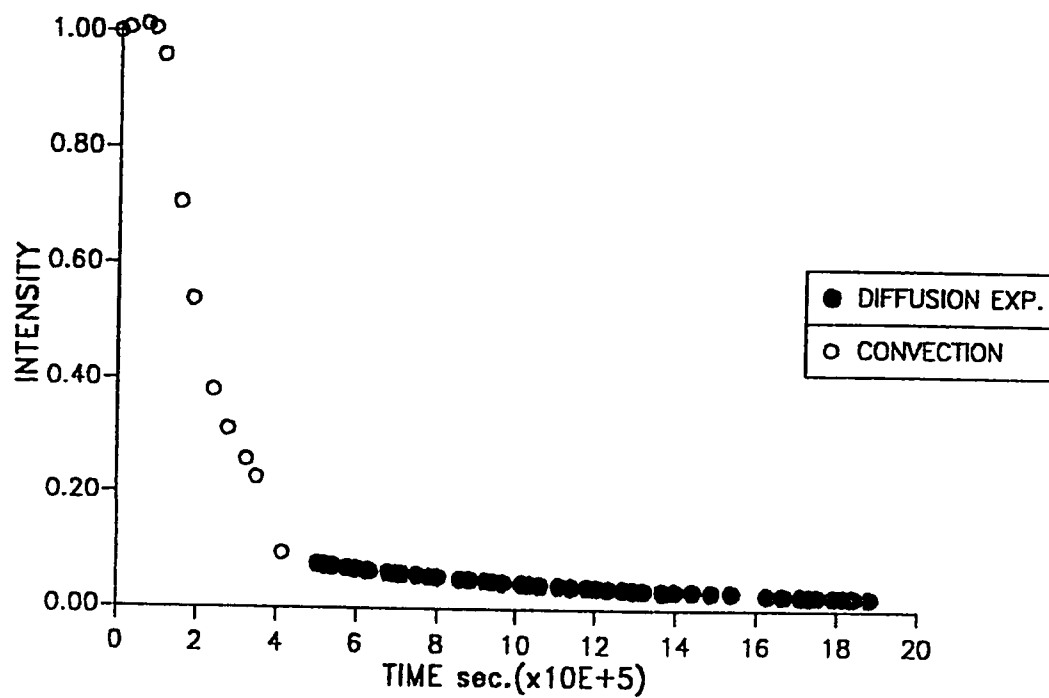


Fig. 5.13 : The effect of high concentration on diffusion.

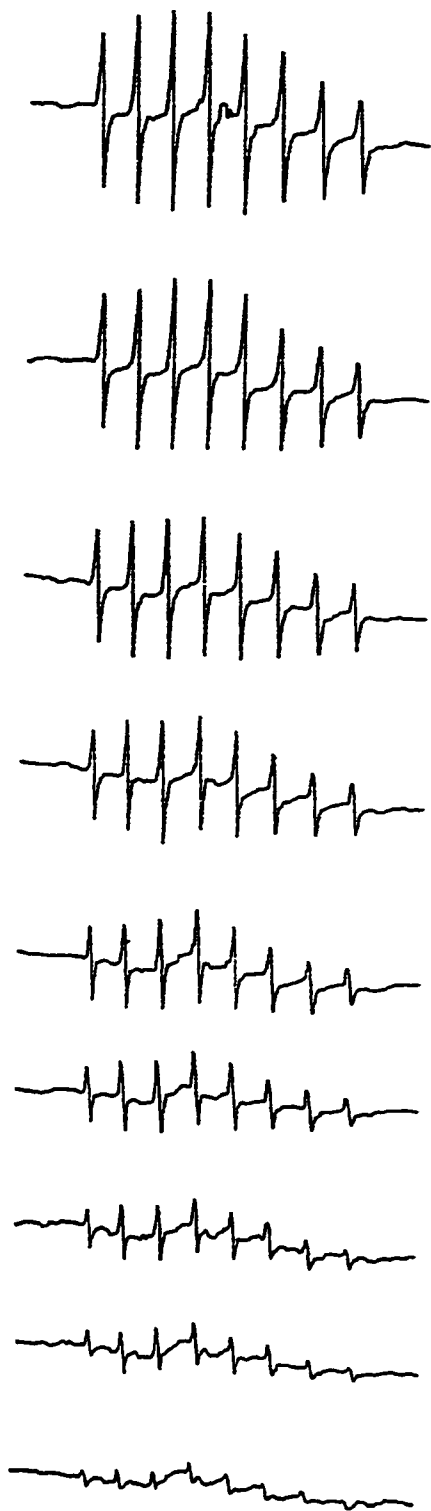


Fig. 5.14 : First derivative ESR spectra of  $\text{Cp}_2\text{VCl}_2$  at different time, the time and total intensity is marked (\*) in table 5.5 .

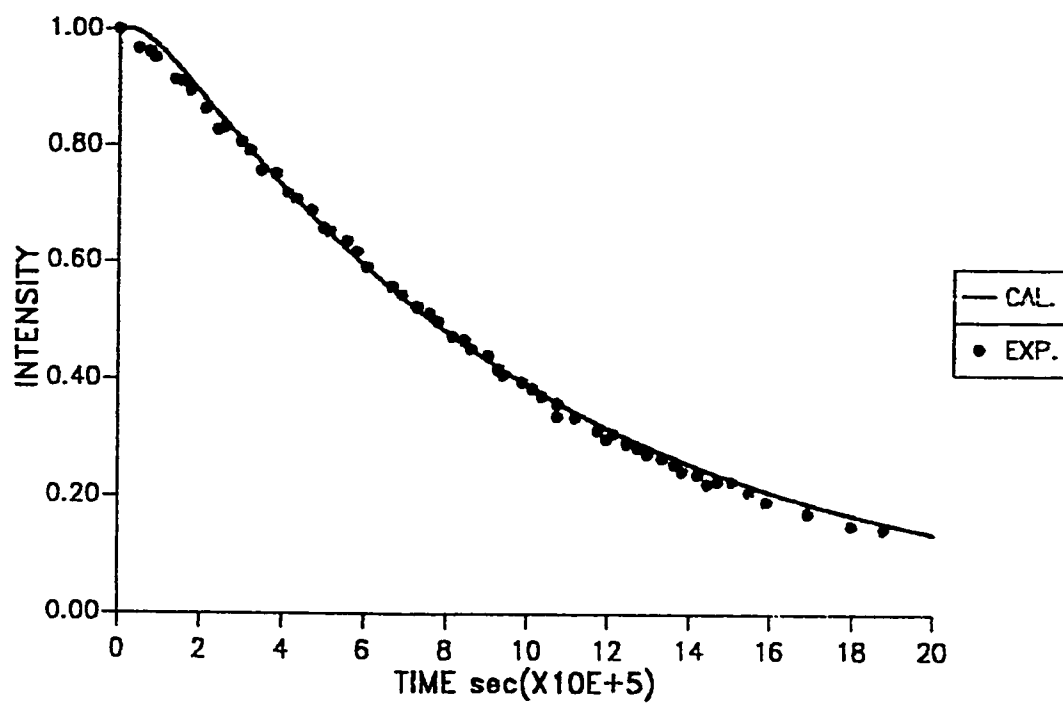


Fig. 5.15 : Total ESR intensities as a function of time for  
 $\text{Cp}_2\text{VCl}_2$

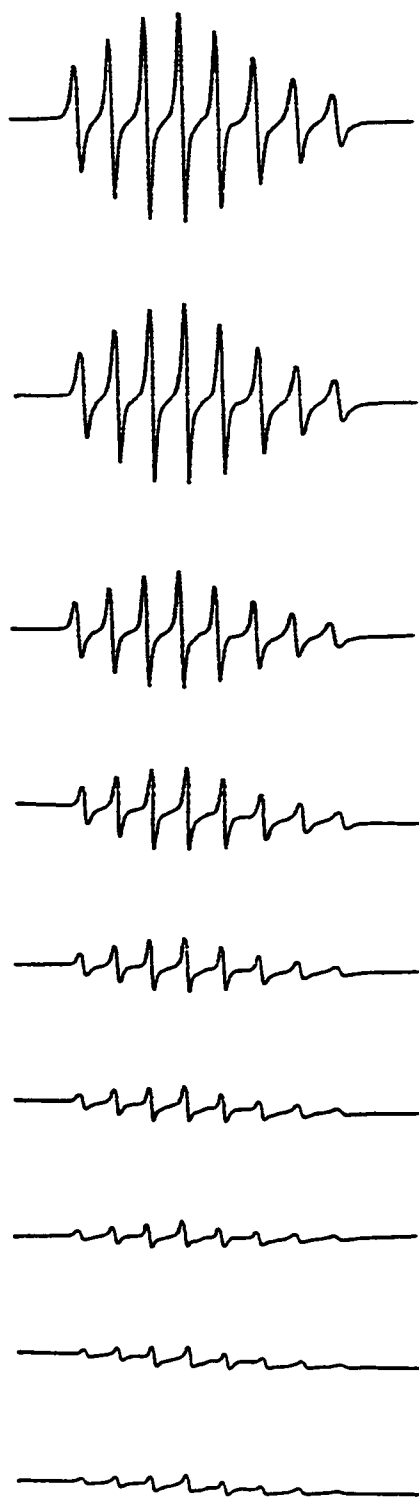


Fig. 5.16 : First derivative ESR spectra of  
VO[5-MeO-Sal SB](Phen) at different time.

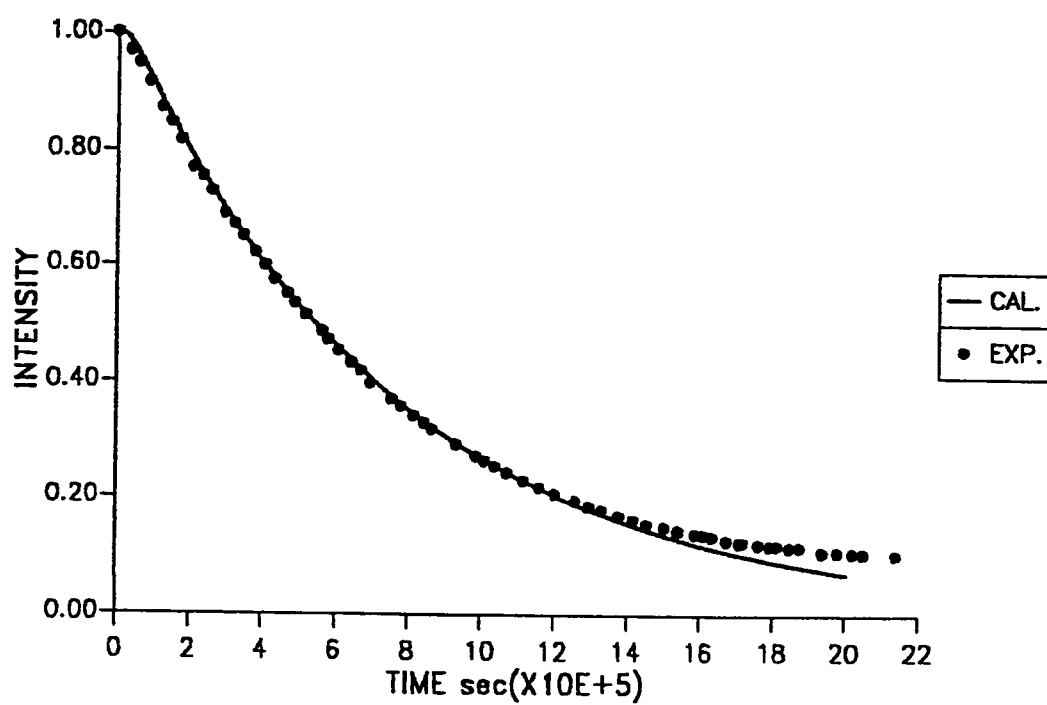


Fig. 5.17 : Total ESR intensities as a function of time for  
VO[5-MeO-Sal SB](Phen)



### 5.3.3 The Hydrodynamic radius

To calculate the hydrodynamic radius of these compound one can use Stokes-Einstein equation(1.1), by substituting the values obtained for the diffusion constant and the other parameter the unit for each are,  $D$  in  $\text{cm}^2/\text{sec}$ ,  $k$  in  $\text{JK}^{-1}$ ,  $T$  room temperature in K,  $\eta$  in  $\text{Pa}\cdot\text{sec}$ , by multiplying with a factor of  $(10^{14})$  the hydrodynamic radius results in ( $\text{\AA}$ ).

This hydrodynamic radius actually is the diameter of the molecules since it is rotating as sphere, so the radius of it will be half of the value obtained from the equation, to get the volume of the molecule by using the formula of the sphere volume  $(4 \pi r_0^3 / 3)$ . The value results are in the next table,

Compound	Hyd.radius( $\text{\AA}$ )	Radius( $\text{\AA}$ )	Volume( $\text{\AA}^3$ )
Compound I	7.820	3.910	250.0
Compound II	8.073	4.036	275.0

Compound I :  $\text{Cp}_2\text{VCl}_2$

Compound II :  $\text{VO}[5\text{-MeO-Sal SB}](\text{Phen})$

### 5.3.4 Simulating The Molecular Volume from Disk Top

#### Molecular Model [DTMM]

A computer program available for drawing the molecular structure and minimizing the energy of the molecules by iteration to minimize the torsional angle, Van der Waals forces, and other effect. Then it produce a structure that can be compare with the real molecule, looking at it from different angel and position and rotation on different axis carefully we can measure at least three lengths to determine the volume of this molecule, the structures are shown in Fig.5.18 for the two compound under study and a module compound vanadyl bisacetylacetonate or [VOAA], Fig.5.19 is the crystal structure of VOAA. The lengths obtained for the compounds are

Compound	$r_1(\text{\AA})$	$r_2(\text{\AA})$	$r_3(\text{\AA})$
Compound I	5.78	7.73	4.27
Compound II	8.6	13.5	6.4

Compound I :  $\text{Cp}_2\text{VCl}_2$ , Compound II :  $\text{VO}[5\text{-MeO Sal SB}](\text{Phen})$

where  $r_1$ ,  $r_2$ ,  $r_3$  are the length between two atoms in the same plane. Looking at the two tables above we can correlate between the hydrodynamic radius of the first compound (I) and the value of  $r_2$ , similarly for compound (II) the hydrodynamic radius is comparable with  $r_1$ , this give some indication that these compounds are rotating along these axis.

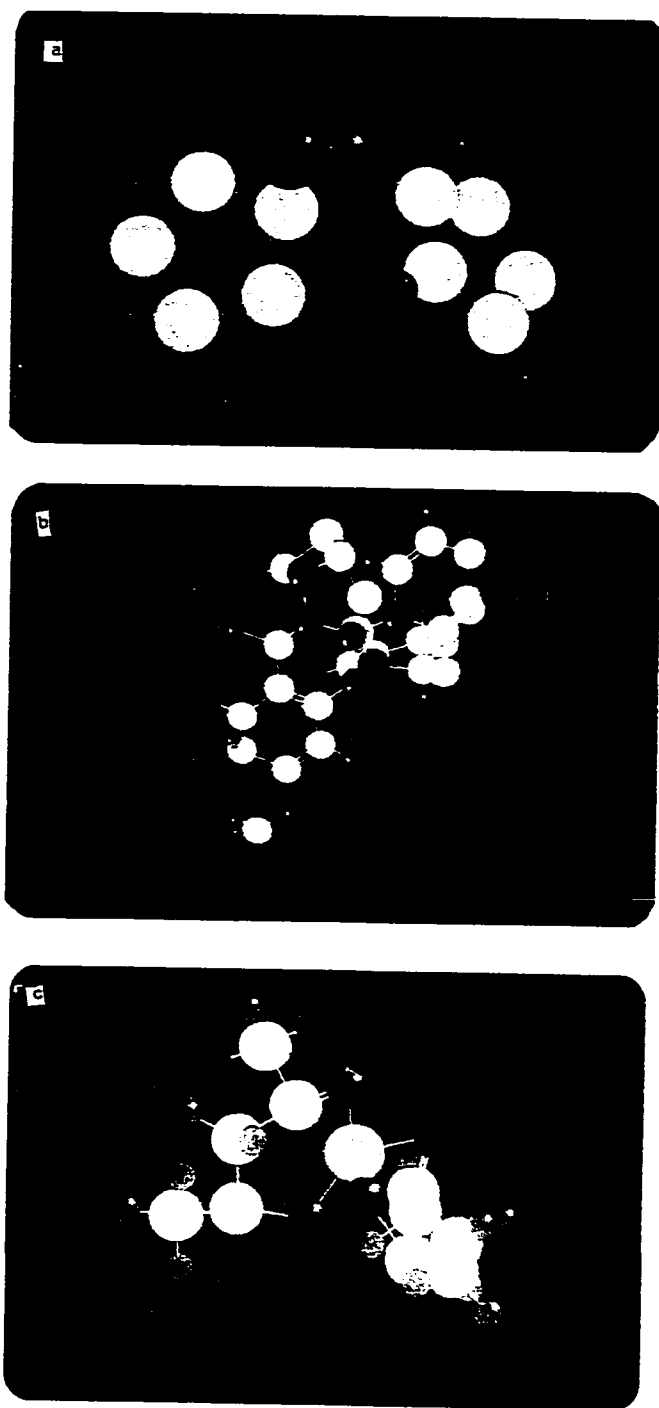


Fig. 5.18 : Structures by DTMM program, (a)  $\text{Cp}_2\text{VCl}_2$ ,  
(b)  $\text{VO}[5\text{-MeO Sal SB}](\text{Phen})$  , (c)  $\text{VOAA}$

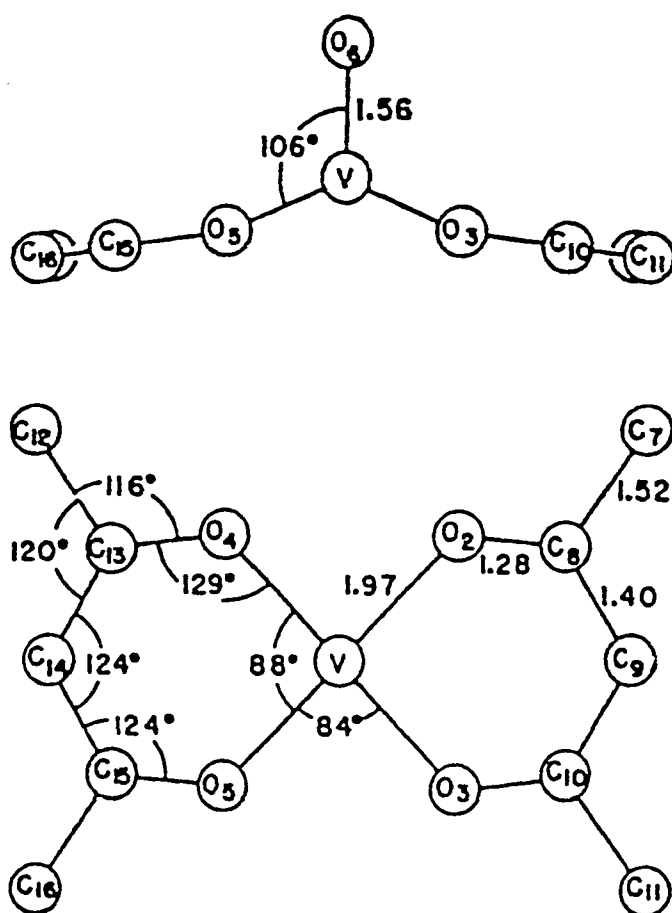
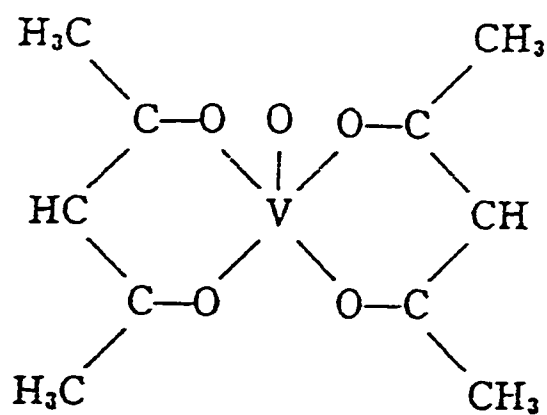


Fig. 5.19 : Crystal Structure of VOAA

#### 5.4 ANISOTROPIC INTERACTION PARAMETER ( $\kappa$ )

In previous discussion in section 5.2 for  $\text{Cp}_2\text{VCl}_2$  we end with the slope value results from plotting  $(\eta/T)$  Vs.  $\tau_p$  which are related by eq.(1.2), the slope is equal to

$$\text{The Slope} = \left(\frac{4}{3} \pi r_o^3\right) \left(\frac{\kappa}{k_B}\right)$$

to get  $\kappa$  we substitute by the volume in ( $\text{\AA}^3$ ), the slope in ( $\text{sec.Pa.K}^{-1}$ ), and  $k_B$  in ( $\text{JK}^{-1}$ ), then a factor of ( $10^{31}$ ) is multiplied to get  $\kappa$  as dimensionless unit,

$$\begin{aligned}\kappa &= (\text{slope}) (k_B) (10^{31}) / \text{volume} \\ &= (2.6823 \times 10^{-1})(1.381 \times 10^{-23})(10^{31}) / (250.0) \\ \kappa &= 0.150\end{aligned}$$

## 5.5 STICKINESS FACTOR (S)

We can analyzed the data for  $\text{Cp}_2\text{VCl}_2$  which we get from DTMM program in term of hydrodynamic free space model for molecular relaxation in liquids[63], a theory based on the existence of free space in the hydrodynamic continuum in which the molecule can rotate. We assign the major axis  $r_{\parallel} = 7.73$   $r_{\perp} = 5.78$  , the symmetry parameter  $\alpha^{\perp}$  is defined as  $(r_{\perp}/r_{\parallel})$ . Under stick boundary conditions, where the fluid at the surface of the rotating body sticks to it [68], the dimensionless rotational friction coefficient can be calculated from Perrin's formula [64],

$$f_{stick}^{\perp} = \frac{(2/3) [1 - (\alpha^{\perp})^4]}{\frac{[2 - (\alpha^{\perp})^2] (\alpha^{\perp})^2}{[1 - (\alpha^{\perp})^2]^{\frac{1}{2}}} \ln \left| \frac{1 + [1 - (\alpha^{\perp})^2]^{\frac{1}{2}}}{\alpha^{\perp}} \right| - (\alpha^{\perp})^2} \quad (5.17)$$

$\alpha^{\perp} = 0.747$  . This gives  $f_{stick}^{\perp} = 1.11$ . The corresponding value of the dimensionless rotational friction coefficient under slip boundary condition,  $f_{slip}^{\perp}$  , is calculated from

$$f_{slip}^{\perp} = f_{stick}^{\perp} [1 - (f_{stick}^{\perp})^{\frac{-2}{3}}] \quad (5.18)$$

the value of  $f_{slip}^{\perp} = 0.0761$

from the relation,

$$f_{slip}^{\perp} = f_{stick}^{\perp} C_{slip}^{hyd} \quad (5.19)$$

we can calculate

$$C_{slip}^{hyd} = 0.0761/1.111 = 0.0685$$

The stickiness factor,  $S$  is independent of molecular geometry and is zero in the slip limit and 1 in the stick limit. For truly hydrodynamic situation  $0 \leq S \leq 1$ , from eq.(1.3) we can calculate

$$S = (0.150 - 0.0685)/(1 - 0.0685) = 0.0875$$

## 5.6 STOKES-EINSTEIN MODEL

The anisotropy for rotation ( $N$ ), defined as  $N = R_{\parallel}/R_{\perp}$  is related to the geometric structure factor by means of the Stokes-Einstein equation[11,59,60],

$$R_i = \frac{k_B T}{(8\pi\eta r_i^3 \sigma_i)}, \quad i = \parallel \text{ or } \perp \quad (5.20)$$

where,

$$N = \frac{R_{\parallel}}{R_{\perp}} = \frac{\sigma_{\perp}}{\sigma_{\parallel}} \quad (5.21)$$

$$\sigma_{\parallel} = \frac{2}{3} \lambda^2 (1+\lambda^2) [(1+\lambda^2) \lambda^{-1} \tan^{-1} \lambda - 1]^{-1} \quad (5.22)$$

$$\sigma_{\perp} = \frac{2}{3} \lambda^2 (2+\lambda^2) [1 - (1-\lambda^2) \lambda^{-1} \tan^{-1} \lambda]^{-1} \quad (5.23)$$



and  $\lambda$  is a dimensionless structure factor,

$$\lambda = \frac{(r_{\parallel}^2 - r_{\perp}^2)^{\frac{1}{2}}}{r_{\perp}} \quad (5.24)$$

From the DTMM program data,

$$r_{\parallel} = 7.73$$

$$r_{\perp} = 4.28,$$

$$\lambda = 1.504$$

$$\sigma_{\parallel} = 0.0405$$

$$\sigma_{\perp} = 0.133$$

$$N = 3.28$$

which is within experimental error, with our observation that  
is  $N = 2.7 \pm 0.5$  .

## 5.7 THE ALLOWED - VALUES EQUATION (AVE)

An alternative approach to the determination of  $N$  is through the use of the allowed-values equation,

$$A \rho_x = B \rho_y + C \quad (5.25)$$

derived by Kowert [45,60] from line width coefficients  $B$  and  $c$  when nonsecular terms are unimportant, which shows that only one value of  $\rho_x = R_x/R_z$  is consistent with a  $C/B$  ratio obtained from experiment for a given  $\rho_y = R_y/R_z$ . If we let  $C = \gamma$  and  $B = \beta$ , then

$$A = (\Delta a - 3\delta a)[(\gamma/\beta)(\Delta g^- - 3\delta g^-) - (5/16)(\Delta a - 3\delta a)] \quad (a)$$

$$B = (\Delta a + 3\delta a)[(5/16)(\Delta a + 3\delta a) - (\gamma/\beta)(\Delta g^- + 3\delta g^-)] \quad (b)$$

$$C = 4\Delta a[(5/16)\Delta a - (\gamma/\beta)\Delta g^-] \quad (c)$$

with

$$\Delta a = a_z - (\frac{1}{2})(a_x + a_y) \quad (d)$$

$$\delta a = (\frac{1}{2})(a_x - a_y) \quad (e)$$

$$\Delta g^- = [g_z - (\frac{1}{2})(g_x + g_y)](\beta H_0/\hbar) \quad (f)$$

$$\delta g^- = (\frac{1}{2})(g_x - g_y)(\beta H_0/\hbar) \quad (g)$$

and

$$\rho_x = R_x/R_z \quad (h)$$

$$\rho_y = R_y/R_z \quad (i)$$

The values of the various expressions in Eqs. (a-g) using the magnetic parameters g's and a's of set (d) in table 5.1 are listed as follows (in Hz),

$$Aa = 1.426 \times 10^9$$

$$\delta a = 3.850 \times 10^8$$

$$Ag^- = 7.956 \times 10^8$$

$$\delta g^- = 1.609 \times 10^8$$

$$A = 2.834 \times 10^{16}$$

$$B = 8.572 \times 10^{16}$$

$$C = -2.040 \times 10^{17}$$

Using C/B value of 0.605 determined experimentally in table 5.2, the allowed-values equation is

$$2.834 \times 10^{16} \rho_x = 8.572 \times 10^{16} \rho_y - 2.040 \times 10^{17} \quad (j)$$

Dividing the equation by  $2.834 \times 10^{16}$  gives

$$\rho_x = 3.02 \rho_y - 7.20 \quad (k)$$

Equation (k) is not consistent with isotropic reorientation and only consistent with axial symmetry if  $\rho_x = 1$ ,  $\rho_y = 2.71$ . The allowed-values equation is very sensitive to variation in the  $\gamma/\beta$  value. The  $N$  obtained by using the allowed-values equation is  $N = 2.71 \pm 0.5$ . This result is consistent with the value determined from motional - narrowing results and the Stokes-Einstein model.

## 5.8 CONCLUSION

The translational diffusion constant of  $\text{Cp}_2\text{VCl}_2$  in chloroform has been measured by the capillary diffusion method at room temperature. The translational diffusion processes were monitored by electron spin resonance spectral intensities as a function of time. A computer program based on Fick's second law of diffusion for the capillary diffusion cell was written to get a single value of the diffusion constant. From the best least squares fit of the total intensities as a function of time a single value of the translational diffusion coefficient of  $0.520 \times 10^{-5} \text{ cm}^2/\text{s}$  was found for  $\text{Cp}_2\text{VCl}_2$  in chloroform measured at room temperature ( $25 \pm 1^\circ\text{C}$ ). Using Stokes-Einstein equation for translational diffusion, the hydrodynamic volume of  $\text{Cp}_2\text{VCl}_2$  in chloroform was found to be  $250 \text{ \AA}^3$ . And for  $\text{VO}[5\text{-MeO-Sal SB}](\text{Phen})$  the diffusion constant is  $0.655 \times 10^{-5} \text{ cm}^2/\text{s}$ , the volume is  $275 \text{ \AA}^3$ .

From an analysis of the temperature dependent ESR line width of  $\text{Cp}_2\text{VCl}_2$  in toluene in the motional narrowing region it was found that (1) of the four sets of the magnetic parameters for  $\text{Cp}_2\text{VCl}_2$  available in the literature, only one set of magnetic parameters by Petersen and Dahl could be used to explain the line width data, and (2)  $\text{Cp}_2\text{VCl}_2$  in toluene was undergoing anisotropic reorientation with  $N = 2.7 \pm 0.5$  at an axis  $z' = Y$ , where  $N$  is the ratio of  $R_{\parallel}/R_{\perp}$  and  $R_{\parallel}$  is the rotational diffusion constant along the molecular  $Z$  axis and

$R_1$  is the rotational diffusion constant perpendicular to the molecular Z axis. This conclusion is also supported by the analysis of anisotropic and nonsecular contributions to the ESR line widths using the allowed-values equation. The AVE analysis for  $\text{Cp}_2\text{VCl}_2$  in toluene gives  $\rho_x = 3.02 \rho_y - 7.20$  which is not consistent with isotropic reorientation ( $\rho_x = \rho_y = 1$ ) and only consistent with  $\rho_x = 1$ ,  $\rho_y = 2.71$ , agrees with our result. This value is also supported by the Stokes-Einstein model. The stickiness factor (S) calculated to be ( 0.0875 ) which is close to the slip limit (  $S = 0$  ).

The anisotropic interaction parameter  $\kappa$  which is a measure of the anisotropy of intermolecular interactions, was found to be 0.15 for  $\text{Cp}_2\text{VCl}_2$  in toluene.  $\kappa$  for VOAA in toluene is 0.53. The small  $\kappa$  value for  $\text{Cp}_2\text{VCl}_2$  in toluene could be due to the absence of oxygen atom at the apex of a five coordination square pyramidal structure present in VOAA, resulting in a more slippery boundary condition. Therefore, it seems that  $\kappa$  value decreases with decreasing solute corrugation. So we expect that  $\kappa$  for VO[ 5-MeO-Sal SB ](Phen) to be higher value.

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## **APPENDIX 1**

```

C THIS PROGRAM IS A SIMULATION FOR CALCULATION OF THE TRANSLATION DIF00010
C DIFFUSION CONSTANTS THE EQUATIONS ARE TAKEN FROM THE FOLLOWING DIF00020
C REFERENCE JOURNAL OF MAGNETIC RESONANCE 22,289-293(1976) " DIF00030
C BY THE TITEL " DIFFUSION COEFFICIENT OF PARAMAGNETIC SPECIES DIF00040
C IN SOLUTION " REPORTED BY " NYONG-KU AMM" . DIF00050
C THE CAPILLARY DIFFUSION METHOD ARE USED AND MONITORING THE DIF00060
C INTENSITIES BY ESR MACHIN AS A FUNCTION OF TIME . DIF00070
C NOTE :- DIF00080
C DIF00090
C 1)THE INPUT ARE :(FIRST READ STATMENT) DIF00100
C READ THE NAME OF THE COMPLUD DIF00110
C 2)THE INPUT ARE :(SECOND READ STATMENT) DIF00120
C SPIN=SPIN Q./ FOR THE ATOM IN THE SAMPLE DIF00130
C IDAY=THE NUMBER OF PLOTS TAKEN DIF00140
C REFIN= TOTAL INTINSESTY OF THE LINES IN THE REFERENCE SPECTROM DIF00150
C NFIT=THE WANTED FIT ORDER DIF00160
C LT= TOTAL LENGTH OF THE CAPILLARY DIF00170
C LS= LENGTH OF THE CAPILLARY IN THE CAVITY DIF00180
C THE INPUT ARE :(THAERED READ STATMENT) DIF00190
C EACH LINE FOR ONE SPECTRON DATA,THE ORDER OF INPUT DATA ARE DIF00200
C IRO=INTINSESTY OF A REFERENCE LINE IN THE SPECTROM DIF00210
C WLO=THE LINE WIEDTH OF THE REFERENCE LINE IN THE SPECTROM DIF00220
C T=TIME(SEC) OF THE PLOT TAKEN WITH RESPECT TO THE START OF EXP. DIF00230
C EI(K)=THE INTINSESTY OF ALL LINES IN THE SPECTROM DIF00250
C DIF00260
C 3) THE OUTPUT RESULTS ARE ON SEPARET TWO FIELDS DIF00270
C 4) THE PROGRAM FINED THE LEAST SQUIRE VALUE OF THE ALL PLOT AND DIF00280
C FROM WHICH IT CALCULATE THE EXP. VALUE OF THE DIF00290
C DIFFUSION COEFFICIENT. DIF00300
C 5) THE CALCULATED DIFFUSION COEFFICIENT VALUE CALCULATED BY CHANGING DIF00310
C THE R-VALUE ( R < 1 ) IN SUBROUTINE DIFSIM UNTIL YOU GET BEST FIT DIF00320
C 6) ON PLOT DATA FILE THE PROGRAM GENARETE THE CAL.INTINSESTY AND DIF00330
C THE TIME WHICH CAN BE USED AS DATA FOR PLOTING BY ICUSYS,THE DIF00340
C TIME IS SCALED BY DIVIDING IT BY 10ES FOR PLOTING ONLY DIF00350
C DIF00360
C ----- DIF00370
C WRITEN BY -----> NASAN,B.BALIDHOYOR DIF00380
C DIF00390
C ----- DIF00400
C DIF00410
C DIF00420
C DOUBLE PRECISION T,ISUM,X,Y,TM,SHDIF,MINSOR,VAR,STNDEV, DIF00430
C 1 ISM,DIF,YCAL,BB,R DIF00440
C CHARACTER NAME=30 DIF00450
C REAL LT,LS,IRO,IO DIF00460
C DIMENSION X(100),Y(100),EI(100),W(100),AVGW(100,20),BB(20), DIF00470
C 2 DIFSOR(100) DIF00480
C PI=3.1415926 DIF00490
C ACCUR=0.00001 DIF00500
C SCAL=100000 DIF00510
C ----- DIF00520
C FIRST READ----- DIF00530
C READ(5,16) NAME DIF00540
C 16 FORMAT(A30) DIF00550
C ----- DIF00560
C SECOND READ----- DIF00570
C READ(5,*) SPIN,IDAY,REFINT,NFIT,LT,LS DIF00580
C R=LS/LT DIF00590
C NP=IDAY DIF00600
C M=(2*SPIN)+1 DIF00610
C WRITE(6,3) DIF00620
C 3 FORMAT('1') DIF00630
C WRITE(7,123) DIF00640
C 123 FORMAT('1',' ','801',' ','2') DIF00650
C WRITE(7,321) IDAY DIF00660
C 321 FORMAT(14) DIF00670
C MM=1 DIF00680
C DO 66 L=1,IDAY DIF00690
C Z=-SPIN DIF00700
C ISUM=0 DIF00710
C ----- DIF00720
C READ STATMENT THREE ----- DIF00730
C READ(5,*) IRO,WLO,T,(EI(K),K=1,M) DIF00740
C IO=IRO/REFINT DIF00750
C DO 77 K=1,M DIF00760
C EI(K)=EI(K)/REFINT DIF00770
C ISUM=ISUM+EI(K). DIF00780

```

77	CONTINUE	DIF00760
	IF(I SUM.GT.1) GO TO 108	DIF00770
	JJ=488	DIF00780
	DO 118 N=1,N	DIF00790
	W(N)=WLO*((10/EI(N))**.5)	DIF00800
	AVGM(JJ,N)=W(N)	DIF00810
118	CONTINUE	DIF00820
	IF(I SUM.GT.1) GO TO 108	DIF00830
	NN=JJ+1	DIF00840
108	CONTINUE	DIF00850
	X(L)=T	DIF00860
	V(L)=I SUM	DIF00870
	IPL=L-1	DIF00880
	TSCAL=T/SCAL	DIF00890
	WRITE(7,100) TSCAL,I SUM	DIF00900
100	FORMAT(2(2X,F16.5))	DIF00910
66	CONTINUE	DIF00920
	N=-SPIN	DIF00930
	WRITE(6,15) NAME	DIF00940
15	FORMAT(/,'THE NAME OF THE SAMPLE IS ',A30)	DIF00950
	WRITE(6,5) N	DIF00960
5	FORMAT(/,'THE NUMBER OF LINES = ',I4)	DIF00970
	WRITE(6,102)	DIF00980
102	FORMAT(/,' ',,'THE SPIN G.# ',',')	DIF00990
	'THE AVERAGE LINE WIDTH ')	DIF01000
	WRITE(6,700)	DIF01010
700	FORMAT(' ')	DIF01020
	DO 88 K=1,N	DIF01030
	AVROM=0	DIF01040
	DO 99 N=1,JJ	DIF01050
	AVROM=AVROM+AVGM(N,K)	DIF01060
99	CONTINUE	DIF01070
	AVROM=AVROM/JJ	DIF01080
	WRITE(6,92) N,AVROM	DIF01090
92	FORMAT(12X,F6.1,7X,F16.5)	DIF01100
	N=N+1	DIF01110
88	CONTINUE	DIF01120
C		DIF01130
C	TO CAL. THE VAR = STW.DIV & LOWEST	DIF01140
C		DIF01150
	SHDIF=0	DIF01160
	-----CALLING SIMULATION TO FIND THE BEST FIT AND LEAST SQUARE	DIF01170
	-----VALUE	DIF01180
	CALL SIM(NP,NFIT,X,Y,BB)	DIF01190
	WRITE(6,607)	DIF01200
607	FORMAT(/,'PL.#',4X,'TIME.SEC',7X,'EXP.INT1',7X	DIF01210
	1,'CAL.INT1',7X,'SOUR.DIFF')	DIF01220
	WRITE(6,608)	DIF01230
608	FORMAT(' ')	DIF01240
	DO 60 K=1,NP	DIF01250
	ISM=0	DIF01260
	DO 70 N=1,NFIT	DIF01270
	YCAL=BB(N)*X(K)**(N-1)	DIF01280
	ISM=ISM+YCAL	DIF01290
70	CONTINUE	DIF01300
	DIF=(Y(K)-ISM)	DIF01310
	DIFSQ(K)=(Y(K)-ISM)**2	DIF01320
	SHDIF=SHDIF+DIFSQ(K)	DIF01330
	WRITE(6,60) K,X(K),Y(K),ISM,DIFSQ(K)	DIF01340
60	FORMAT(15,4D15.3)	DIF01350
60	CONTINUE	DIF01360
	MINSQR=DIFSQ(1)	DIF01370
	J=1	DIF01380
	DO 111 I=2,NP	DIF01390
	IF(DIFSQ(I).GT.MINSQR) GO TO 12	DIF01400
	MINSQR=DIFSQ(I)	DIF01410
	J=I	DIF01420
12	CONTINUE	DIF01430
111	CONTINUE	DIF01440
	WRITE(6,81) J	DIF01450
81	FORMAT(/,'THE LEAST SQUARES FIT IS FOR THE PLT.# = ',I3)	DIF01460
	VAR=(SHDIF/(N-1))	DIF01470
	STNDEV=VAR**.5	DIF01480
	WRITE(6,90) VAR	DIF01490
90	FORMAT(/,'THE V A R I A N C E = ',D15.5)	DIF01500
	WRITE(6,91) STNDEV	DIF01510
91	FORMAT(/,'THE STANDARD DEVIATION = ',D15.5)	DIF01520
	I SUM=Y(J)	DIF01530
	T=X(J)	DIF01540

```

C-----CALLING TRSDIFF TO CAL.THE EXP.DIFFUSION COEFF. -----
      CALL TRSDIFF (PI,T,LT,R,ACCUR,ISUM,SCAL,DIFF,RINT)
      WRITE(6,23)
23  FORMAT(//,'          TIME . SEC    INTENSITY    R-VALUE
1    DIFF.COEF    ')
      WRITE(6,24) T,ISUM,R,DIFF
24  FORMAT(//,' EXP.VAL ',E10.4,F13.6,D20.6,D18.6)
C-----CALLING DIFSIM TO GENERATE THE CALCULATED INTINSETY AND TIME
      CALL DIFSIM(R,LT,PI,ACCUR,SCAL,DIFF,BB)
      CALL TRSDIFF (PI,T,LT,R,ACCUR,ISUM,SCAL,DIFF,RINT)
      WRITE(6,996) T,RINT,R,DIFF
996  FORMAT(//,' CAL.VAL ',E10.4,F13.6,D20.6,D18.6)
      WRITE(6,2)
2    FORMAT('1')
      STOP
      END
C -----
      SUBROUTINE TRSDIFF (PI,T,LT,R,ACCUR,ISUM,SCAL,DIFF,RINT)
      DOUBLE PRECISION T,ISUM,R
      REAL LT
      DIFF=0.0
      K=100000000
      A=8/(PI*PI*R)
      DO 11 I=1,K
        B=0
        SUM=0
        DO 22 J=1,100
          EEE=((((2*N)+1)**2)*(PI**2)*DIFF*T)/(4*(LT**2))
          B=((EXP(-EEE))/((((2*N)+1)**2))
          IF (B.LE.ACCUR) GO TO 33
          FFF=((((2*N)+1)*PI*R)/2
          C=SIN(FFF)
          D=((-1)**N)*B*C
          SUM=SUM+D
          N=N+1
22      CONTINUE
33      CONTINUE
      RINT=A*SUM
      IF (RINT.LE.ISUM) GO TO 44
      DIFF=DIFF+0.1E-9
11      CONTINUE
44      CONTINUE
      RETURN
      END
C -----
C -----
      SUBROUTINE DIFSIM(R,LT,PI,ACCUR,SCAL,DIFF,BB)
      DOUBLE PRECISION BB,R
      DIMENSION BB(20)
      REAL LT
      WRITE(6,613)
C 613  FORMAT(//,'THIS IS SIMULATION DATA TIME(E+5) VS INTENSITIES ')
      WRITE(6,999)
999  FORMAT(' ')
      WRITE(7,300)
300  FORMAT(' ','R01')
C
C CHANGE R-VALUE TO GET BEST FIT
C
      R=0.5700+00
C
      T=0
      A=8/(PI*PI*R)
      DO 11 I=1,401
        B=0
        SUM=0
        DO 22 J=1,50
          EEE=((((2*N)+1)**2)*(PI**2)*DIFF*T)/(4*(LT**2))
          B=((EXP(-EEE))/((((2*N)+1)**2))
          IF (B.LE.ACCUR) GO TO 33
          FFF=((((2*N)+1)*PI*R)/2
          C=SIN(FFF)
          D=((-1)**N)*B*C
          SUM=SUM+D
          N=N+1
22      CONTINUE
33      CONTINUE
      RINT=A*SUM
      TSCAL=T/SCAL

```

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D1F01550
D1F01560
D1F01570
D1F01580
D1F01590
D1F01600
D1F01610
D1F01620
D1F01630
D1F01640
D1F01650
D1F01660
D1F01670
D1F01680
D1F01690
D1F01700
D1F01710
D1F01720
D1F01730
D1F01740
D1F01750
D1F01760
D1F01770
D1F01780
D1F01790
D1F01800
D1F01810
D1F01820
D1F01830
D1F01840
D1F01850
D1F01860
D1F01870
D1F01880
D1F01890
D1F01900
D1F01910
D1F01920
D1F01930
D1F01940
D1F01950
D1F01960
D1F01970
D1F01980
D1F01990
D1F02000
D1F02010
D1F02020
D1F02030
D1F02040
D1F02050
D1F02060
D1F02070
D1F02080
D1F02090
D1F02100
D1F02110
D1F02120
D1F02130
D1F02140
D1F02150
D1F02160
D1F02170
D1F02180
D1F02190
D1F02200
D1F02210
D1F02220
D1F02230
D1F02240
D1F02250
D1F02260
D1F02270
D1F02280
D1F02290
D1F02300
D1F02310
D1F02320
D1F02330
D1F02340
D1F02350

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      WRITE(7,100) TSCAL,RINT          DIF02360
100   FORMAT(2(AX,F16.5))              DIF02370
C     WRITE(6,101) TSCAL,RINT         DIF02380
C 101   FORMAT(2(AX,F16.5))            DIF02390
      T=T+(0.05*SCAL)                  DIF02400
11    CONTINUE                          DIF02410
      RETURN                             DIF02420
      END                                DIF02430
-----DIF02440
C                                           DIF02450
C                                           DIF02460
C     PROGRAM TO SOLVE SYSTEMS OF LINEARSIMUL- DIF02470
C     TANEOUS EQUATIONS USING IBM SUBROUTINE DGELS DIF02480
C                                           DIF02490
C     INPUT: N - NUMBER OF DATA POINTS       DIF02500
C           ND - MAXIMUM DEGREE OF POLYNOMIAL  DIF02510
C           (X(I),Y(I), I=1,N) - THE DATA POINTS
C-----DIF02520
C-----DIF02530
      SUBROUTINE SIM(NP,NFIT,X,Y,BB)        DIF02540
      DOUBLE PRECISION X,Y,B,AUX,POL,A,BB  DIF02550
      DIMENSION X(100),Y(100),B(20),A(210),AUX(20),BB(20) DIF02560
      EXTERNAL POL                           DIF02570
      NB=NP                                  DIF02580
      ND=NFIT                                 DIF02590
      M=2                                     DIF02600
150   MZ=M-1                               DIF02610
      IF(M.NE.NFIT) GOTO 202                DIF02620
      WRITE(6,55) M,MZ                      DIF02630
202   CONTINUE                              DIF02640
      CALL FIT(M,N,X,Y,A,B,AUX,POL)         DIF02650
      IF(M.NE.NFIT) GOTO 203                DIF02660
      WRITE(6,33)                            DIF02670
203   CONTINUE                              DIF02680
      DO 200 I=1,M                           DIF02690
        J=I-1                                DIF02700
        IF(M.NE.NFIT) GOTO 201               DIF02710
        BB(I)=B(I)                           DIF02720
200   WRITE(6,44) J,B(I)                     DIF02730
201   CONTINUE                              DIF02740
        M=M+1                                DIF02750
        IF(M.LE.(ND+1)) GO TO 150-          DIF02760
11    FORMAT(2I5)                            DIF02770
15    FORMAT(/,' THE # OF DATA POINTS IS ',I4,' AND THE MAXIMUM DEGREE OF DIF02780
      1 POLYNOMIAL FIT IS ',I3,/)           DIF02790
20    FORMAT('      I             X(I)             Y(I)',/) DIF02800
22    FORMAT(2D15.3)                         DIF02810
24    FORM='(I10,1PD15.3,1PD15.3)'          DIF02820
26    FORMAT(/,' THE COEFFICIENTS OF POLYNOMIAL ARE',/) DIF02830
28    FORM='( ,AX,'A' ,I1,' )=' ,1PD11.3)' DIF02840
30    FORMAT(/,' THE # OF DATA POINTS IS ',I3,AX,' THE DESIRED POLYNOMI DIF02850
      IAL FIT IS ',I2)                       DIF02860
      RETURN                                   DIF02870
      END                                    DIF02880
C                                           DIF02890
      SUBROUTINE FIT(M,N,X,Y,A,B,AUX,POL)   DIF02900
      DOUBLE PRECISION X,Y,B,AUX,POL,A     DIF02910
      DIMENSION X(1),Y(1),A(1),B(1),AUX(1) DIF02920
      ITRAN(1,J)=(J*(J-1))/2+1              DIF02930
      DO 100 K=1,M                           DIF02940
        B(K)=0.00                             DIF02950
        DO 100 I=1,N                           DIF02960
          B(K)=B(K)+Y(I)*POL(K,X(I))          DIF02970
100   CONTINUE                              DIF02980
        DO 200 J=1,M                           DIF02990
          DO 200 K=J,M                           DIF03000
            A(ITRAN(J,K))=0.00                 DIF03010
            DO 200 I=1,N                           DIF03020
              A(ITRAN(J,K))=A(ITRAN(J,K))+POL(K,X(I))*POL(J,X(I)) DIF03030
200   CONTINUE                              DIF03040
      CALL DGELS(B,A,M,1,0.0001,IER,AUX)    DIF03050
      IF(M.NE.NFIT) GOTO 202                DIF03060
      WRITE(6,11) IER                        DIF03070
11    FORMAT(/,' MESSAGE ERROR # FROM DGELS IS ',I3,/) DIF03080
202   CONTINUE                              DIF03090
      RETURN                                   DIF03100
      END                                    DIF03110
C                                           DIF03120
C                                           DIF03130

```



FUNCTION POL(I,X)	DIF03140
DOUBLE PRECISION X,POL	DIF03150
POL=X**2(I-1)	DIF03160
RETURN	DIF03170
END	DIF03180
	DIF03190
	DIF03200
	DIF03210
.....	DIF03220
SUBROUTINE DGELS	DIF03230
	DIF03240
	DIF03250
PURPOSE	DIF03260
TO SOLVE A SYSTEM OF SIMULTANEOUS LINEAR EQUATIONS WITH	DIF03270
SYMMETRIC COEFFICIENT MATRIX UPPER TRIANGULAR PART OF WHICH	DIF03280
IS ASSUMED TO BE STORED COLUMNWISE.	DIF03290
	DIF03300
USAGE	DIF03310
CALL DGELS(R,A,M,N,EPS,IER,AUX)	DIF03320
	DIF03330
DESCRIPTION OF PARAMETERS	DIF03340
R - DOUBLE PRECISION M BY N RIGHT HAND SIDE MATRIX	DIF03350
(DESTROYED). ON RETURN R CONTAINS THE SOLUTION OF	DIF03360
THE EQUATIONS.	DIF03370
A - UPPER TRIANGULAR PART OF THE SYMMETRIC DOUBLE	DIF03380
PRECISION M BY N COEFFICIENT MATRIX. (DESTROYED)	DIF03390
M - THE NUMBER OF EQUATIONS IN THE SYSTEM.	DIF03400
N - THE NUMBER OF RIGHT HAND SIDE VECTORS.	DIF03410
EPS - SINGLE PRECISION INPUT CONSTANT WHICH IS USED AS	DIF03420
RELATIVE TOLERANCE FOR TEST ON LOSS OF	DIF03430
SIGNIFICANCE.	DIF03440
IER - RESULTING ERROR PARAMETER CODED AS FOLLOWS	DIF03450
IER=0 - NO ERROR,	DIF03460
IER=-1 - NO RESULT BECAUSE OF M LESS THAN 1 OR	DIF03470
PIVOT ELEMENT AT ANY ELIMINATION STEP	DIF03480
EQUAL TO 0,	DIF03490
IER=K - WARNING DUE TO POSSIBLE LOSS OF SIGNIFI-	DIF03500
CANCE INDICATED AT ELIMINATION STEP K+1,	DIF03510
WHERE PIVOT ELEMENT WAS LESS THAN OR	DIF03520
EQUAL TO THE INTERNAL TOLERANCE EPS TIMES	DIF03530
ABSOLUTELY GREATEST MAIN DIAGONAL	DIF03540
ELEMENT OF MATRIX A.	DIF03550
AUX - DOUBLE PRECISION AUXILIARY STORAGE ARRAY	DIF03560
WITH DIMENSION M-1.	DIF03570
	DIF03580
REMARKS	DIF03590
UPPER TRIANGULAR PART OF MATRIX A IS ASSUMED TO BE STORED	DIF03600
COLUMNWISE IN M*(M+1)/2 SUCCESSIVE STORAGE LOCATIONS, RIGHT	DIF03610
HAND SIDE MATRIX R COLUMNWISE IN M*N SUCCESSIVE STORAGE	DIF03620
LOCATIONS. ON RETURN SOLUTION MATRIX R IS STORED COLUMNWISE	DIF03630
TOO.	DIF03640
THE PROCEDURE GIVES RESULTS IF THE NUMBER OF EQUATIONS M IS	DIF03650
GREATER THAN 0 AND PIVOT ELEMENTS AT ALL ELIMINATION STEPS	DIF03660
ARE DIFFERENT FROM 0. HOWEVER WARNING IER=K - IF GIVEN -	DIF03670
INDICATES POSSIBLE LOSS OF SIGNIFICANCE. IN CASE OF A WELL	DIF03680
SCALED MATRIX A AND APPROPRIATE TOLERANCE EPS, IER=K MAY BE	DIF03690
INTERPRETED THAT MATRIX A HAS THE RANK K. NO WARNING IS	DIF03700
GIVEN IN CASE M=1.	DIF03710
ERROR PARAMETER IER=-1 DOES NOT NECESSARILY MEAN THAT	DIF03720
MATRIX A IS SINGULAR, AS ONLY MAIN DIAGONAL ELEMENTS	DIF03730
ARE USED AS PIVOT ELEMENTS. POSSIBLY SUBROUTINE DGELS (WHICH	DIF03740
WORKS WITH TOTAL PIVOTING) WOULD BE ABLE TO FIND A SOLUTION.	DIF03750
	DIF03760
SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED	DIF03770
NONE	DIF03780
	DIF03790
METHOD	DIF03800
SOLUTION IS DONE BY MEANS OF GAUSS-ELIMINATION WITH	DIF03810
PIVOTING IN MAIN DIAGONAL, IN ORDER TO PRESERVE	DIF03820
SYMMETRY IN REMAINING COEFFICIENT MATRICES.	DIF03830
	DIF03840
.....	DIF03850

C	SUBROUTINE DGELS(R,A,M,N,EPS,IER,AUX)	DIF03860
C		DIF03870
C		DIF03880
C	DIMENSION A(1),R(1),AUX(1)	DIF03890
	DOUBLE PRECISION R,A,AUX,PIV,TB,TOL,PIVI	DIF03900
	IF(M)24,24,1	DIF03910
		DIF03920
C	SEARCH FOR GREATEST MAIN DIAGONAL ELEMENT	DIF03930
1	IER=0	DIF03940
	PIV=0.DO	DIF03950
	L=0	DIF03960
	DO 3 K=1,M	DIF03970
	L=L+K	DIF03980
	TB=DABS(A(L))	DIF03990
	IF(TB-PIV)3,3,2	DIF04000
2	PIV=TB	DIF04010
	I=L	DIF04020
	J=K	DIF04030
3	CONTINUE	DIF04040
	TOL=EPS*PIV	DIF04050
C	MAIN DIAGONAL ELEMENT A(I)=A(J,J) IS FIRST PIVOT ELEMENT.	DIF04060
C	PIV CONTAINS THE ABSOLUTE VALUE OF A(I).	DIF04070
C		DIF04080
C		DIF04090
C	START ELIMINATION LOOP	DIF04100
	LST=0	DIF04110
	NN=M*M	DIF04120
	LEND=M-1	DIF04130
	DO 18 K=1,M	DIF04140
		DIF04150
C		DIF04160
C	TEST ON USEFULNESS OF SYMMETRIC ALGORITHM	DIF04170
	IF(PIV)24,24,4	DIF04180
4	IF(IER)7,5,7	DIF04190
5	IF(PIV-TOL)6,6,7	DIF04200
6	IER=K-1	DIF04210
7	LT=J-K	DIF04220
	LST=LST+K	DIF04230
C		DIF04240
C	PIVOT ROW REDUCTION AND ROW INTERCHANGE IN RIGHT HAND SIDE R	DIF04250
	PIVI=1.DO/A(I)	DIF04260
	DO 8 L=K,NN,M	DIF04270
	LL=L+LT	DIF04280
	TB=PIVI*R(LL)	DIF04290
	R(LL)=R(L)	DIF04300
8	R(L)=TB	DIF04310
C		DIF04320
C	IS ELIMINATION TERMINATED	DIF04330
	IF(K-M)9,19,19	DIF04340
C		DIF04350
C	ROW AND COLUMN INTERCHANGE AND PIVOT ROW REDUCTION IN MATRIX A.	DIF04360
C	ELEMENTS OF PIVOT COLUMN ARE SAVED IN AUXILIARY VECTOR AUX.	DIF04370
9	LR=LST+(LT*(K+J-1))/2	DIF04380
	LL=LR	DIF04390
	L=LST	DIF04400
	DO 14 I=K,LEND	DIF04410
	L=L+1	DIF04420
	LL=LL+1	DIF04430
	IF(L-LR)12,10,11	DIF04440
10	A(LL)=A(LST)	DIF04450
	TB=A(L)	DIF04460
	GO TO 13	DIF04470
11	LL=L+LT	DIF04480
12	TB=A(LL)	DIF04490
	A(LL)=A(L)	DIF04500
13	AUX(I)=TB	DIF04510
14	A(L)=PIVI*TB	DIF04520
C		DIF04530
C	SAVE COLUMN INTERCHANGE INFORMATION	DIF04540
	A(LST)=I	DIF04550
C		DIF04560

C	ELEMENT REDUCTION AND SEARCH FOR NEXT PIVOT	
	PIV=0.00	D1F04570
	LLST=LST	D1F04580
	LT=0	D1F04590
	DO 18 I=K,LEND	D1F04600
	PIV=-ABS(I)	D1F04610
	LL=LLST	D1F04620
	LT=LT+1	D1F04630
	DO 15 LLD=I,LEND	D1F04640
	LL=LL+LLD	D1F04650
	L=LL+LT	D1F04660
	15 A(L)=A(L)+PIV*A(LL)	D1F04670
	LLST=LLST+1	D1F04680
	LR=LLST+LT	D1F04690
	TB=DABS(A(LR))	D1F04700
	IF(TB-PIV)17,17,16	D1F04710
	16 PIV=TB	D1F04720
	I=LR	D1F04730
	J=I+1	D1F04740
	17 DO 18 LR=J,MM,M	D1F04750
	LL=LR+LT	D1F04760
	18 R(LL)=R(LL)+PIV*R(LR)	D1F04770
C	END OF ELIMINATION LOOP	D1F04780
C		D1F04790
C		D1F04800
C	BACK SUBSTITUTION AND BACK INTERCHANGE	D1F04810
	19 IF(LEND)24,23,20	D1F04820
	20 I=M	D1F04830
	DO 22 J=2,M	D1F04840
	LST=LST-I	D1F04850
	I=I-1	D1F04860
	L=A(LST)+.500	D1F04870
	DO 22 J=1,MM,M	D1F04880
	TB=R(J)	D1F04890
	LL=J	D1F04900
	K=LST	D1F04910
	DO 21 LT=1,LEND	D1F04920
	LL=LL+1	D1F04930
	K=K+LT	D1F04940
	21 TB=TB-A(K)*R(LL)	D1F04950
	K=J+L	D1F04960
	R(J)=R(K)	D1F04970
	22 R(K)=TB	D1F04980
	23 RETURN	D1F04990
C		D1F05000
C		D1F05010
C	ERROR RETURN	D1F05020
	24 IER=-1	D1F05030
	RETURN	D1F05040
	END	D1F05050
		D1F05060

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CPZYCL2
3.5 57 11985 4 3.500 2.000
1869 6.2622 240 1442 1756 1869 1816 1678 1427 1102 895
1775 6.2622 48000 1381 1671 1775 1774 1563 1406 1149 840
1765 6.2622 74400 1438 1666 1765 1752 1562 1346 1118 852
1805 6.2622 87600 1410 1644 1805 1716 1567 1363 1080 809
1693 6.6536 137700 1343 1642 1693 1660 1554 1253 1031 756
1666 6.6536 153300 1330 1614 1666 1676 1533 1272 1012 791
1648 6.2622 173100 1313 1546 1648 1597 1485 1282 1027 770
1616 5.8708 211500 1272 1508 1616 1620 1461 1157 976 726
1504 6.2622 240600 1242 1438 1504 1537 1345 1185 928 722
1528 5.8708 259200 1205 1477 1528 1546 1345 1161 935 723
1409 6.2622 297300 1168 1391 1509 1467 1363 1152 911 694
1442 6.2622 318300 1158 1381 1442 1489 1321 1117 868 700
1425 6.2622 345600 1134 1356 1425 1363 1221 1053 816 687
1346 6.2622 381300 1120 1352 1346 1394 1235 1026 839 651
1302 6.1057 409200 1070 1262 1302 1309 1194 1006 784 674
1313 5.9491 430500 1055 1237 1313 1311 1172 1003 781 595
1304 5.6360 468600 1011 1201 1304 1256 1181 936 772 577
1196 6.1057 498000 985 1160 1196 1195 1129 885 741 580
1196 5.9491 511800 988 1176 1196 1190 1092 917 691 555
1121 6.1057 555000 946 1128 1121 1178 1021 917 726 564
1104 5.7926 579000 883 1092 1104 1149 1057 837 695 558
1151 5.6360 604200 880 1044 1151 1073 1003 786 625 507
1090 6.1057 666900 796 979 1090 997 924 783 620 480
1002 5.4795 690000 824 934 1002 1061 885 751 605 443
1013 5.6360 726900 800 911 1013 964 862 724 569 415
933 6.2622 757500 803 912 933 927 819 752 525 446
918 5.4795 777600 758 875 918 931 820 688 586 383
863 5.9491 815100 721 813 863 908 796 665 512 375
877 5.0098 843000 644 837 877 849 798 662 527 392
839 6.1057 859800 651 767 839 840 790 643 505 370
847 5.9491 900600 628 765 847 776 780 628 466 377
726 5.7926 928700 615 799 726 757 692 558 472 369
754 5.6360 939600 597 777 754 719 689 565 445 331
706 6.2622 986400 619 725 706 700 673 539 445 320
720 5.6360 1011600 549 693 720 712 683 500 438 304
695 5.1663 1035000 555 675 695 672 651 470 412 310
690 5.7926 1072800 514 642 690 649 623 481 388 310
669 5.7926 1072800 480 627 669 620 542 446 387 269
700 5.9491 1116900 453 563 700 588 544 478 361 283
609 5.4795 1174500 407 615 609 580 539 432 284 275
568 5.7926 1194900 426 563 568 563 496 380 332 244
543 5.1663 1211400 443 555 543 581 503 436 345 278
546 5.0098 1246200 468 497 546 518 495 378 340 287
510 5.7926 1272600 421 508 510 516 499 411 316 219
521 5.1663 1295700 403 492 521 491 444 406 287 222
445 5.7926 1332300 367 470 445 528 429 379 324 202
498 5.6360 1362300 365 408 498 500 417 376 264 244
477 5.3229 1380900 368 457 477 458 363 320 262 220
444 5.4795 1419900 318 439 444 455 405 348 249 198
421 5.1663 1443900 358 410 421 414 328 328 222 177
343 5.1663 1468500 335 363 343 435 387 343 246 213
423 5.1663 1505100 359 388 423 407 365 335 230 200
416 5.4795 1547700 313 357 416 370 365 257 252 162
343 5.4795 1592100 306 356 343 337 308 244 227 180
286 5.6360 1691500 261 313 286 319 301 230 209 144
292 6.4188 1797360 234 286 292 229 210 194 190 174
235 5.6360 1878360 224 237 235 279 247 212 163 138

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VO(5-MEO-SAL SB)(PHEN)

3.5 62 41684 5 3.500 2.000

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 7542 12.3940 36180 3841 5772 7288 7542 6381 4561 3069 1950  
 7828 12.3940 59580 3784 5637 7161 7828 6200 4466 3001 1892  
 7151 12.3940 87480 3647 5444 6874 7151 6010 4312 2881 1843  
 6854 12.3940 121980 3441 5191 6568 6854 5738 4104 2733 1733  
 6649 12.3940 147880 3345 5014 6376 6649 5600 3995 2648 1669  
 6410 12.3940 173580 3227 4872 6149 6410 5395 3859 2548 1604  
 6036 12.3940 208080 3051 4593 5799 6036 5091 3600 2409 1495  
 5914 12.3940 233880 2960 4511 5676 5914 4983 3552 2369 1505  
 5735 12.3940 259380 2871 4343 5515 5735 4793 3421 2279 1441  
 5423 12.3940 294780 2705 4098 5214 5423 4572 3254 2146 1343  
 5272 12.3940 321780 2632 4000 5087 5272 4423 3166 2084 1322  
 5078 12.3940 345780 2584 3896 4878 5078 4304 3052 2064 1248  
 4883 12.3940 379080 2435 3719 4710 4883 4095 2919 1938 1208  
 4719 12.3940 405780 2342 3569 4559 4719 3951 2810 1862 1159  
 4515 12.3940 432180 2269 3434 4355 4515 3771 2682 1812 1119  
 4339 12.3940 466680 2180 3265 4169 4339 3609 2584 1705 1071  
 4211 12.3940 487080 2101 3175 4032 4211 3531 2486 1659 1047  
 4034 12.3940 517980 2007 3066 3885 4034 3381 2406 1608 1018  
 3858 12.3940 562080 1879 2879 3678 3858 3227 2255 1516 915  
 3716 12.3940 577080 1843 2807 3578 3716 3100 2204 1472 919  
 3547 12.3940 604980 1767 2696 3434 3547 2976 2130 1436 877  
 3417 12.3940 640080 1673 2572 3283 3417 2869 2033 1318 841  
 3334 11.7417 665580 1604 2492 3169 3334 2800 1948 1267 817  
 3141 11.7417 692280 1528 2361 3020 3141 2629 1840 1221 773  
 2929 12.3940 751980 1422 2206 2831 2929 2468 1711 1154 701  
 2831 11.7417 777780 1373 2126 2733 2831 2377 1664 1110 689  
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 1785 11.7417 1159380 844 1287 1662 1785 1404 979 697 438  
 1660 11.7417 1201980 807 1216 1562 1660 1366 968 636 380  
 1575 11.7417 1257780 756 1158 1486 1575 1279 868 614 347  
 1451 12.3940 1296780 695 1121 1398 1451 1226 847 582 356  
 1426 11.7417 1330380 690 1057 1341 1426 1177 822 564 338  
 1146 11.7417 1376280 650 1032 1262 1146 1119 761 500 331  
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 1266 11.7417 1454280 590 925 1132 1266 998 701 484 300  
 1231 11.7417 1503480 556 855 1136 1231 964 682 475 302  
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 1047 11.7417 1611480 526 797 1044 1047 900 641 403 274  
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 987 11.7417 1676280 490 744 950 987 828 615 391 255  
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 979 12.3940 1720680 491 739 930 979 809 531 409 225  
 945 11.7417 1762380 448 717 907 945 769 589 359 219  
 917 11.7417 1794780 463 686 893 917 786 545 373 224  
 944 11.7417 1814280 449 697 883 944 768 541 375 230  
 921 11.7417 1849080 450 699 847 921 749 520 361 231  
 924 11.7417 1875480 456 681 885 924 740 513 356 225  
 845 11.0894 1938480 387 640 797 845 729 476 354 208  
 827 11.7417 1980780 423 638 804 827 721 482 337 205  
 870 11.7417 2021580 419 605 806 870 665 489 313 209  
 820 12.3940 2050080 386 623 784 820 692 494 310 224  
 813 11.7417 2140080 388 605 761 813 692 471 330 179

FILE: DIFFVL OUTPUT A1 KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAM

THE NAME OF THE SAMPLE IS CP2VCL2

THE NUMBER OF LINES = 8

THE SPIN 0.0 THE AVERAGE LINE WIDTH

-1.5 6.52724  
-2.5 5.95845  
-1.5 5.82001  
-0.5 5.85751  
0.5 6.18422  
1.5 6.72765  
2.5 7.52331  
3.5 8.53251

THE # OF DATA POINTS IS 57 THE DESIRED POLYNOMIAL FIT IS 3

THE COEFFICIENTS OF POLYNOMIAL ARE

A(0)= 1.0170+00  
A(1)= -7.8520-07  
A(2)= 1.3850-13  
A(3)= 1.8270-20

FILE: DIFFVL OUTPUT A1 KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAM

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PL.#	TIME, SEC	EXP. INTI	CAL. INTI	SQR. DIFF
1	0.2400+03	0.1000+01	0.1020+01	0.2690-03
2	0.4800+03	0.9640+00	0.9790+00	0.1710-03
3	0.7440+03	0.9610+00	0.9590+00	0.4780-03
4	0.8760+03	0.9510+00	0.9490+00	0.3320-03
5	0.1380+06	0.9120+00	0.9110+00	0.1020-05
6	0.1530+06	0.9090+00	0.9000+00	0.8920-04
7	0.1730+06	0.8930+00	0.8930+00	0.7310-04
8	0.2120+06	0.8620+00	0.8570+00	0.1070-04
9	0.2410+06	0.8260+00	0.8360+00	0.9630-04
10	0.2590+06	0.8310+00	0.8230+00	0.6990-04
11	0.2970+06	0.8060+00	0.7960+00	0.9470-04
12	0.3180+06	0.7910+00	0.7810+00	0.8810-04
13	0.3460+06	0.7560+00	0.7630+00	0.4870-04
14	0.1810+06	0.7510+00	0.7360+00	0.1630-03
15	0.4090+06	0.7180+00	0.7200+00	0.4280-03
16	0.4110+06	0.7060+00	0.7060+00	0.6290-06
17	0.4490+06	0.6870+00	0.6810+00	0.4140-04
18	0.4980+06	0.6570+00	0.6620+00	0.2930-04
19	0.5120+06	0.6510+00	0.6510+00	0.4870-03
20	0.5550+06	0.6340+00	0.6270+00	0.5830-04
21	0.5790+06	0.6150+00	0.6120+00	0.1180-04
22	0.6040+06	0.5970+00	0.5970+00	0.4800-04
23	0.6670+06	0.5600+00	0.5600+00	0.1220-04
24	0.6900+06	0.5430+00	0.5470+00	0.1570-04

THE LEAST SQUARES FIT IS FOR THE PLT.# = 34

THE VARIANCE = 0.687690-03

THE STANDARD DEVIATION = 0.2622ND-01

EXP. VAL	0.9846E+06	0.194410	0.5714290+00	0.5203900-03
THE	TIME, SEC	INTENSITY	R-VALUE	DIFF. COEF

EXP. VAL	0.9846E+06	0.194410	0.5714290+00	0.5203900-03
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1  
THE NAME OF THE SAMPLE IS VO(S-MEO-SAL 58)(PHEN)  
THE NUMBER OF LINES = 6

THE SPIN Q,  $\beta$  THE AVERAGE LINE WIDTH

-1.5	17.19795
-2.5	13.69622
-1.5	12.32136
-0.5	12.02574
0.5	13.19674
1.5	15.71477
2.5	19.20445
3.5	24.37782

THE  $\beta$  OF DATA POINTS IS 62 THE DESIRED POLYNOMIAL FIT IS 4

THE COEFFICIENTS OF POLYNOMIAL ARE

A(0)= 1.0150+00  
A(1)= -1.2700+06  
A(2)= 6.5180-13  
A(3)= -1.4120-19  
A(4)= 9.6000-27

PL.#	TIME, SEC	EXP. INTI	CAL. INTI	SQR. DIFF
1	0.4800+03	0.1000+01	0.1010+01	0.2100+03
2	0.3620+03	0.9690+00	0.9700+00	0.5010+06
3	0.5960+03	0.9480+00	0.9420+00	0.4340+04
4	0.8750+03	0.9160+00	0.9080+00	0.4390+04
5	0.1220+06	0.8720+00	0.8700+00	0.7360+05
6	0.4470+06	0.8470+00	0.8420+00	0.2750+04
7	0.1740+06	0.8170+00	0.8140+00	0.1340+04
8	0.2080+06	0.7690+00	0.7780+00	0.6910+04
9	0.2340+06	0.7550+00	0.7520+00	0.9380+05
10	0.2590+06	0.7270+00	0.7270+00	0.4750+05
11	0.2950+06	0.6900+00	0.6940+00	0.1550+04
12	0.1220+06	0.6710+00	0.6690+00	0.4500+05
13	0.3460+06	0.6500+00	0.6480+00	0.4430+05
14	0.3790+06	0.6220+00	0.6200+00	0.3050+05
15	0.4060+06	0.5990+00	0.5980+00	0.1510+05
16	0.4120+06	0.5750+00	0.5770+00	0.4380+05
17	0.4670+06	0.5500+00	0.5500+00	0.2300+06
18	0.4670+06	0.5340+00	0.5350+00	0.2830+05
19	0.5180+06	0.5140+00	0.5130+00	0.1620+06
20	0.5620+06	0.4850+00	0.4830+00	0.3300+05
21	0.5770+06	0.4710+00	0.4730+00	0.1760+05
22	0.6050+06	0.4530+00	0.4550+00	0.7310+05
23	0.6400+06	0.4320+00	0.4340+00	0.2980+05

24	0.6660+06	0.4180+00	0.4190+00	0.2480+06
25	0.6920+06	0.3960+00	0.4030+00	0.5390+04
26	0.7520+06	0.3700+00	0.3720+00	0.2420+05
27	0.7780+06	0.3580+00	0.3590+00	0.1020+05
28	0.8130+06	0.3410+00	0.3410+00	0.1350+07
29	0.8410+06	0.3290+00	0.3280+00	0.9460+07
30	0.8640+06	0.3190+00	0.3180+00	0.1820+07
31	0.9300+06	0.2920+00	0.2910+00	0.1240+05
32	0.9850+06	0.2710+00	0.2700+00	0.4660+06
33	0.1010+07	0.2630+00	0.2620+00	0.4120+06
34	0.1040+07	0.2540+00	0.2530+00	0.9610+06
35	0.1070+07	0.2430+00	0.2420+00	0.1020+05
36	0.1120+07	0.2290+00	0.2280+00	0.1100+05
37	0.1160+07	0.2170+00	0.2160+00	0.1900+05
38	0.1200+07	0.2060+00	0.2050+00	0.1760+05
39	0.1260+07	0.1940+00	0.1920+00	0.4960+05
40	0.1300+07	0.1840+00	0.1830+00	0.8000+06
41	0.1330+07	0.1780+00	0.1760+00	0.1990+05
42	0.1380+07	0.1680+00	0.1680+00	0.2740+06
43	0.1420+07	0.1610+00	0.1610+00	0.1230+10
44	0.1450+07	0.1530+00	0.1530+00	0.2580+05
45	0.1500+07	0.1490+00	0.1480+00	0.7670+06
46	0.1540+07	0.1430+00	0.1430+00	0.4430+08
47	0.1590+07	0.1370+00	0.1370+00	0.3030+07
48	0.1610+07	0.1360+00	0.1350+00	0.1060+05
49	0.1640+07	0.1330+00	0.1320+00	0.8250+06
50	0.1680+07	0.1260+00	0.1280+00	0.3810+05
51	0.1710+07	0.1220+00	0.1230+00	0.5920+05
52	0.1720+07	0.1230+00	0.1240+00	0.1020+05
53	0.1760+07	0.1190+00	0.1210+00	0.8070+05
54	0.1790+07	0.1170+00	0.1180+00	0.1210+05
55	0.1810+07	0.1170+00	0.1170+00	0.1130+06
56	0.1850+07	0.1150+00	0.1150+00	0.5090+08
57	0.1880+07	0.1150+00	0.1130+00	0.3220+05
58	0.1940+07	0.1060+00	0.1090+00	0.7750+05
59	0.1980+07	0.1040+00	0.1070+00	0.2590+06
60	0.2020+07	0.1050+00	0.1030+00	0.4380+08
61	0.2050+07	0.1040+00	0.1040+00	0.1570+06
62	0.2140+07	0.1020+00	0.9950+01	0.4670+05

THE LEAST SQUARES FIT IS FOR THE PLT.# = 43

THE VARIANCE = 0.118930+03

THE STANDARD DEVIATION = 0.109040+01

EXP. VAL	TIME, SEC	INTENSITY	R-VALUE	DIFF. COEF
0.10117E+07	0.160973	0.5714290+00	0.6763190+05	
0.10117E+07	0.160970	0.680000+00	0.6551420+05	

## **APPENDIX 2**



FILE: CPT2A OUTPUT AT KING FAND UNIVERSITY OF PETROLEUM AND MINERALS, DUMMAN.

THE CAL. B & C VALUES FOR 2°K USING STEWART & PORTER MAG. PAR FOR  
DIFF. VALUES IN M, THE FREQ. 9.511 GHZ

CHMR. TIME	B	C	
6. BAYLINE-12	.1641	.054078	
7. BAYLINE-12	.704407	.076332	
8. BAYLINE-11	.49759	.77455	N = 1
9. BAYLINE-11	.76448	.48481	
10. BAYLINE-11	1.04135	.69165	
11. BAYLINE-11	1.323	.89902	
SAME ORDER OF			
CHMR. TIME AS	.18137	.04914	
ABOVE	.177534	.043142	N = 2
	.61962	.2346	
	.6642	.41236	
	.8954	.50926	
	1.1319	.765434	
	.18429	.0430063	
	.15765	.054191	
	.3964	.203451	N = 3
	.594556	.340116	
	.79616	.5161	
	.99682	.67187	
	.1129	.03072	
	.10296	.09062	N = 4
	.2483	.181913	
	.5452	.32381	
	.7227	.4641	
	.904594	.604454	

FILE: CPT2A OUTPUT AT KING FAND UNIVERSITY OF PETROLEUM AND MINERALS, DUMMAN.

THE CAL. B & C VALUES FOR 2°K USING STEWART & PORTER MAG. PAR FOR  
DIFF. VALUES IN M, THE FREQ. 9.511 GHZ

CHMR. TIME	B	C	
6. BAYLINE-12	.1641	.05406	
7. BAYLINE-12	.7045	.07632	
8. BAYLINE-11	.4976	.77455	N = 1
9. BAYLINE-11	.76488	.48481	
10. BAYLINE-11	1.04135	.69165	
11. BAYLINE-11	1.323	.8999	
SAME ORDER OF			
CHMR. TIME AS	.2033	.078033	
ABOVE	.24982	.10315	N = 2
	.60655	.37201	
	.94185	.64567	
	1.35433	.91676	
	1.6921	1.18012	
	.22996	.094192	
	.28013	.12083	N = 3
	.69531	.483711	
	1.11695	.7677237	
	1.54647	1.04926	
	1.9693	1.31183	
	.25045	.100342	
	.304077	.14372	N = 4
	.77282	.501559	
	1.25132	.87196	
	1.72758	1.236812	
	2.2181	1.6937	

FILE: CPT2A OUTPUT AT KING FAND UNIVERSITY OF PETROLEUM AND MINERALS, DUMMAN.

THE CAL. B & C VALUES FOR 2°K USING STEWART & PORTER MAG. PAR FOR  
DIFF. VALUES IN M, THE FREQ. 9.511 GHZ

CHMR. TIME	B	C	
6. BAYLINE-12	.1646	.072322	
7. BAYLINE-12	.24416	.09515	
8. BAYLINE-11	.70547	.346568	N = 2
9. BAYLINE-11	.91419	.600999	
10. BAYLINE-11	1.29078	.854265	
11. BAYLINE-11	1.64822	1.10789	
SAME ORDER OF			
CHMR. TIME AS	.22173	.084227	
ABOVE	.27034	.11173	N = 3
	.67316	.39953	
	1.0161	.60997	
	1.4876	.919926	
	1.9615	1.27056	
	.24004	.0950754	
	.29187	.12594	N = 4
	.741164	.44344	
	1.1961	.767255	
	1.65832	1.0034	
	.26904	.11394	
	.3223	.15066	N = 6
	.65929	.370255	
	1.0033	.60726	
	1.39382	1.27991	
	.292532	.13034	N = 8
	.33403	.172593	
	.94185	.501485	
	1.58272	1.01794	
	.38101	.172024	N = 14
	.42443	.266775	
	1.2163	.78435	
	.363876	.183978	N = 16
	.4449	.241789	
	1.2896	.61351	
	.39293	.20594	N = 20
	.46293	.2701	
	1.425	.9041	
	.65412	.2531	N = 24
	.56418	.35043	
	1.7133	1.09995	

THE CAL. B & C VALUES FOR  $Z' = X$  USING BAKALIK & HAYES MAG. PAR FOR  
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88710E-12	.142395	.069719	
8.88710E-12	.17748	.091617	
2.44710E-11	.431849	.33435	N = 1
4.88710E-11	.66343	.58424	
6.88710E-11	.903776	.832302	
8.88710E-11	1.1484	1.0795	
SAME ORDER OF	.13581	.079244	
CORR.TIME AS	.16832	.104483	
ABOVE	.43074	.37152	N = 4
	.6671	.64568	
	.90598	.919643	
	1.14861	1.19354	
	.13576	.094162	
	.16755	.12409	
	.456016	.42868	N = 8
	.7201	.740543	
	.98047	1.05478	
	1.2426	1.37079	

THE CAL. B & C VALUES FOR  $Z' = Y$  USING BAKALIK & HAYES MAG. PAR FOR  
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88710E-12	.140063	.0602662	
8.88710E-12	.171574	.0796071	
2.44710E-11	.4170408	.2880488	N = 2
4.88710E-11	.6629788	.4991402	
6.88710E-11	.9159114	.709363	
8.88710E-11	1.171911	.9197114	
SAME ORDER OF	.181371	.0808904	
CORR.TIME AS	.2189958	.1072044	
ABOVE	.5358927	.3789803	N = 4
	.9071857	.6522152	
	1.264468	.925948	
	1.624466	1.201343	
	.2507477	.1261714	
	.3030084	.1645478	N = 10
	.8414513	.5683533	
	1.402059	.9735658	
	.3792556	.221347	
	.470338	.289466	N = 30
	1.425019	.9633199	
	.6403737	.4096137	
	.8137168	.5374634	N = 100
	2.579395	1.770348	

THE CAL. B & C VALUES FOR  $Z' = Z$  USING BAKALIK & HAYES MAG. PAR FOR  
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88710E-12	.14239	.069719	
8.88710E-12	.17748	.091617	
2.44710E-11	.43185	.33435	N = 1
4.88710E-11	.66383	.58424	
6.88710E-11	.903776	.832302	
8.88710E-11	1.1484	1.07947	
SAME ORDER OF	.13586	.064047	
CORR.TIME AS	.168377	.08414	
ABOVE	.43088	.30049	N = 4
	.66732	.52604	
	.906342	.75144	
	1.14908	.976399	
	.13584	.066919	
	.16765	.087896	
	.45627	.30514	N = 8
	.72055	.53168	
	.98111	.76034	
	1.2434	.98986	

FILE: CP2XC OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR  $Z'=X$  USING CASEY & RAYNOR MAG. PAR. FOR  
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.20938	.051146	
8.88X10E-12	.26096	.067211	
2.88X10E-11	.63499	.24529	N = 1
4.88X10E-11	.97608	.42869	
6.88X10E-11	1.3289	.61084	
8.88X10E-11	1.6886	.7925	
SAME ORDER OF	.1457807	.03365571	
CORR.TIME AS	.1845777	.04388695	
ABOVE	.4753854	.1581491	N = 4
	.705093	.280999	
	.9357591	.4038153	
	1.172129	.5260099	
	.1118138	.02519779	
	.1425162	.03273453	
	.3889812	.1155009	N = 8
	.5730618	.2068584	
	.7467171	.2996174	
	.9226121	.3921521	

FILE: CP22C OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES AT  $Z'=Z$  USING CASEY & RAYNOR MAG. PAR. FOR  
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.20938	.051146	
8.88X10E-12	.260964	.067211	
2.88X10E-11	.63499	.245297	N = 1
4.88X10E-11	.9761	.428694	
6.88X10E-11	1.3289	.61084	
8.88X10E-11	1.68859	.7924662	
SAME ORDER OF	.3298447	.09292165	
CORR.TIME AS	.3994029	.1232105	N = 4
ABOVE	1.014741	.4352556	
	1.64664	.748131	
	.4116183	.1314404	
	.4971032	.1738935	N = 8
	1.350281	.5976089	

FILE: CP2YC OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR  $Z'=Y$  USING CASEY & RAYNOR MAG. PAR FOR  
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.20938	.051146	
8.88X10E-12	.260964	.067211	
2.88X10E-11	.63499	.2453	N = 1
4.88X10E-11	.9761	.4287	
6.88X10E-11	1.3289	.61084	
SAME ORDER OF	.2947512	.0866697	
CORR.TIME AS	.358445	.1148427	N = 4
ABOVE	.911908	.4060238	
	1.467125	.698853	
	.3544574	.1202329	
	.4294979	.1590027	N = 8
	1.1670	.5467515	
	1.915462	.9349394	

FILE: CP2XD OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'=X USING PETERSEN & DAHL MAG.PAR AT  
DIFF. VALUES OF N THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.17439	.052638	
8.88X10E-12	.21736	.069172	
2.88X10E-11	.52889	.252457	N = 1.0
4.88X10E-11	.81300	.441216	
6.88X10E-11	1.10687	.6287	
8.88X10E-11	1.40645	.81567	
SAME ORDER OF	.123124	.0349899	
CORR.TIME AS	.155723	.0456389	
ABOVE	.400941	.164411	N = 4.0
	.595986	.291979	
	.791918	.419518	
	.992593	.546438	
	.0959023	.0265643	
	.121981	.0345285	
	.332873	.121737	N = 8.0
	.492641	.217747	
	.6438824	.315219	
	.796998	.412501	

FILE: CP2YD OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'=Y USING PETERSEN & DAHL MAG.PAR AT  
DIFF. VALUES OF N THE FREQ.= 9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.174392	.052638	
8.88X10E-12	.21736	.069172	N = 1
2.88X10E-11	.52889	.252457	
4.88X10E-11	.812997	.441216	
6.88X10E-11	1.10687	.5287	
8.88X10E-11	1.40646	.81567	
SAME ORDER OF	.21706	.073773	
CORR.TIME AS	.26607	.097589	
ABOVE	.65864	.35019	N = 2.7
	1.04358	.605337	
	1.437498	.859878	
	.282285	.119276	
	.342421	.157710	N = 8
	.930494	.542482	
	1.52378	.930171	

FILE: CP2ZD OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'=Z USING PETERSEN & DAHL MAG.PAR FOR  
DIFF. VALUES OF N , THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.17439	.052638	
8.88X10E-12	.21736	.0691718	
2.88X10E-11	.528891	.252457	N = 1
4.88X10E-11	.812997	.441216	
6.88X10E-11	1.10686	.6287	
8.88X10E-11	1.406455	.8156679	
SAME ORDER OF	.280983	.097772	
CORR.TIME AS	.3399639	.1296689	N = 4
ABOVE	.86351	.497974	
	1.403485	.78691	
	.353024	.139111	
	.4261	.18406	N = 8
	1.157812	.632495	
	1.91514	1.082584	

FILE: CP2V27AL OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR  $Z'=Y$  AT  $N=2.7$  USING DIFFERANTE MAG. PAR

CORR. TIME	B	C	
6.88X10E-12	.21542	.081018	
8.88X10E-12	.26338	.107161	
2.88X10E-11	.6.111	.384558	STEWART & FORT
4.88X10E-11	1.03699	.6648	
6.88X10E-11	1.431913	.944341	
SAME ORDER	.15701	.068077	
CORR. TIME	.19095	.090098	
AS ABOVE	.470725	.323174	BAKALIK & HAYES
.	.75775	.558183	
.	1.05155	.792746	
8.88X10E-11	1.34814	1.02792	
	.267064	.073245	
	.32698	.09492	
	.80891	.34763	CASEY & RAYBOR
	1.28475	.60063	
	1.7717	.85295	
	.21706	.073773	
	.26607	.097589	
	.65864	.35019	PETERSEN & DAHL
	1.04358	.605337	
	1.437498	.859878	

FILE: CP2VDFRE OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR  $Z'=Y$  AND  $N=2.7$  AT DIFF. FREQ USING MAG. PAR FOR PETERSEN & DAHL

CORR. TIME	B	C	
6.88X10E-12	.2131	.0734264	
8.88X10E-12	.261461	.097368	
2.88X10E-11	.64653	.34979	AT 9.3 GHZ
4.88X10E-11	1.0225	.60501	
6.88X10E-11	1.4072	.85962	
SAME ORDER	.21498	.073696	
OF CORR. TIME	.26365	.097473	
AS ABOVE FOR	.65227	.349982	AT 9.4 GHZ
ALL SETS	1.03247	.605168	
	1.4216	.85975	
	.21706	.073773	
	.26607	.097589	
	.65864	.35019	AT 9.511 GHZ
	1.04358	.605337	
	1.437498	.859878	
	.21872	.0738345	
	.268003	.097682	
	.663739	.350357	AT 9.6 GHZ
	1.05249	.605468	
	1.45028	.85998	
	.22058	.0739035	
	.270165	.0977856	
	.66951	.350538	AT 9.7 GHZ
	1.06251	.605612	
	1.4646	.860094	